

ON THE STRONG ATTRACTION LIMIT FOR A CLASS OF NONLOCAL INTERACTION ENERGIES

ALMUT BURCHARD, RUSTUM CHOKSI, AND ELIAS HESS-CHILDS

ABSTRACT. This note concerns the problem of minimizing a certain family of non-local energy functionals over measures on \mathbb{R}^n , subject to a mass constraint, in a strong attraction limit. In these problems, the total energy is an integral over pair interactions of attractive-repulsive type. The interaction kernel is a sum of competing power law potentials with attractive powers $\alpha \in (0, \infty)$ and repulsive powers associated with Riesz potentials. The strong attraction limit $\alpha \rightarrow \infty$ is addressed via Gamma-convergence, and minimizers of the limit are characterized in terms of an isodiametric capacity problem. We also provide evidence for symmetry-breaking in high dimensions.

1. INTRODUCTION AND STATEMENT OF THE RESULTS

We consider mass-constrained variational problems of the form

$$\begin{cases} \text{Minimize} & \mathcal{E}_{\alpha,\lambda}(\mu) := \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} K_{\alpha,\lambda}(x-y) d\mu(x) d\mu(y) \\ \text{over } \mathcal{P} := & \{\mu \text{ Borel measure on } \mathbb{R}^n : \mu(\mathbb{R}^n) = 1\}, \end{cases} \quad (1)$$

where the interaction kernel is given by

$$K_{\alpha,\lambda}(x-y) := |x-y|^\alpha + |x-y|^{-\lambda} \quad \text{with } \alpha \in (0, \infty) \text{ and } \lambda \in (0, n). \quad (2)$$

These kernels are strongly repulsive at short range, with the repulsion controlled by the exponent λ , and attractive at long range, with the attraction controlled by α , see Fig. 1. Since the kernels are lower semicontinuous, locally integrable, and grow at infinity, by the results of [15, 6], Problem (1) has a global minimizer.

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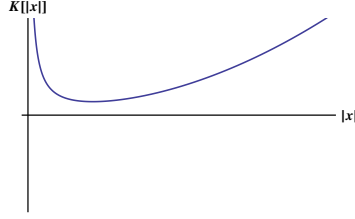


FIGURE 1. Shape of the interaction kernel $K_{\alpha,\lambda}(|\cdot|)$ for $\alpha \in (0, \infty)$ and $\lambda \in (0, n)$.

Variational problems of the form (1) arise in connection with a class of models for aggregation and self-assembly that have recently received much attention (see for example, [2] and the references therein). In those models, a population density ρ evolves according to the equation

$$\rho_t + \nabla \cdot (\rho \mathbf{v}) = 0, \quad \mathbf{v} = -\nabla K_{\alpha,\lambda} * \rho,$$

which is the gradient flow of the energy $\mathcal{E}_{\alpha,\lambda}(\mu)$ on absolutely continuous measures $\mu = \rho dx$ in the 2-Wasserstein metric (cf. [5]). Energy minimizers represent stable steady-states of the aggregation process.

Here, we study the minimization problem (1) in the strong attraction regime where $\alpha \rightarrow \infty$. In this limit, finite energy alone restricts the support of a measure to have diameter no larger than one.

In Fig. 2, we present a few particle simulations in dimension $n = 2$ which suggest that as α increases, minimizers concentrate on the boundary of the ball of diameter 1 for some values of λ ; but spread out (non-uniformly) over the ball for larger values of λ . A broader range of behaviour is expected for other parameters and in higher dimensions (see for example Fig. 3).

Our first result is that in the limit as $\alpha \rightarrow \infty$, Problem (1) approaches the problem of minimizing

$$\mathcal{E}_{\infty,\lambda}(\mu) := \begin{cases} \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} |x - y|^{-\lambda} d\mu(x) d\mu(y) & \text{if } \text{diam}(\text{supp } \mu) \leq 1 \\ +\infty & \text{otherwise} \end{cases} \quad (3)$$

over \mathcal{P} . The limit is understood in the sense of Gamma-convergence.

Theorem 1 (Strong attraction limit). *Let $\lambda \in (0, n)$. Then $\mathcal{E}_{\alpha,\lambda} \xrightarrow{\Gamma} \mathcal{E}_{\infty,\lambda}$ as $\alpha \rightarrow \infty$ in the weak topology of measures.*

The limiting problem admits a solution:

Theorem 2 (Existence). *The functional $\mathcal{E}_{\infty,\lambda}$ has a global minimizer in \mathcal{P} .*

The proofs of Theorems 1 and 2 are presented in Section 3.

Remark. In the literature, the interaction kernel is sometimes normalized to

$$\tilde{K}_{\alpha,\lambda}(x - y) := \frac{1}{\alpha} |x - y|^\alpha + \frac{1}{\lambda} |x - y|^{-\lambda}, \quad (4)$$

