

# Scientific Computing: Ordinary Differential Equations

**Aleksandar Donev**  
*Courant Institute, NYU<sup>1</sup>*  
*donev@courant.nyu.edu*

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# Outline

- 1 Initial Value Problems
- 2 One-step methods for ODEs
- 3 MATLAB ode suite
- 4 Stability
- 5 Conclusions

# Initial Value Problems

- We want to numerically approximate the solution to the **ordinary differential equation**

$$\frac{dx}{dt} = x'(t) = \dot{x}(t) = f(x(t), t),$$

with **initial condition**  $x(t=0) = x(0) = x_0$ .

- This means that we want to generate an approximation to the **trajectory**  $x(t)$ , for example, a sequence  $x(t_k = k\Delta t)$  for  $k = 1, 2, \dots, N = T/\Delta t$ , where  $\Delta t$  is the **time step** used to discretize time.
- If  $f$  is independent of  $t$  we call the system **autonomous**.
- Note that second-order equations can be written as a **system** of first-order equations:

$$\frac{d^2x}{dt^2} = \ddot{x}(t) = f[x(t), t] \quad \equiv \quad \begin{cases} \dot{x}(t) = v(t) \\ \dot{v}(t) = f[x(t), t] \end{cases}$$

# Relation to Numerical Integration

- If  $f$  is independent of  $x$  then the problem is equivalent to numerical integration

$$x(t) = x_0 + \int_0^t f(s) ds.$$

- More generally, we cannot compute the integral because it depends on the unknown answer  $x(t)$ :

$$x(t) = x_0 + \int_0^t f(x(s), s) ds.$$

- Numerical methods are based on approximations of  $f(x(s), s)$  into the “future” based on knowledge of  $x(t)$  in the “past” and “present”.

# Convergence

- Consider a trajectory numerically discretized as a sequence that **approximates** the exact solution at a **discrete** set of points:

$$x^{(k)} \approx x(t_k = k\Delta t), \quad k = 1, \dots, T/\Delta t.$$

- A method is said to **converge with order**  $p > 0$ , or to have **order of accuracy**  $p$ , if for any finite  $T$  for which the ODE has a solution,

$$\left| x^{(k)} - x(k\Delta t) \right| = O(\Delta t^p) \text{ for all } 0 \leq k \leq T/\Delta t.$$

- All methods are recursions that compute a new  $x^{(k+1)}$  from previous  $x^{(k)}$  by evaluating  $f(x)$  several times. For example, **one-step methods** have the form

$$x^{(k+1)} = G(x^{(k)}, \Delta t; f).$$

# Consistency

- The **local truncation error** (LTE) of a method is the amount by which the *exact* solution does not satisfy the numerical scheme at the end of the time step if started from the correct solution  $x^{(k)} = x(k\Delta t)$ :

$$e_k = x[(k+1)\Delta t] - G[x(k\Delta t), \Delta t; f],$$

- A method is **consistent with order**  $q > 1$  if  $|e_k| = O(\Delta t^q)$ .
- The **global truncation error** is the actual error

$$E_{t=k\Delta t} = \left| x^{(k)} - x(t = k\Delta t) \right|.$$

- Numerical analysis question: Can the global error be bounded by  $O(\Delta t^p)$  if the local one is  $O(\Delta t^q)$ ?

# Propagation of errors

- *Crude* estimate: If one makes an error  $O(\Delta t^q)$  at each time step, the global error after  $T/\Delta t$  time steps can become on the order of

$$\left| x^{(k)} - x(k\Delta t) \right| = O\left(\Delta t^q \cdot \frac{T}{\Delta t}\right) = O(\Delta t^{q-1}) = O(\Delta t^p),$$

and we must have  $p = q - 1 > 0$  for convergence.

- This result is often the right one, but it has a hidden assumption that *errors made at previous steps do not grow* but rather stay of the same order so they can be added.
- In practice, errors made in previous time steps will either grow or shrink with time. If they grow “too fast” we are in trouble.
- So we arrive for the first time at a recurring theme: **Convergence requires stability in addition to consistency**. We discuss stability later on, after we give some basic methods for solving ODEs.

