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SPLITTING METHODS

$$u_t + \nabla \cdot (\underline{a} u) = \nabla \cdot (D \nabla u) + f(u)$$

w
non-stiff
advection

s
shift
diffusion

v
very stiff
chemistry

Splitting methods use different
temporal integrators for different
parts.

Main caveat:

Only "easy" to do up to
second order accuracy!

SPLITTING
ERRORS

②

$$\omega' = A \omega = (A_1 + A_2) \omega$$

$$\begin{aligned} \omega(t_{n+1}) &= e^{\tau A} \omega(t_n) \\ &= e^{\tau (A_1 + A_2)} \omega(t_n) \end{aligned}$$

SPLITTING} $\approx e^{\tau A_2} e^{\tau A_1} w_n$

$$\boxed{w^{n+1} = e^{\tau A_2} e^{\tau A_1} w^n} \leftarrow \text{Lie SPLITTING}$$

This is in general only first-order accurate even if A_1/A_2 exponentials handled exactly!

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BCH formula:

$$e^{\bar{i}A_2} e^{\bar{i}A_1} \underset{\sim}{=} e^{\bar{i}\tilde{A}}$$

$$\tilde{A} = A + \frac{\bar{i}}{2} [A_2, A_1] + O(\bar{i}^2)$$

$$[A_2, A_1] = A_2 A_1 - A_1 A_2 \quad \text{commutator}$$

\Rightarrow Lie splitting exactly solves
the modified equation

$$\left\{ \begin{array}{l} \omega' = \tilde{A} \omega = Aw - \frac{\bar{i}}{2} [A_2, A_1] w \\ \text{so error is } O(\bar{i}) \end{array} \right.$$

④ One can do second-order splitting by symmetrizing the order of splitting to eliminate odd powers of τ :

$$w^{n+1} = e^{\frac{\tau A_1}{2}} e^{\tau A_2} e^{\frac{\tau A_1}{2}} w^n$$

STRANG splitting

Local truncation error

$$S^n = \frac{\tau^2}{24} ([A_1, [A_1, A_2]] + 2[A_2, [A_1, A_2]]) w^{n+1/2}$$

or

$$w^{n+1} = \frac{1}{2} [e^{\tau A_1} e^{\tau A_2} + e^{\tau A_2} e^{\tau A_1}] w^n$$

⑤ One can do similar tricks with multi-component splitting

$$A = A_1 + A_2 + A_3$$

$$\omega^{n+1} = e^{\bar{\tau}A_1/2} e^{\bar{\tau}A_2/2} e^{\bar{\tau}A_3/2} e^{\bar{\tau}A_2/2} e^{\bar{\tau}A_1/2} \omega^n$$

Some results for advection-diffusion

① Advection commutes with diffusion only if $\underline{a} = \text{const}$ and $D = \text{const}$

② Advection commutes with local terms (reaction, forcing, etc.) if $\nabla \cdot \underline{a} = 0$ and forcing is autonomous $f = f(u)$

⑥

Non linear case

$$u_t(x, t) = f[x, u(x, t)]$$

$$f = f_1 + f_2$$

$$u^{n+1} = S_{\bar{\tau}}^2 [S_{\bar{\tau}}^1(u^n)] \leftarrow \begin{array}{l} \text{lie} \\ \text{splitting} \end{array}$$

one-step update
for $u_t = f_1$ and $u_t = f_2$

Local truncation error

$$S_n = \frac{\bar{\tau}}{2} \left[\frac{\partial f_1}{\partial u} f_2 - \frac{\partial f_2}{\partial u} f_1 \right] u^n = \frac{\bar{\tau}}{2} [f_1 f_2] u^n$$

commutator for
functions

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Problems with splitting:

- 1) Higher-order integrators are complicated
- 2) For PDEs, the physical boundary conditions are given for the whole problem, NOT for each of the split problems
(think advection-diffusion splitting)

$$w' = Aw + g(t) = (A_1 w + g_1) + (A_2 w + g_2)$$

$$\begin{aligned} w^{n+1} &= e^{\bar{\tau} A_2} e^{\bar{\tau} A_1} w^n + e^{\bar{\tau} A_2} \int_0^{\bar{\tau}} e^{(\bar{\tau}-s) A_1} g_1(s) ds \\ &\quad + \int_0^{\bar{\tau}} e^{(\bar{\tau}-s) A_2} g_2(s) ds \neq \text{exact} \end{aligned}$$

⑧ So now we have an additional splitting error due to BCs; even if $[A_1, A_2] = 0$!

$$\text{error} \approx \int_0^T e^{(t-s)A} g(s) ds -$$

$$\left[e^{\tau A_2} \int_0^{\tau} e^{(\tau-s)A_1} g_1(s) ds + \int_0^{\tau} e^{(t-s)A_2} g_2(s) ds \right]$$

3) If A_1 has eigenvalues proportional to $-1/\epsilon$, $\tau/\epsilon \gg 1$ (very stiff problem), then even Strang splitting is first order (reduction of order)