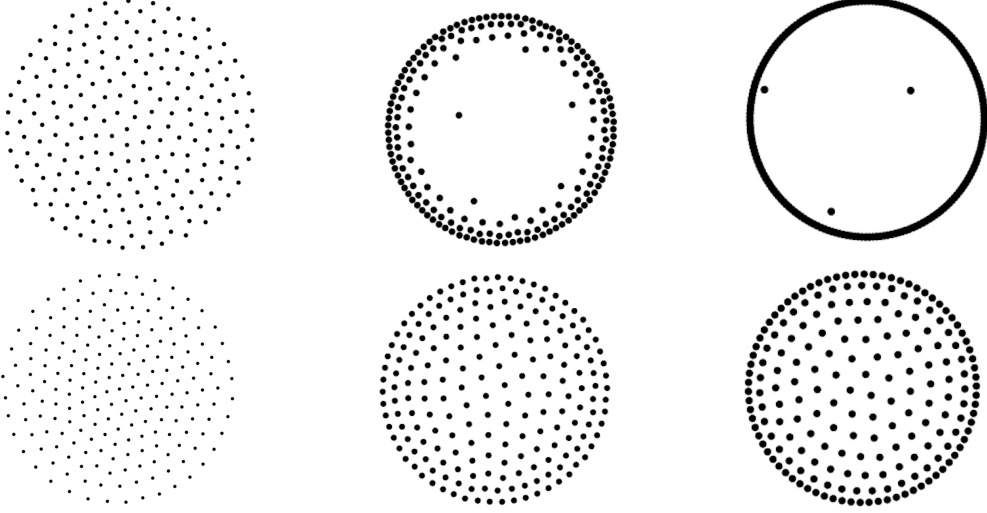


FIGURE 2. Particle simulations associated with minimizers of (1) in dimension $n = 2$. Each particle $i = 1, \dots, N$ is tracked via the system of ODEs

$$\frac{dX_i}{dt} = -\frac{1}{N} \sum_{j=1}^N \nabla K_{\alpha,\lambda}(X_i - X_j)$$

until the configuration stabilizes. The interaction kernel is given by Eq. (2), where the exponent of attraction ranges through $\alpha = 2, 20, 200$ (from left to right). *Top row:* Repulsive term replaced with the logarithmic term $-\log|x - y|$ that corresponds to the Newton potential ($\lambda = n - 2$) in two dimensions. *Bottom row:* Exponent of repulsion $\lambda = 1$, which lies in the super-Newtonian regime.

which assumes its minimum when $|x - y| = 1$ (cf. [4]). This normalization can be achieved by acting on \mathcal{P} with a suitable dilation. For the normalized kernel, the conclusions of Theorem 1 hold with $\frac{1}{\lambda}\mathcal{E}_{\infty,\lambda}$ as the limiting functional, and Theorem 2 applies without change.

We then consider the nature of minimizers for the limiting problem $\mathcal{E}_{\infty,\lambda}$. This turns out to be a rather subtle question; indeed, due to the diameter constraint, the functional $\mathcal{E}_{\infty,\lambda}$ is non-convex on \mathcal{P} . Our approach is to rephrase the limiting problem as an isodiametric capacity problem. More precisely, for $\lambda \in (0, n)$ and a set $A \subset \mathbb{R}^n$, we define the λ -capacity of A to be

$$C_\lambda(A) := \left(\inf_{\nu \in \mathcal{P}} \left\{ I_\lambda(\nu) \mid \text{supp } \nu \subset A \right\} \right)^{-1}, \quad (5)$$

where

$$I_\lambda(\nu) := \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} |x - y|^{-\lambda} d\nu(x) d\nu(y).$$

In the special case where $n = 3$ and $\lambda = 1$, C_λ agrees (up to a multiplicative constant) with the electrostatic capacity of A . It is straightforward (cf. Lemma 7) to show that

$$\inf_{\nu \in \mathcal{P}} \mathcal{E}_{\infty, \lambda}(\nu) = \left(\sup_{A \subset \mathbb{R}^n} \{C_\lambda(A) \mid \text{diam}(A) \leq 1\} \right)^{-1},$$

with a direct relationship between the optimal measure ν on the left and optimal set A on the right. This allows us to exploit tools from potential theory (cf. [10]) to partially characterize the support of minimizers.

Theorem 3 (Properties of minimizers of the limit problem). *Let $n \geq 3$, $\lambda \in (0, n)$, and assume that μ minimizes $\mathcal{E}_{\infty, \lambda}$ on \mathcal{P} . Then there exists a convex body W of constant width 1 such that*

$$\begin{cases} \text{supp } \mu \subset \partial W, & \lambda \in (0, n-2) \text{ (sub-Newtonian)}, \\ \text{supp } \mu = \partial W, & \lambda = n-2 \text{ (Newtonian)}, \\ \text{supp } \mu = W, & \lambda \in (n-2, n) \text{ (super-Newtonian)}. \end{cases}$$

The set W may depend on μ as well as λ . We do not know whether minimizers are unique up to translation, and whether $\mathcal{E}_{\infty, \lambda}$ admits additional critical points, including local minima. The proof of Theorem 3 is presented in Section 4.

The theorem extends to lower dimensions as follows. For $n = 1$, the entire range $\lambda \in (0, 1)$ is super-Newtonian, and the support of any minimizing measure is an interval of length one. In dimension $n = 2$, the entire range $\lambda \in (0, 2)$ is super-Newtonian as well, and the support of any minimizing measure is a planar convex set W of constant width 1. The role of the Newton potential $|x - y|^{2-n}$ is played by the logarithmic kernel $-\log|x - y|$; in this case, the support of a minimizer is the boundary of a planar convex set of constant width 1.

In Section 5, we prove the following result pertaining to the asymmetry of minimizers in high space dimensions. Precisely, we prove

Theorem 4 (Asymmetry of minimizers in high dimensions). *For every $\lambda > 0$ there exists N such that for all $n \geq N$,*

$$\sup \{C_\lambda(A) \mid A \subset \mathbb{R}^n, \text{diam}(A) \leq 1\} > C_\lambda(B_{1/2}^{(n)}),$$

where $B_{1/2}^{(n)} = \{x \in \mathbb{R}^n : |x| \leq \frac{1}{2}\}$.