# Euler's Method

- Assume that we have our approximation  $x^{(k)}$  and want to move by one time step:

$$x^{(k+1)} \approx x^{(k)} + \int_{k\Delta t}^{(k+1)\Delta t} f(x(s), s) ds.$$

- The simplest possible thing is to use a piecewise constant approximation:

$$f(x(s), s) \approx f(x^{(k)}) = f^{(k)},$$

which gives the **forward Euler method**

$$x^{(k+1)} = x^{(k)} + f^{(k)} \Delta t.$$

- This method requires only one function evaluation per time step.



# Euler's Method

$$\text{Scheme: } x^{(k+1)} - x^{(k)} - f^{(k)} \Delta t = 0$$

- The local truncation error is easy to find using a Taylor series expansion:

$$\begin{aligned} e_k &= x[(k+1)\Delta t] - x(k\Delta t) - f[x(k\Delta t)] \Delta t = \\ &= x[(k+1)\Delta t] - x(k\Delta t) - [x'(k\Delta t)] \Delta t = \frac{x''(\xi)}{2} \Delta t^2, \end{aligned}$$

for some  $k\Delta t \leq \xi \leq (k+1)\Delta t$ .

- Therefore the LTE is  $O(\Delta t^2)$ ,  $q = 2$ .
- The global truncation error, however, is of order  $O(\Delta t)$ ,  $p = q + 1$ , so this is a **first-order accurate** method.

# Backward Euler

$$x^{(k+1)} \approx x^{(k)} + \int_{k\Delta t}^{(k+1)\Delta t} f[x(s), s] ds.$$

- How about we use a piecewise constant-approximation, but based on the end-point:

$$f[x(s), s] \approx f(x^{(k+1)}) = f^{(k+1)},$$

which gives the first-order **backward Euler method**

$$x^{(k+1)} = x^{(k)} + f(x^{(k+1)})\Delta t.$$

- This **implicit method** requires **solving a non-linear equation** at every time step, which is expensive and hard.  
We will understand why implicit methods are needed when we examine absolute stability later on.

# Runge-Kutta Methods

- **Runge-Kutta methods** are a powerful class of **one-step methods** similar to Euler's method, but more accurate.
- As an example, consider using a trapezoidal rule to approximate the integral

$$x^{(k)} + \int_{k\Delta t}^{(k+1)\Delta t} f[x(s), s] ds \approx x^{(k)} + \frac{\Delta t}{2} [f(k\Delta t) + f((k+1)\Delta t)],$$

$$x^{(k+1)} = x^{(k)} + \frac{\Delta t}{2} \left[ f\left(x^{(k)}, t^{(k)}\right) + f\left(x^{(k+1)}, t^{(k+1)}\right) \right]$$

which requires solving a nonlinear equation for  $x^{(k+1)}$ .

- This is the simplest **implicit Runge-Kutta method**, usually called the **implicit trapezoidal method**.
- The local truncation error is  $O(\Delta t^3)$ , so the global error is **second-order accurate**  $O(\Delta t^2)$ .

# Midpoint/Trapezoidal Methods

- Schemes that **treat beginning and end of time step in a symmetric fashion** will lead to a **cancellation of first-order error terms** in Taylor series and will thus be **second order** (Lesson: **second order is easy**).
- In addition to trapezoidal one can do **implicit midpoint** scheme:

$$x^{(k+1)} = x^{(k)} + \frac{\Delta t}{2} f \left( \frac{x^{(k)} + x^{(k+1)}}{2}, t^{(k)} + \frac{\Delta t}{2} \right)$$

Observe this is the same as trapezoidal for linear problems (why?).

