

EnKF and Catastrophic filter divergence

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The filtering problem

We have a **deterministic model**

$$\frac{d\mathbf{v}}{dt} = F(\mathbf{v}) \quad \text{with } \mathbf{v}_0 \sim N(m_0, C_0).$$

We will denote $\mathbf{v}(t) = \Psi_t(\mathbf{v}_0)$. Think of this as **very high dimensional** and **nonlinear**.

We want to **estimate** $\mathbf{v}_j = \mathbf{v}(jh)$ for some $h > 0$ and $j = 0, 1, \dots, J$ given the **observations**

$$\mathbf{y}_j = H\mathbf{v}_j + \xi_j \quad \text{for } \xi_j \text{ iid } N(0, \Gamma).$$

We can write down the conditional density using **Bayes' formula** ...

But for high dimensional nonlinear systems it's **horrible**.

Alternatively, we can use **EnKF** to draw **approximate samples** from the posterior.

The set-up for EnKF

Suppose we are given the ensemble $\{u_j^{(1)}, \dots, u_j^{(K)}\}$ at time j . For each particle, we create an **artificial observation**

$$y_{j+1}^{(k)} = y_{j+1} + \xi_{j+1}^{(k)} \quad , \quad \xi_{j+1}^{(k)} \text{ iid } N(0, \Gamma).$$

We update each particle using the **Kalman update**

$$u_{j+1}^{(k)} = \Psi_h(u_j^{(k)}) + G(u_j) \left(y_{j+1}^{(k)} - H\Psi_h(u_j^{(k)}) \right) \quad ,$$

where $G(u_j)$ is the **Kalman gain** computed using the **forecasted ensemble covariance**

$$\hat{C}_{j+1} = \frac{1}{K} \sum_{k=1}^K (\Psi_h(u_j^{(k)}) - \overline{\Psi_h(u_j)})^T (\Psi_h(u_j^{(k)}) - \overline{\Psi_h(u_j)}) \quad .$$

There aren't many **theorems** about
EnKF.

Ideally, we would like a theorem about
long time behaviour of the filter for a
finite ensemble size.

Filter divergence

It has been observed (★) that when observations are **very frequent** the ensemble can **blow-up** (ie. reach machine-infinity) in **finite time**, even when the model has nice bounded solutions.

This is known as **catastrophic filter divergence**.

It is suggested in (★) that this is caused by numerically integrating a stiff-system. Our aim is to “prove” this.

★ Harlim, Majda (2010), Gottwald (2011), Gottwald, Majda (2013).

Assumptions on the dynamics

We make a **dissipativity** assumption on F . Namely that

$$F(\cdot) = A\cdot + B(\cdot, \cdot)$$

with A linear elliptic and B bilinear, satisfying certain estimates and symmetries.

Eg. 2d-Navier-Stokes, Lorenz-63, Lorenz-96.

Discrete time results

For a fixed observation frequency $h > 0$ we can prove

Theorem (AS,DK)

If $H = \Gamma = Id$ then there exists constant $\beta > 0$ such that

$$\mathbf{E}|u_j^{(k)}|^2 \leq e^{2\beta jh} \mathbf{E}|u_0^{(k)}|^2 + 2K\gamma^2 \left(\frac{e^{2\beta jh} - 1}{e^{2\beta h} - 1} \right)$$

Rmk. This becomes useless as $h \rightarrow 0$

For observations with $h \ll 1$, we need another approach.

The EnKF equations look like a discretization

Recall the ensemble update equation

$$\begin{aligned} \mathbf{u}_{j+1}^{(k)} &= \Psi_h(\mathbf{u}_j^{(k)}) + G(\mathbf{u}_j) \left(\mathbf{y}_{j+1}^{(k)} - H\Psi_h(\mathbf{u}_j^{(k)}) \right) \\ &= \Psi_h(\mathbf{u}_j^{(k)}) + \hat{\mathbf{C}}_{j+1} H^T (H^T \hat{\mathbf{C}}_{j+1} H + \Gamma)^{-1} \left(\mathbf{y}_{j+1}^{(k)} - H\Psi_h(\mathbf{u}_j^{(k)}) \right) \end{aligned}$$

Subtract $\mathbf{u}_j^{(k)}$ from both sides and divide by h

$$\begin{aligned} \frac{\mathbf{u}_{j+1}^{(k)} - \mathbf{u}_j^{(k)}}{h} &= \frac{\Psi_h(\mathbf{u}_j^{(k)}) - \mathbf{u}_j^{(k)}}{h} \\ &\quad + \hat{\mathbf{C}}_{j+1} H^T (hH^T \hat{\mathbf{C}}_{j+1} H + h\Gamma)^{-1} \left(\mathbf{y}_{j+1}^{(k)} - H\Psi_h(\mathbf{u}_j^{(k)}) \right) \end{aligned}$$

Clearly we need to rescale the noise (ie. Γ).