This demonstrates that for any fixed value of $\lambda > 0$, the ball ceases to be optimal when n is sufficiently large. Thus, in this regime optimal measures are supported on the boundary of sets that are not radially symmetric. As a result, minimizers of $\mathcal{E}_{\alpha, \lambda}$ in high dimensions must also be asymmetric when α is large.

The prospect of symmetry-breaking presents an interesting, largely open, question. Even in low space dimensions, we suspect that when $0 < \lambda \ll n - 2$ the maximal capacity among bodies of given diameter may be achieved by non-symmetric sets, and that the equilibrium measure may be supported on a proper subset of the boundary. For example, in Fig. 3 we present the results of 3D particle simulations for $\lambda = 0.01$ and respectively, $\alpha = 2, 20, 200$. The simulations suggest that minimizers are asymmetric

for large α . However, the number of particles is too small to draw conclusions about the supports of minimizing measures.

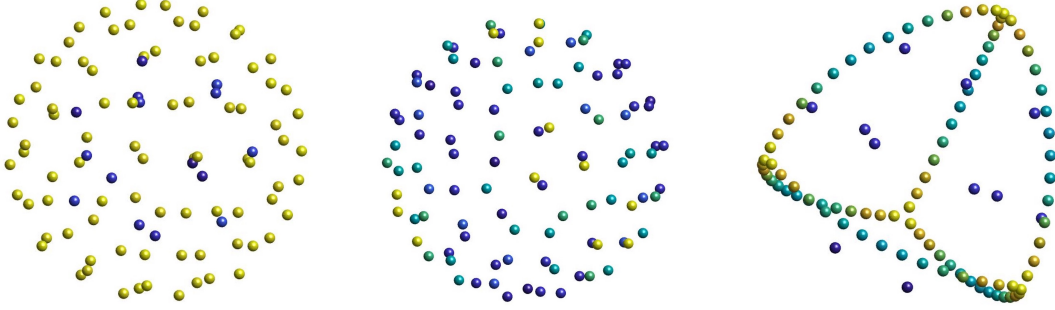


FIGURE 3. Results of particle simulations in dimension $n = 3$. The exponent of attraction ranges through $\alpha = 2, 20, 200$ (from left to right). The exponent of repulsion is $\lambda = 0.01$, which lies in the sub-Newtonian regime.

2. RELATED WORK AND FURTHER QUESTIONS

2.1. Comparisons with Related Results. According to Theorem 3, every minimizer μ of the functional $\mathcal{E}_{\infty, \lambda}$ is supported on a convex body W_μ of constant width 1, and the following relations, summarized in the table below, hold true.

TABLE 1. Characteristics of minimizers of $\mathcal{E}_{\infty, \lambda}$ in terms of a body W_μ of constant width.

| Repulsion | Geometry |
|-------------------|---|
| $\lambda < n - 2$ | $\text{supp } \mu \subset \partial W_\mu$ |
| $\lambda = n - 2$ | $\text{supp } \mu = \partial W_\mu$ |
| $\lambda > n - 2$ | $\text{supp } \mu = W_\mu$ |

In particular, the Hausdorff dimension of μ satisfies

$$\dim(\text{supp } \mu) \begin{cases} \leq n - 1, & \lambda \in (0, n - 2), \\ = n - 1, & \lambda = n - 2, \\ = n, & \lambda \in (n - 2, n). \end{cases}$$

To offer some perspective, note that classical results of geometric measure theory imply that $\dim(\text{supp } \mu) \geq \lambda$ for every Borel measure μ with $\mathcal{E}_{\infty, \lambda}(\mu) < \infty$ (see for example Theorem 4.13 in [8]). For minimizers of energy functionals defined by attractive-repulsive pair interaction kernels, a stronger lower bound was obtained in [1, Theorem 1]. Specifically, minimizers of $\mathcal{E}_{\alpha, \lambda}$ in the sub-Newtonian regime $\lambda \in (0, n - 2)$ satisfy

$$\dim(\text{supp } \mu) \geq \lambda + 2. \quad (6)$$

When $\lambda \in (n-3, n-2)$ this lower bound exceeds $n-1$, and in particular exceeds the dimension of the support of the corresponding minimizer of $\mathcal{E}_{\infty, \lambda}$. The results of [1] apply more generally to *local minimizers*, in an optimal transport topology, for a larger class of attractive-repulsive functionals with integrable singularities at the origin. In light of (6), which holds for all $\alpha > 0$, the dimensional reduction of the support for $\lambda \leq n-2$ (cf. Theorem 3 or Table 1) is only achieved in the limit. In this limit, the dimension of minimizers are strictly smaller than those in the finite regime, a consequence of the strength of the diameter constraint.

Through Theorem 1 and Lemma 7 (below), the question of what the minimizers of the limiting functional look like is transformed into an isodiametric capacity problem: *For a given $\lambda \in (0, n)$, which sets of diameter 1 have the largest λ -capacity?* Although for any given set $W \subset \mathbb{R}^n$ the equilibrium measure that realizes the capacity is unique, there could be more than one capacity-maximizing set.