- In an explicit method, we would approximate  $x^* \approx x^{(k+1)}$  first using Euler's method, to get the simplest **explicit Runge-Kutta method**, usually called **Heun's or explicit trapezoidal method**

$$x^{(k+1,*)} = x^{(k)} + f \left( x^{(k)}, t^{(k)} \right) \Delta t$$

$$x^{(k+1)} = x^{(k)} + \frac{\Delta t}{2} \left[ f \left( x^{(k)}, t^{(k)} \right) + f \left( x^{(k+1,*)}, t^{(k+1)} \right) \right].$$

# Explicit Midpoint

- **Explicit midpoint** rule

$$x^{(k+\frac{1}{2},*)} = x^{(k)} + f\left(x^{(k)}, t^{(k)}\right) \frac{\Delta t}{2}$$

$$x^{(k+1)} = x^{(k)} + \frac{\Delta t}{2} f\left(x^{(k+\frac{1}{2},*)}, t^{(k)} + \frac{\Delta t}{2}\right).$$

- Explicit midpoint/trapezoidal are a representative of a powerful class of second-order methods called **predictor-corrector methods**: Euler (forward or backward) method is the predictor, and then (implicit or explicit) trapezoidal/midpoint method is the corrector.
- One can also consider these as examples of **multi-stage one-step** methods: the predictor is the first stage, the corrector the second.

# Higher-Order Runge-Kutta Methods

- The idea in RK methods is to evaluate the function  $f(x, t)$  several times and then take a time-step based on an average of the values.
- In practice, this is done by performing the calculation in **stages**: Calculate an intermediate approximation  $x^*$ , evaluate  $f(x^*)$ , and go to the next stage.
- The most celebrated Runge-Kutta method is a **four-stage** fourth-order accurate RK4 method based on **Simpson's rule** for the integral:

$$\begin{aligned}
 x^{(k)} + \int_{k\Delta t}^{(k+1)\Delta t} f[x(s), s] ds \\
 &\approx x^{(k)} + \frac{\Delta t}{6} \left[ f(x^{(k)}) + 4f(x^{(k+1/2)}) + f(x^{(k+1)}) \right] \\
 &= x^{(k)} + \frac{\Delta t}{6} \left[ f^{(k)} + 4f^{(k+1/2)} + f^{(k+1)} \right],
 \end{aligned}$$

and we approximate  $4f^{(k+1/2)} = 2f^{(k+1/2;1)} + 2f^{(k+1/2;2)}$ .

## RK4 Method

$$f^{(k)} = f(x^{(k)}), \quad x^{(k+1/2;1)} = x^{(k)} + \frac{\Delta t}{2} f^{(k)}$$

$$f^{(k+1/2;1)} = f(x^{(k+1/2;1)}, t^{(k)} + \Delta t/2)$$

$$x^{(k+1/2;2)} = x^{(k)} + \frac{\Delta t}{2} f^{(k+1/2;1)}$$

$$f^{(k+1/2;2)} = f(x^{(k+1/2;2)}, t^{(k)} + \Delta t/2)$$

$$x^{(k+1;1)} = x^{(k)} + \Delta t f^{(k+1/2;2)}$$

$$f^{(k+1)} = f(x^{(k+1;1)}, t^{(k)} + \Delta t)$$

$$x^{(k+1)} = x^{(k)} + \frac{\Delta t}{6} \left[ f^{(k)} + 2f^{(k+1/2;1)} + 2f^{(k+1/2;2)} + f^{(k+1)} \right]$$

## Intro to multistep Methods

$$x^{(k+1)} \approx x^{(k)} + \int_{k\Delta t}^{(k+1)\Delta t} f[x(s), s] ds.$$

- Euler's method was based on a piecewise constant approximation (extrapolation) of  $f(s) \equiv f[x(s), s]$ .
- If we instead integrate the linear extrapolation

$$f(s) \approx f(x^{(k)}, t^{(k)}) + \frac{f(x^{(k)}, t^{(k)}) - f(x^{(k-1)}, t^{(k-1)})}{\Delta t} (s - t_k),$$

we get the second-order **two-step Adams-Bashforth** method

$$x^{(k+1)} = x^{(k)} + \frac{\Delta t}{2} \left[ 3f(x^{(k)}, t^{(k)}) - f(x^{(k-1)}, t^{(k-1)}) \right].$$

- This is an example of a **multi-step method**, which requires keeping previous values of  $f$ .