Continuous-time limit

If we set $\Gamma = h^{-1}\Gamma_0$ and substitute $y_{j+1}^{(k)}$, we obtain

$$\frac{u_{j+1}^{(k)} - u_j^{(k)}}{h} = \frac{\Psi_h(u_j^{(k)}) - u_j^{(k)}}{h} + \widehat{C}_{j+1}H^T(hH^T\widehat{C}_{j+1}H + \Gamma_0)^{-1} \\ \left(H\mathbf{v} + h^{-1/2}\Gamma_0^{1/2}\xi_{j+1} + h^{-1/2}\Gamma_0^{1/2}\xi_{j+1}^{(k)} - H\Psi_h(u_j^{(k)}) \right)$$

But we know that

$$\Psi_h(u_j^{(k)}) = u_j^{(k)} + O(h)$$

and

$$\widehat{C}_{j+1} = \frac{1}{K} \sum_{k=1}^K (\Psi_h(u_j^{(k)}) - \overline{\Psi_h(u_j)})^T (\Psi_h(u_j^{(k)}) - \overline{\Psi_h(u_j)}) \\ = \frac{1}{K} \sum_{k=1}^K (u_j^{(k)} - \overline{u_j})^T (u_j^{(k)} - \overline{u_j}) + O(h) = C(u_j) + O(h)$$

Continuous-time limit

We end up with

$$\begin{aligned} \frac{u_{j+1}^{(k)} - u_j^{(k)}}{h} &= \frac{\Psi_h(u_j^{(k)}) - u_j^{(k)}}{h} - C(u_j)H^T\Gamma_0^{-1}H(u_j^{(k)} - v_j) \\ &\quad + C(u_j)H^T\Gamma_0^{-1} \left(h^{-1/2}\xi_{j+1} + h^{-1/2}\xi_{j+1}^{(k)} \right) + O(h) \end{aligned}$$

This looks like a **numerical scheme** for **Itô S(P)DE**

$$\begin{aligned} \frac{du^{(k)}}{dt} &= F(u^{(k)}) - C(u)H^T\Gamma_0^{-1}H(u^{(k)} - v) \quad (\bullet) \\ &\quad + C(u)H^T\Gamma_0^{-1/2} \left(\frac{dW^{(k)}}{dt} + \frac{dB}{dt} \right) . \end{aligned}$$

First observation: nudging

$$\begin{aligned} \frac{d\mathbf{u}^{(k)}}{dt} = & F(\mathbf{u}^{(k)}) - C(\mathbf{u})H^T\Gamma_0^{-1}H(\mathbf{u}^{(k)} - \mathbf{v}) \quad (\bullet) \\ & + C(\mathbf{u})H^T\Gamma_0^{-1/2} \left(\frac{d\mathbf{W}^{(k)}}{dt} + \frac{d\mathbf{B}}{dt} \right) . \end{aligned}$$

- 1 - Extra dissipation term only sees **differences in observed space**
- 2 - Extra dissipation only occurs in the **space spanned by ensemble**

Second observation: Kalman-Bucy limit

If F were **linear** and we write $m(t) = \frac{1}{K} \sum_{k=1}^K u^{(k)}(t)$ then

$$\begin{aligned} \frac{dm}{dt} = & F(m) - C(u)H^T \Gamma_0^{-1} H(m - v) \\ & + C(u)H^T \Gamma_0^{-1/2} \frac{dB}{dt} + O(K^{-1/2}). \end{aligned}$$

This is the equation for the **Kalman-Bucy** filter, with empirical covariance $C(u)$. The remainder $O(K^{-1/2})$ can be thought of as a **sampling error**.

Continuous-time results

Theorem (AS,DK)

Suppose that $\{u^{(k)}\}_{k=1}^K$ satisfy (\bullet) with $H = \Gamma = Id$. Let

$$e^{(k)} = u^{(k)} - v .$$

Then there exists constant $\beta > 0$ such that

$$\frac{1}{K} \sum_{k=1}^K \mathbf{E} |e^{(k)}(t)|^2 \leq \left(\frac{1}{K} \sum_{k=1}^K \mathbf{E} |e^{(k)}(0)|^2 \right) \exp(\beta t) .$$

Why do we need $H = \Gamma = Id$?

In the equation

$$\begin{aligned} \frac{d\mathbf{u}^{(k)}}{dt} = & F(\mathbf{u}^{(k)}) - C(\mathbf{u})H^T\Gamma_0^{-1}H(\mathbf{u}^{(k)} - \mathbf{v}) \\ & + C(\mathbf{u})H^T\Gamma_0^{-1/2} \left(\frac{dW^{(k)}}{dt} + \frac{dB}{dt} \right). \end{aligned}$$

The **energy** pumped in by the noise must be balanced by **contraction** of $(\mathbf{u}^{(k)} - \mathbf{v})$. So the operator

$$C(\mathbf{u})H\Gamma_0^{-1}H$$

must be **positive-definite**.

Both $C(\mathbf{u})$ and $H\Gamma_0^{-1}H$ are pos-def, but this doesn't guarantee the same for the **product**!

Summary + Future Work

- (1) Writing down an SDE/SPDE allows us to see the **important quantities** in the algorithm.
 - (2) Does not “prove” that catastrophic filter divergence is a numerical phenomenon, but is a decent starting point.
- (1) Improve the condition on H .
 - (2) If we can **measure** the important quantities, then we can test the performance during the algorithm.