One candidate for a set that maximizes capacity among sets of diameter 1 is the ball of radius $\frac{1}{2}$, which uniquely maximizes volume under the diameter restriction. For each $\lambda \in (n-2, n)$, the equilibrium measure on the ball is a well-known positive, radially symmetric density, and for $\lambda \leq n-2$ it is the uniform measure on the boundary sphere [10, p. 163]. Note, however, that the ball *minimizes* capacity among sets of given volume, indicating competition between size and shape in the isodiametric problem.

There are a number of related results for the weak repulsion regime (corresponding to $\lambda < 0$) which imply that the support of minimizers has dimension zero [1, Theorem 2] provided that the pair interaction kernel vanishes of higher order as $|x-y| \rightarrow 0$. In particular, the variance is maximized, among probability measures on \mathbb{R}^n whose support has diameter one, by the uniform measure on the vertices of the unit simplex [12].

2.2. Restricting Problem (1) to Densities and Sets. In an interesting variant of Problem (1), the minimization is restricted to absolutely continuous probability measures $\mu = \rho dx$ with density bounded by $\rho \leq m^{-1}$ for some $m > 0$.

$$\begin{cases} \text{Minimize} & \mathcal{E}'_{\alpha, \lambda}(\rho) := \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} K_{\alpha, \lambda}(x-y) \rho(x) \rho(y) dx dy \\ \text{over } \mathcal{A}_m := & \left\{ \rho \in L^1(\mathbb{R}^n) \mid 0 \leq \rho \leq m^{-1}, \int_{\mathbb{R}^n} \rho dx = 1 \right\}. \end{cases} \quad (7)$$

The density constraint plays the role of an additional repulsive term in the energy. This is relevant for biological aggregation problems, where the density of individuals cannot exceed a certain critical value. By rescaling, Problem (7) is equivalent to minimizing $\mathcal{E}'_{\alpha, \lambda}(\rho)$ among measures of mass m , subject to the density constraint $\rho \leq 1$. Unlike Problem (1), the mass m does not scale out of the problem. It is known that for each $\alpha > 0$ and $\lambda \in (0, n)$, the functional $\mathcal{E}'_{\alpha, \lambda}$ has a minimizer on \mathcal{A}_m for any $m > 0$ (cf. [7]).

Since the set of probability measures of density at most m^{-1} is a closed convex subspace of \mathcal{P} , Theorems 1 and 2 continue to hold.

Corollary 5 (Strong attraction limit with density constraint). *For $\lambda \in (0, n)$ and $\mu \in \mathcal{P}$, let $\mathcal{E}'_{\alpha, \lambda}$ be as in Problem (7), and define $\mathcal{E}'_{\infty, \lambda}(\rho) := \mathcal{E}_{\infty, \lambda}(\rho dx)$ for $\rho \in L^1$. Then*

- (1) $\mathcal{E}'_{\alpha, \lambda} \xrightarrow{\Gamma} \mathcal{E}'_{\infty, \lambda}$ as $\alpha \rightarrow \infty$ in the weak topology on L^1 .
- (2) For each $m \leq |B_{\frac{1}{2}}|$, the functional $\mathcal{E}'_{\infty, \lambda}$ attains a global minimum on \mathcal{A}_m .

The assumption on m guarantees that the energy of the uniform measure on $B_{\frac{1}{2}}$ remains bounded as $\alpha \rightarrow \infty$ (see the proof of Theorem 2). As $m \rightarrow 0$, the measures corresponding to a sequence of minimizers converge (up to translations, along suitable subsequences, weakly in \mathcal{P}) to minimizers of $\mathcal{E}_{\infty, \lambda}$.

Problem 7 is of interest also when m is large. Under certain assumptions on λ and α , $\mathcal{E}'_{\alpha, \lambda}$ is minimized for m sufficiently large by the uniform probability density on a set S of volume m ([4, 9, 13]). In the context of aggregation models, this indicates the formation of a swarm. A minimizing set is the solution of the purely geometric, non-local shape optimization problem

$$\begin{cases} \text{Minimize} & \mathcal{E}''_{\alpha, \lambda}(S) := \mathcal{E}_{\alpha, \lambda}(\nu_S) \\ \text{over } \mathcal{S}_m := \{S \subset \mathbb{R}^n \mid |S| = m\}, \end{cases} \quad (8)$$

where ν_S is the uniform probability measure on S . It turns out that the infimum in Problem (8) agrees with Problem (7), but it is not always attained. If the density of a minimizer of $\mathcal{E}'_{\alpha, \lambda}$ on \mathcal{A}_m falls strictly between 0 and m^{-1} on all or part of its support, then the shape optimization problem (8) has no solution [4, Theorem 4.4], indicating a failure to fully aggregate. In this case, minimizing sequences for Problem (8) diverge due to oscillations. When m is too small, typically $\rho < m^{-1}$ everywhere (cf. [4, 9, 13]), preventing even partial aggregation.

All known solutions of the shape optimization problem (8) are radially symmetric, and in many cases they are large balls (cf. [4, 9, 13]). It may be possible to discover interesting examples of symmetry-breaking in the strong-attraction limit, using Corollary 5 and the known relation between Problems (7) and (8).

We are not aware of any explicit characterization of the minimizers for $\mathcal{E}'_{\infty, \lambda}$ on \mathcal{A}_m , even in the Newtonian case. Suppose that W maximizes capacity among sets of given diameter. Since the density constraint prevents minimizers to concentrate on a lower-dimensional set, one may wonder whether a thin neighborhood of ∂W might appear as a solution to Problem (8), and whether such a solution persists for sufficiently large finite values of α ? When W is not a ball, this could give rise to symmetry-breaking in Problems (7) and (8).