# In MATLAB

- In MATLAB, there are several functions whose names begin with

$$[\mathbf{t}, \mathbf{x}] = \text{ode}(f, [t_0, t_e], x_0, \text{odeset}(\dots)).$$

- *ode23* is a second-order **adaptive explicit** Runge-Kutta method, while *ode45* is a fourth-order version (try it first).
- *ode23tb* is a second-order **implicit** RK method.
- *ode113* is a **variable-order explicit** multi-step method that can provide very high accuracy.
- *ode15s* is a **variable-order implicit** multi-step method.
- For implicit methods the Jacobian can be provided using the *odeset* routine – very important!

# Rigid body motion

```
function dy = rigid(t,y)
dy = zeros(3,1);    % a column vector
dy(1) = y(2) * y(3);
dy(2) = -y(1) * y(3);
dy(3) = -0.51 * y(1) * y(2);
%—————

opts = odeset('RelTol',1e-3,'AbsTol',[1e-4 1e-4 1e-5]);
[T,Y] = ode45(@rigid, [0 12], [0 1 1], opts);

plot(T,Y(:,1),'o-r', T,Y(:,2),'s-b', T,Y(:,3),'d-g');
xlabel('t'); ylabel('y'); title('RelTol=1e-3');
```

## van der Pol equation

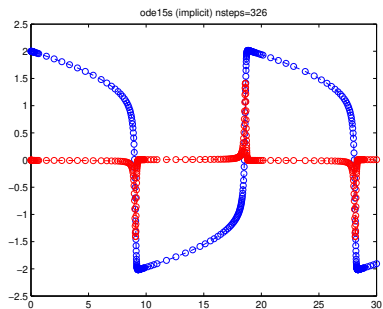
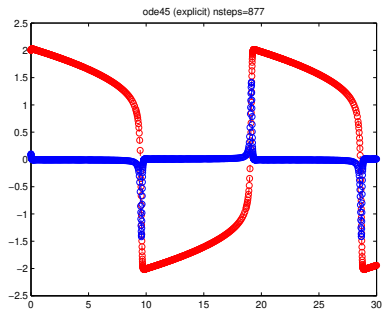
```

r=10; % Try r=100
f = @(t,y) [y(2); r*(1 - y(1)^2)*y(2) - y(1)];

figure(2); clf
[T,Y] = ode45(f,[0 3*r],[2 1]);
plot(T,Y(:,1),'o—r', T,Y(:,2)/r,'o—b')
title(['ode45 (explicit) nsteps=', int2str(size(T,1))]);

figure(3); clf
[T,Y] = ode15s(f,[0 3*r],[2 0]);
plot(T,Y(:,1),'o—b', T,Y(:,2)/r,'o—r')
title(['ode15s (implicit) nsteps=', int2str(size(T,1))]);

```

Stiff van der Pol system ( $r = 10$ )

A stiff problem is one where  $\Delta t$  has to be small even though the solution is smooth and a large  $\Delta t$  is OK for accuracy.

# Zero Stability

- Consistency is not enough — we must also examine **error propagation** from one time step to another.
- A method is called **zero stable** if for all sufficiently small but finite  $\Delta t$ , introducing perturbations at each step (e.g., roundoff errors, errors in evaluating  $f$ ) with magnitude less than some small  $\epsilon$  perturbs the solution by at most  $O(\epsilon)$ .
- This simply means that errors do not increase but rather decrease from step to step.
- Denote our numerical approximation with time step size  $\tau$ :

$$\mathbf{x}^{(k)} \approx \mathbf{x}(k\Delta t).$$

- A method is **convergent** if applying it to any system of ODEs where  $\mathbf{f}$  is Lipschitz over a **finite time interval**  $T > 0$  during which the ODE has a solution, the numerical approximation converges to that solution,

$$\lim_{\Delta t \rightarrow 0} \mathbf{x}^{(N=T/\Delta t)} = \mathbf{x}(T).$$

# Convergence

- A central theorem in numerical methods for differential equations is that  
*Any consistent method is convergent if and only if it is zero stable, or*  
**consistency + (zero) stability = convergence.**
- Note that we haven't given a precise definition to **zero stability** because in some sense it is defined as: the extra conditions that are needed to make a consistent method convergent.
- **Convergence is a statement about a limit, and does not imply a method will give reasonable answers for finite  $\Delta t > 0$ .** For that we will later introduce **absolute stability**.
- It can be shown that **one-step methods are zero-stable** if  $f$  is well-behaved (Lipschitz continuous w.r.t. second argument).