3. CONVERGENCE

We begin by recalling a few definitions. Given a topological space X , let $(G_n)_n$ be a sequence of functions on X . We say that (G_n) **Gamma-converges** to a function G ($G_n \xrightarrow{\Gamma} G$) if the following two conditions hold for every $x \in X$:

- *Lower bound inequality:* for all sequences $(x_n)_n \subset X$ such that $x_n \rightarrow x \in X$,

$$\liminf_{n \rightarrow \infty} G_n(x_n) \geq G(x);$$

- *Upper bound inequality:* for all $x \in X$ there exists a sequence $(x_n)_n \subset X$ such that $x_n \rightarrow x$ and

$$\limsup_{n \rightarrow \infty} G_n(x_n) \leq G(x).$$

Gamma-convergence has many useful implications, the most important of which is that if x_n minimizes G_n over X , then every cluster point of the sequence (x_n) minimizes G over X (cf. [3]).

Given a sequence of measures $(\mu_n)_n \subset \mathcal{P}$, we say $(\mu_n)_n$ **converge weakly** to $\mu \in \mathcal{P}$ ($\mu_n \rightharpoonup \mu$) if

$$\lim_{n \rightarrow \infty} \int \phi d\mu_n = \int \phi d\mu$$

for every bounded continuous function ϕ on \mathbb{R}^n . This induces the weak topology on \mathcal{P} .

Proof of Theorem 1. Let $\mu \in \mathcal{P}$ be given. In the case where $\text{diam}(\text{supp } \mu) > 1$, choose two points $p, q \in \text{supp } \mu$ with $|p - q| > 1$. By continuity of the distance function, there exist open neighborhoods U, V of p and q such that $\text{dist}(U, V) > 1$. For any sequence of measures (μ_n) with $\mu_n \rightharpoonup \mu$ in \mathcal{P} , we have

$$\begin{aligned} \mathcal{E}_{\alpha, \lambda}(\mu_n) &= \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} |x - y|^\alpha + |x - y|^{-\lambda} d\mu_n(x) d\mu_n(y) \\ &\geq (\text{dist}(U, V))^\alpha \mu_n(U) \mu_n(V). \end{aligned}$$

Since $\liminf \mu_n(U) \geq \mu(U) > 0$ and likewise for V , it follows that $\mathcal{E}_{\alpha_n, \lambda}(\mu_n) \rightarrow \infty$ along every sequence (α_n) with $\alpha_n \rightarrow \infty$, verifying simultaneously the lower and upper bound inequalities for this case.

Otherwise, $\text{diam}(\text{supp } \mu) \leq 1$. To see the lower bound inequality, let (μ_n) be a sequence in \mathcal{P} that converges weakly to μ , and let $t > 0$. For every $\alpha > 0$,

$$\mathcal{E}_{\alpha, \lambda}(\mu_n) \geq \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} \min\{|x - y|^{-\lambda}, t\} d\mu_n(x) d\mu_n(y).$$

Since $\mathbb{R}^n \times \mathbb{R}^n$ is separable, the product measures $\mu_n \times \mu_n$ converge weakly to $\mu \times \mu$, and thus for any sequence (α_n) ,

$$\liminf_{n \rightarrow \infty} \mathcal{E}_{\alpha_n, \lambda}(\mu_n) \geq \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} \min\{|x - y|^{-\lambda}, t\} d\mu(x) d\mu(y).$$

By monotone convergence, taking $t \rightarrow \infty$ yields the lower bound inequality.

The upper bound inequality is achieved by a sequence of properly chosen dilations of μ . Given a sequence $\alpha_n \rightarrow \infty$, set $\beta_n = e^{\frac{1}{\sqrt{\alpha_n}}}$, and define a sequence of Borel measures by

$$\mu_n(A) = \mu(\beta_n A), \quad n \geq 1.$$

Since $\beta_n \rightarrow 1$, clearly $\mu_n \rightarrow \mu$. We estimate

$$\begin{aligned} \mathcal{E}_{\alpha_n, \lambda}(\mu_n) &= \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} |x - y|^{\alpha_n} + |x - y|^{-\lambda} d\mu_n(x) d\mu_n(y) \\ &= \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} \left(\frac{|x - y|}{\beta_n} \right)^{\alpha_n} + \left(\frac{|x - y|}{\beta_n} \right)^{-\lambda} d\mu(x) d\mu(y) \\ &\leq e^{-\sqrt{\alpha_n}} + e^{\frac{\lambda}{\sqrt{\alpha_n}}} \mathcal{E}_{\infty, \lambda}(\mu). \end{aligned}$$

We have used that $|x - y| \leq 1$ on the support of μ to bound the first summand of the integrand, and inserted the definition of the limiting functional into the second summand. The desired inequality follows upon taking $n \rightarrow \infty$. \square

The proof of Theorem 2 requires a compactness argument. To this end one often resorts to an application of Lions' concentration compactness principle for probability measures (cf. [16, Section 4.3]) which asserts that every sequence $(\mu_n)_n$ in \mathcal{P} has a subsequence $(\mu_{n_k})_k$ satisfying one of the three following alternatives: (i) tightness up to translation (ii) vanishing (mass sent to infinity) or (iii) dichotomy (splitting). A standard technique is to show that (ii) and (iii) can not happen, yielding (i) which, precisely, means: There exists a sequence $(y_k)_k \subset \mathbb{R}^n$ such that for all $\varepsilon > 0$ there exists $R > 0$ with the property that $\mu_{n_k}(B_R(y_k)) \geq 1 - \varepsilon$ for all k .

However, in our simpler case we may just as well directly prove tightness and obtain compactness.

Lemma 6. *Let $\mathcal{E}_{\alpha, \lambda}$ be as in Eq. (1), let (α_n) be a sequence with $\alpha_n \rightarrow \infty$, and fix $\lambda \in (0, n)$. Then every sequence (μ_n) in \mathcal{P} such that $\mathcal{E}_{\alpha_n, \lambda}(\mu_n)$ is bounded has a subsequence that converges weakly, up to translations, to some $\mu \in \mathcal{P}$.*

Proof. Let (μ_n) be such that

$$\sup_{n \in \mathbb{N}} \mathcal{E}_{\alpha_n, \lambda}(\mu_n) < \infty.$$

Fix an $R > 1$. We have the lower bounds

$$\begin{aligned} \mathcal{E}_{\alpha_n, \lambda}(\mu_n) &\geq \iint_{|x - y| \geq R} R^{\alpha_n} d\mu_n(x) d\mu_n(y) \\ &\geq R^{\alpha_n} \int_{\mathbb{R}^n} \mu_n(\mathbb{R}^n \setminus B_R(y)) d\mu_n(y) \\ &\geq R^{\alpha_n} \left(1 - \sup_{y \in \mathbb{R}^n} \mu_n(B_R(y)) \right). \end{aligned}$$

Since the left hand side is bounded by assumption while $\alpha_n \rightarrow \infty$, it follows that $\sup_{y \in \mathbb{R}^n} \mu_n(B_R(y)) \rightarrow 1$. This establishes the first alternative of Lions' concentration compactness principle.

Choose a sequence $(y_n) \subset \mathbb{R}^n$ such that

$$\lim_{n \rightarrow \infty} \mu_n(B_2(y_n)) = 1.$$

Given $\varepsilon > 0$, let N be so large that $\mu_n(B_2(y_n)) \geq 1 - \varepsilon$ for all $n > N$. Then choose R so large that $\mu_n(B_R(y_n)) \geq 1 - \varepsilon$ for $n = 1, \dots, N$. Taking $R \geq 2$ ensures that $\mu_n(B_R(y_n)) \geq \mu_n(B_2(y_n)) \geq 1 - \varepsilon$ also for $n > N$.

Let $(\tilde{\mu}_n)_n$ be the sequence of translates of μ_n defined by

$$\tilde{\mu}_n(A) = \mu_n(y_n + A), \quad n \geq 1$$

for each Borel set $A \subset \mathbb{R}^n$. Since $(\tilde{\mu}_n)$ is tight, Prokhorov's theorem yields a subsequence $(\tilde{\mu}_{n_k})_k$ that converges weakly in \mathcal{P} . \square

Proof of Theorem 2. Let (α_n) be a nonnegative sequence with $\alpha_n \rightarrow \infty$, and let (μ_n) be a sequence of measures such that each μ_n minimizes $\mathcal{E}_{\alpha_n, \lambda}$. We will prove that $(\mathcal{E}_{\alpha_n, \lambda}(\mu_{\alpha_n}))_n$ is bounded, and then apply Lemma 6.

Let ν be the uniform probability measure on the ball of radius $\frac{1}{2}$. Since μ_n minimizes $\mathcal{E}_{\alpha_n, \lambda}$ for each n , we have

$$\begin{aligned} \mathcal{E}_{\alpha_n, \lambda}(\mu_n) &\leq \mathcal{E}_{\alpha_n, \lambda}(\nu) \\ &= \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} |x - y|^\alpha + |x - y|^{-\lambda} d\nu(x) d\nu(y) \\ &\leq 1 + \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} |x - y|^{-\lambda} d\nu(x) d\nu(y) \\ &< \infty. \end{aligned}$$

In the last two inequalities, we have used that the support of ν has diameter one, and that the kernel is locally integrable.

By Lemma 6 there exists a subsequence μ_{n_k} that converges weakly up to translation, to some measure $\mu \in \mathcal{P}$. Since the functionals are translation invariant, we may assume that the sequence of minimizers itself that has a subsequence converging weakly to μ . By the properties of the Gamma-limit, μ is a global minimizer of $\mathcal{E}_{\infty, \lambda}$. \square

4. CHARACTERIZATION OF MINIMIZERS

We recall some classical results from potential theory. First, recall the λ -capacity of a set $A \subset \mathbb{R}^n$ previously defined in (5) as the reciprocal of the minimum of the repulsive energy I_λ over measures supported in A . If A is a compact set of positive Lebesgue measure, the λ -capacity is finite by the local integrability of the Riesz-potential, and the supremum is achieved by some measure $\mu \in \mathcal{P}$. Since I_λ is positive definite, the minimizer is unique.

The next lemma relates the minimization problem for $\mathcal{E}_{\infty, \lambda}$ to an isodiametric capacity problem.

Lemma 7. *Let $n \geq 1$, $\lambda \in (0, n)$. Then*

$$\inf_{\nu \in \mathcal{P}} \mathcal{E}_{\infty, \lambda}(\nu) = \left(\sup_{A \subset \mathbb{R}^n} \left\{ C_\lambda(A) \mid \text{diam}(A) \leq 1 \right\} \right)^{-1}.$$

Furthermore, the infimum on the left hand side is attained for some measure μ with $\text{diam}(\text{supp } \mu) = 1$, and the supremum on the right hand side is attained for some convex body $W \subset \mathbb{R}^n$ of constant width 1 containing the support of μ . Conversely, if W maximizes λ -capacity among bodies of constant width, then the equilibrium measure on W attains the minimum on the left hand side.