# Stiff example

- In section 7.1 LeVeque discusses

$$x'(t) = \lambda(x - \cos t) - \sin t.$$

with solution  $x(t) = \cos t$  if  $x(0) = 1$ .

- If  $\lambda = 0$  then this is very simple to solve using Euler's method, for example,  $\Delta t = 10^{-3}$  up to time  $T = 2$  gives error  $\sim 10^{-3}$ .
- For  $\lambda = -10$ , one gets an even smaller error with the same time step size.
- But for  $\lambda = -2100$ , results for  $\Delta t > 2/2100 = 0.000954$  are completely useless: **method is unstable**.

# Conditional Stability

- Consider the model problem for  $\lambda < 0$ :

$$x'(t) = \lambda x(t)$$

$$x(0) = 1,$$

with an exact solution that **decays exponentially**,  $x(t) = e^{\lambda t}$ .

- Applying Euler's method to this model equation gives:

$$x^{(k+1)} = x^{(k)} + \lambda x^{(k)} \Delta t = (1 + \lambda \Delta t) x^{(k)} \Rightarrow$$

$$x^{(k)} = (1 + \lambda \Delta t)^k$$

- The numerical solution will **decay** if the time step satisfies the **stability criterion**

$$|1 + \lambda \Delta t| \leq 1 \Rightarrow \Delta t < -\frac{2}{\lambda}.$$

- Otherwise, the numerical solution will eventually blow up!



# Unconditional Stability

- The above analysis shows that **forward Euler is conditionally stable**, meaning it is stable if  $\Delta t < 2/|\lambda|$ .
- Let us examine the stability for the model equation  $x'(t) = \lambda x(t)$  for **backward Euler**:

$$x^{(k+1)} = x^{(k)} + \lambda x^{(k+1)} \Delta t \quad \Rightarrow \quad x^{(k+1)} = x^{(k)} / (1 - \lambda \Delta t)$$

$$x^{(k)} = x^{(0)} / (1 - \lambda \Delta t)^k$$

- We see that the implicit **backward Euler is unconditionally stable**, since for any time step

$$|1 - \lambda \Delta t| > 1.$$

# Stiff Systems

- An ODE or a system of ODEs is called **stiff** if the solution evolves on widely-separated timescales and the fast time scale decays (dies out) quickly.
- We can make this precise for linear systems of ODEs,  $\mathbf{x}(t) \in \mathbb{R}^n$ :

$$\mathbf{x}'(t) = \mathbf{A} [\mathbf{x}(t)].$$

- Assume that  $\mathbf{A}$  has an eigenvalue decomposition, with potentially **complex eigenvalues**:

$$\mathbf{A} = \mathbf{X}\mathbf{\Lambda}\mathbf{X}^{-1},$$

and express  $\mathbf{x}(t)$  in the basis formed by the eigenvectors  $\mathbf{x}_i$ :

$$\mathbf{y}(t) = \mathbf{X}^{-1} [\mathbf{x}(t)].$$

contd.

$$\mathbf{x}'(t) = \mathbf{A}[\mathbf{x}(t)] = \mathbf{X}\mathbf{\Lambda}[\mathbf{X}^{-1}\mathbf{x}(t)] = \mathbf{X}\mathbf{\Lambda}[\mathbf{y}(t)] \quad \Rightarrow$$

$$\mathbf{y}'(t) = \mathbf{\Lambda}[\mathbf{y}(t)]$$

- The different  $y$  variables are now **uncoupled**: each of the  $n$  ODEs is independent of the others:

$$y_i = y_i(0)e^{\lambda_i t}.$$

- Assume for now that all eigenvalues are **real and negative**,  $\lambda < 0$ , so each component of the solution decays:

$$\mathbf{x}(t) = \sum_{i=1}^n y_i(0)e^{\lambda_i t} \mathbf{x}_i \quad \rightarrow \quad 0 \text{ as } t \rightarrow \infty.$$