Proof. We split the minimization problem for $\mathcal{E}_{\infty,\lambda}$ into two steps,

$$\begin{aligned} \inf_{\nu \in \mathcal{P}} \mathcal{E}_{\infty,\lambda}(\nu) &= \inf_{A \subset \mathbb{R}^n} \left\{ \inf_{\nu \in \mathcal{P}} \{I_\lambda(\nu) \mid \text{supp } \nu \subset A\} \mid \text{diam}(A) \leq 1 \right\} \\ &= \left(\sup_{A \subset \mathbb{R}^n} \{C_\lambda(A) \mid \text{diam}(A) = 1\} \right)^{-1}. \end{aligned}$$

By Theorem 2, the infimum on the left hand side is attained for some measure $\mu \in \mathcal{P}$. Clearly, $\text{diam}(\text{supp } \mu) = 1$, since otherwise μ could be rescaled to lower the value of $\mathcal{E}_{\infty,\lambda}$. Moreover, $A = \text{supp } \mu$ achieves the supremum on the right hand side, and μ is the equilibrium measure for the capacity $C_\lambda(A)$. Since the capacity increases monotonically under inclusion, we may replace A by its convex hull. The last claim follows since every closed convex set of diameter 1 is contained in a convex body W of constant width 1 (cf. [14]). Since $C_\lambda(W) = C_\lambda(\text{supp } \mu)$, it follows that μ is the equilibrium measure also for W . \square

We can now appeal to known properties of equilibrium measures in classical potential theory. Given a probability measure μ on \mathbb{R}^n and $\lambda \in (0, n)$, we define the corresponding potential by

$$\phi_\lambda^\mu(x) := \int_{\mathbb{R}^n} |x - y|^{-\lambda} d\mu(y).$$

For any $x \in \mathbb{R}^n$, the integral is well-defined and strictly positive, though possibly infinite. The function has the following regularity property outside the support of μ .

Lemma 8. *Let μ be a probability measure on \mathbb{R}^n . On $\mathbb{R}^n \setminus \text{supp } \mu$, the potential ϕ_μ^λ is smooth and*

$$\begin{cases} \text{strictly subharmonic} & \lambda \in (0, n-2), \\ \text{harmonic} & \lambda = n-2, \\ \text{strictly superharmonic} & \lambda \in (n-2, n). \end{cases}$$

Proof. By direct computation,

$$\Delta \phi_\lambda^\mu(x) = \lambda(\lambda + 2 - n) \int_{\mathbb{R}^n} |x - y|^{-\lambda-2} d\mu(y)$$

away from the support of μ . \square

In the super-Newtonian regime, the equilibrium measure has the following property.

Lemma 9. [10, p.137] *Let $\lambda \geq n-2$, and let $W \subset \mathbb{R}^n$ be a compact set of positive capacity. If $\mu \in \mathcal{P}$ minimize I_λ among probability measures supported on W , then*

$$\begin{aligned} \phi_\lambda^\mu(x) &= I_\lambda(\mu) \quad \text{approximately everywhere on } W \\ \phi_\lambda^\mu(x) &\leq I_\lambda(\mu) \quad \text{throughout } \mathbb{R}^n \end{aligned}$$

where approximately everywhere means everywhere except on a set of capacity zero.

We are ready for the proof of Theorem 3.

Proof. Let μ be a minimizer of $\mathcal{E}_{\infty,\lambda}$. By Lemma 7, μ is the equilibrium measure that achieves the λ -capacity of some convex body W of constant width 1. When $\lambda \leq n-2$, classical results of potential theory (cf. [10, p.162]) ensure that $\text{supp } \mu \subset \partial W$. This proves the claim in the sub-Newtonian regime.

Let now $\lambda \geq n-2$, and $p \in \partial W$. Since W is a convex body, every neighborhood of p intersects the interior of W in a set of positive volume (and hence positive capacity). Again by classical results of potential theory (cf. [10, p.164]), p lies in the support of μ . Therefore $\partial W \subset \text{supp } \mu$. Together with the result for $\lambda \leq n-2$, this completes the proof in the Newtonian case.

For $\lambda > n-2$ Lemma 8 yields that the potential ϕ_μ^λ is strictly subharmonic outside the support of μ . By the strong maximum principle, ϕ_μ^λ is non-constant on every non-empty open set U with $\mu(U) = 0$. On the other hand, ϕ_μ^λ is constant on the interior of W by Lemma 9. Therefore $\mu(U) > 0$ for every non-empty open subset of the interior of W , and we conclude that $W \subset \text{supp } \mu$. This proves the claim in the super-Newtonian regime. \square

5. CAPACITY ESTIMATES

We close with some simple capacity estimates which will prove Theorem 4.