# Stiffness

- If we solve the original system using Euler's method, the time step must be smaller than the smallest stability limit,

$$\Delta t < \frac{2}{\max_i |\operatorname{Re}(\lambda_i)|}.$$

- A system is **stiff** if there is a strong **separation of time scales** in the eigenvalues:

$$r = \frac{\max_i |\operatorname{Re}(\lambda_i)|}{\min_i |\operatorname{Re}(\lambda_i)|} \gg 1.$$

- For non-linear problems  $\mathbf{A}$  is replaced by the Jacobian  $\nabla_{\mathbf{x}}\mathbf{f}(\mathbf{x}, t)$ .
- In general, the Jacobian will have **complex eigenvalues** as well.

# Absolute Stability

- We see now that for systems we need to allow  $\lambda$  to be a **complex number** but we can still look at scalar equations.
- A method is called **absolutely stable** if for  $\text{Re}(\lambda) < 0$  the numerical solution of the **scalar model equation**

$$x'(t) = \lambda x(t)$$

decays to zero, like the actual solution.

- We call the **region of absolute stability** the set of complex numbers

$$z = \lambda \Delta t$$

for which the numerical solution decays to zero.

- For systems of ODEs **all scaled eigenvalues of the Jacobian  $\lambda_i \Delta t$  should be in the stability region.**

# Stability regions

- For Euler's method, the stability condition is

$$|1 + \lambda\Delta t| = |1 + z| = |z - (-1)| \leq 1 \quad \Rightarrow$$

which means that  $z$  must be in a unit disk in the complex plane centered at  $(-1, 0)$ :

$$z \in \mathcal{C}_1(-1, 0).$$

- A general one-step method of order  $p$  applied to the **model equation**  $x' = \lambda x$  where  $\lambda \in \mathbb{C}$  gives:

$$x^{n+1} = R(z = \lambda\Delta t)x^n.$$

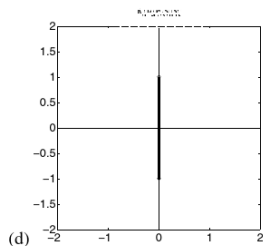
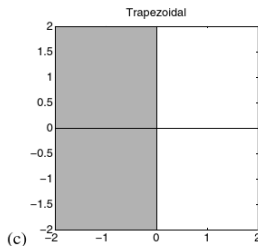
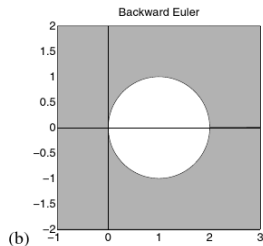
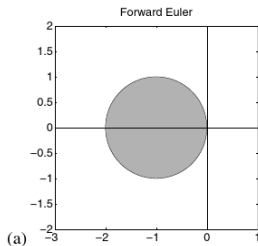
$$R(z) = e^z + O(z^{p+1}) \text{ for small } |z|.$$

- The **region of absolute stability** is the set

$$\mathcal{S} = \{z \in \mathbb{C} : |R(z)| \leq 1\}.$$

# Simple Schemes

Forward/backward Euler, **implicit trapezoidal**, and leapfrog schemes



# A-Stable methods

- A method is **A-stable** if its stability region contains the entire left half plane.
- The backward Euler and the implicit midpoint scheme are both A-stable, but they are also both implicit and thus expensive in practice!
- Theorem: **No explicit one-step method can be A-stable** (discuss in class why).
- Theorem: All explicit RK methods with  $r$  stages and of order  $r$  have the same stability region (discuss why).