Lemma 10. *Let $n \geq 1$, $\lambda \in (0, n)$. Then*

$$\sup_{A \subset \mathbb{R}^n} \left\{ C_\lambda(A) \mid \text{diam}(A) = 1 \right\} < 1.$$

Proof. By Lemma 7, there is a set $A \subset \mathbb{R}^n$ that maximizes the capacity C_λ among sets of diameter 1. Let μ be the equilibrium measure on A that achieves the capacity. We estimate

$$\mathcal{E}_{\infty,\lambda}(\mu) - 1 \geq \int \left(\int_{B_{\frac{1}{2}}(x)} (|x-y|^{-\lambda} - 1) d\mu(y) \right) d\mu(x) > 0,$$

where the first inequality holds since the integrand is nonnegative for every pair of points $x, y \in A$, and the second inequality uses that $\mu(B_{\frac{1}{2}}(x)) > 0$ for x in the support of μ . By Lemma 7, $C_\lambda(A) = (\mathcal{E}_{\infty,\lambda}(\mu))^{-1} < 1$, as claimed. \square

We next consider the capacity of balls in high dimensions.

Lemma 11. *For every $\lambda > 0$*

$$\lim_{n \rightarrow \infty} C_\lambda(B_{1/2}^{(n)}) = 2^{-\frac{\lambda}{2}}.$$

Proof. This follows by direct computation of $C_\lambda(B_{1/2}^{(n)})$ (cf. [10, p.163]) and Stirling's approximation. \square

Finally, we construct sets of larger capacity in high dimensions.

Lemma 12. *For every $\lambda > 0$,*

$$\lim_{n \rightarrow \infty} \left(\sup \left\{ C_\lambda(A) \mid A \subset \mathbb{R}^n, \text{diam}(A) \leq 1 \right\} \right) = 1.$$

Proof. Since $C_\lambda(A) < 1$ for all n by Lemma 10, it suffices to establish the corresponding lower bound on the capacity.

We will construct a family of subsets $(A_n)_n$ of diameter 1 in the sphere of radius $\frac{1}{2}\sqrt{2}$ in \mathbb{R}^n that achieves this limit. For each $n > \lambda + 1$, the spherical cap of diameter 1 in this sphere has positive λ -capacity. Let A_n be a set of maximal capacity among such subsets, and let μ_n be the equilibrium measure on A_n that attains the capacity.

For $m, n > \lambda + 1$, consider a convex combination

$$\mu = (1 - t)(\mu_m \otimes \delta) + t(\delta \otimes \mu_n)$$

on \mathbb{R}^{m+n} , where δ denotes the unit mass at 0 in \mathbb{R}^n and \mathbb{R}^m , respectively, and $t \in (0, 1)$ will be chosen below. By definition, μ is supported on $(A_m \times \{0\}) \cup (\{0\} \times A_n)$, which lies in the sphere of radius $\frac{1}{2}\sqrt{2}$ in \mathbb{R}^{m+n} and has diameter 1. We estimate

$$\begin{aligned} \mathcal{E}_{\infty, \lambda}(\mu_{m+n}) - 1 &\leq \mathcal{E}_{\infty, \lambda}(\mu) - 1 \\ &= \int \int (|x - y|^{-\lambda} - 1) d\mu(x) d\mu(y) \\ &= (1 - t)^2(\mathcal{E}_{\infty, \lambda}(\mu_m) - 1) + t^2(\mathcal{E}_{\infty, \lambda}(\mu_n) - 1); \end{aligned}$$

the mixed terms vanish because $|x - y| = 1$ whenever $x \in A_m \times \{0\}$ and $y \in \{0\} \times A_n$. Minimization over t yields

$$\mathcal{E}_{\infty, \lambda}(\mu_{m+n}) - 1 \leq \frac{(\mathcal{E}_{\infty, \lambda}(\mu_m) - 1)(\mathcal{E}_{\infty, \lambda}(\mu_n) - 1)}{\mathcal{E}_{\infty, \lambda}(\mu_m) + \mathcal{E}_{\infty, \lambda}(\mu_n) - 2}.$$

Since $\mathcal{E}_{\infty, \lambda}(\mu_n) > 1$ for all n by Lemma 10, we can pass to reciprocals and conclude that $(\mathcal{E}_{\infty, \lambda}(\mu_n) - 1)^{-1}$ is superadditive in n . By Fekete's superadditivity lemma

$$\lim_{n \rightarrow \infty} \frac{1}{n} (\mathcal{E}_{\infty, \lambda}(\mu_n) - 1)^{-1} = \sup_n \frac{1}{n} (\mathcal{E}_{\infty, \lambda}(\mu_n) - 1)^{-1} > 0.$$

It follows that $\lim_{n \rightarrow \infty} C_\lambda(A_n) = \left(\lim_{n \rightarrow \infty} \mathcal{E}_{\infty, \lambda}(\mu_n) \right)^{-1} = 1$. \square

The proof of Theorem 4 is an immediate corollary of Lemma 11 and Lemma 12 since $2^{-\frac{\lambda}{2}} < 1$ for every $\lambda > 0$. Note that the near-maximizers constructed in the proof of Lemma 12 have dimension much below n , but this need not be true for actual maximizers.

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DEPT. OF MATHEMATICS, UNIVERSITY OF TORONTO, TORONTO, ON, CANADA
E-mail address: almut@math.toronto.edu

DEPT. OF MATHEMATICS AND STATISTICS, MCGILL UNIVERSITY, MONTRÉAL, QC, CANADA
E-mail address: rustum.choksi@mcgill.ca

DEPT. OF MATHEMATICS AND STATISTICS, MCGILL UNIVERSITY, MONTRÉAL, QC, CANADA
 CURRENT ADDRESS: COURANT INSTITUTE OF MATHEMATICAL SCIENCES NEW YORK, NY
E-mail address: elias.hess-childs@courant.nyu.edu