# One-Step Methods

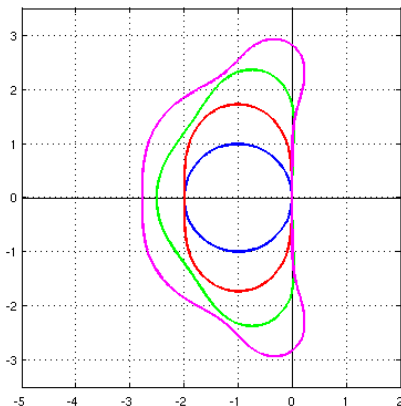
- Any  $r$ -stage **explicit** RK method will produce  $R(z)$  that is a **polynomial** of degree  $r$ .
- Any  $r$ -stage **implicit** RK method has **rational**  $R(z)$  (ratio of polynomials).  
The degree of the denominator cannot be larger than the **number of linear systems that are solved** per time step.
- RK methods give polynomial or rational approximations  $R(z) \approx e^z$ .
- A 4-stage explicit RK method therefore has

$$R(z) = 1 + z + \frac{1}{2}z^2 + \frac{1}{6}z^3 + \frac{1}{24}z^4$$

# Explicit RK Methods

Stability regions for **all**  $r$ -stage explicit RK methods

Runge-Kutta orders 1,2,3,4



**One needs at least 3 stages to be stable for purely imaginary eigenvalues (hyperbolic PDEs later on).**

# Implicit Methods

- **Implicit methods are generally more stable** than explicit methods, and solving stiff problems generally requires using an implicit method.
- The price to pay is solving a system of non-linear equations at every time step (linear if the ODE is linear):  
This is best done using **Newton-Raphson's** method, where the solution at the previous time step is used as an initial guess.
- For PDEs, the linear systems become large and implicit methods can become very expensive...

# Implicit-Explicit Methods

- When solving PDEs, we will often be faced with problems of the form

$$\frac{dx}{dt} = \mathbf{f}(\mathbf{x}, t) + \mathbf{g}(\mathbf{x}, t) = \text{stiff} + \text{non-stiff}$$

where the stiffness comes only from  $\mathbf{f}$ .

- These problems are treated using **implicit-explicit (IMEX)** or **semi-implicit** schemes, which only treat  $\mathbf{f}(\mathbf{x})$  implicitly (see HW4 for KdV equation).
- A very simple example of a second-order scheme is to treat  $\mathbf{g}(\mathbf{x})$  using the **Adams-Bashforth** multistep method and treat  $\mathbf{f}(\mathbf{x})$  using the implicit trapezoidal rule (**Crank-Nicolson** method), the **ABCN** scheme:

$$x^{(k+1)} = x^{(k)} + \frac{\Delta t}{2} \left[ \mathbf{f}(x^{(k)}, t^{(k)}) + \mathbf{f}(x^{(k+1)}, t^{(k+1)}) \right] \\ + \Delta t \left[ \frac{3}{2} \mathbf{g}(x^{(k)}, t^{(k)}) - \frac{1}{2} \mathbf{g}(x^{(k-1)}, t^{(k-1)}) \right].$$

# Which Method is Best?

- As expected, there is no universally “best” method for integrating ordinary differential equations: It depends on the problem:
  - How stiff is your problem (may demand implicit method), and does this change with time?
  - How many variables are there, and how long do you need to integrate for?
  - How accurately do you need the solution, and how sensitive is the solution to perturbations (chaos).
  - How well-behaved or not is the function  $f(x, t)$  (e.g., sharp jumps or discontinuities, large derivatives, etc.).
  - How costly is the function  $f(x, t)$  and its derivatives (Jacobian) to evaluate.
  - Is this really ODEs or a something coming from a PDE integration (next lecture)?

# Conclusions/Summary

- Time stepping methods for ODEs are **convergent if and only if they are consistent and stable**.
- We distinguish methods based on their **order of accuracy** and on whether they are **explicit** (forward Euler, Heun, RK4, Adams-Bashforth), or **implicit** (backward Euler, Crank-Nicolson), and whether they are **adaptive**.
- **Runge-Kutta methods** require more evaluations of  $f$  but are more robust, especially if adaptive (e.g., they can deal with sharp changes in  $f$ ). Generally the recommended first-try (*ode45* or *ode23* in MATLAB).
- **Multi-step methods** offer high-order accuracy and require few evaluations of  $f$  per time step. They are not very robust however. Recommended for well-behaved non-stiff problems (*ode113*).
- For **stiff problems** an **implicit method** is necessary, and it requires solving (linear or nonlinear) systems of equations, which may be complicated (evaluating Jacobian matrices) or costly (*ode15s*).