

SERIES-EXPANSION METHODS

(1)

CFD SPRING 2013, A. DONEV

This is based on ch. 4 in "Numerical methods for wave equations" by Durran.

We want to solve PDE

$$\frac{\partial \psi}{\partial t} \neq F(\psi) = 0$$

We have a numerical approximation

$$\psi \approx \tilde{\psi}$$

and want to expand / represent it in some finite basis for the functional space of interest.

$$\Psi(x, t) = \sum_{k=1}^N a_k(t) \Psi_k(x)$$

(2)

↑
basis functions

Residual

$$R(\Psi) = \frac{\partial \Psi}{\partial t} + F(\Psi)$$

should somehow be

minimized

① Equivalent: Minimize $R(\Psi)$ in L_2 norm
 or require that $R(\Psi)$ be orthogonal
 to the finite-dimensional subspace
 spanned by $\{\Psi_k(x)\}$: GALERKIN

② Collocation: $R[\Psi(j\Delta x)] = 0, j=1, \dots, N$

We adopt the GALERKIN approach. (3)

Minimize over $\dot{a}_k = \frac{da_k(t)}{dt}$ s.t.

$$\int \mathcal{R}[\psi(x)] \psi_k(x) dx = 0 =$$

$$\int \left[\sum_{n=1}^N \dot{a}_n \psi_n + F\left(\sum_{n=1}^N a_n \psi_n\right) \right] \psi_k dx$$

$$\Rightarrow \sum_n M_{nk} \dot{a}_n = - \int F\left(\sum_n a_n \psi_n\right) \psi_k dx$$

for $k = 1, \dots, N$

LINEAR SYSTEM of eqs.

$$M_{nk} = \int \psi_n \psi_k dx \quad (\text{mass matrix}) \quad (4)$$

Note that if the basis functions are orthonormal, then M is the identity matrix, $M = I$

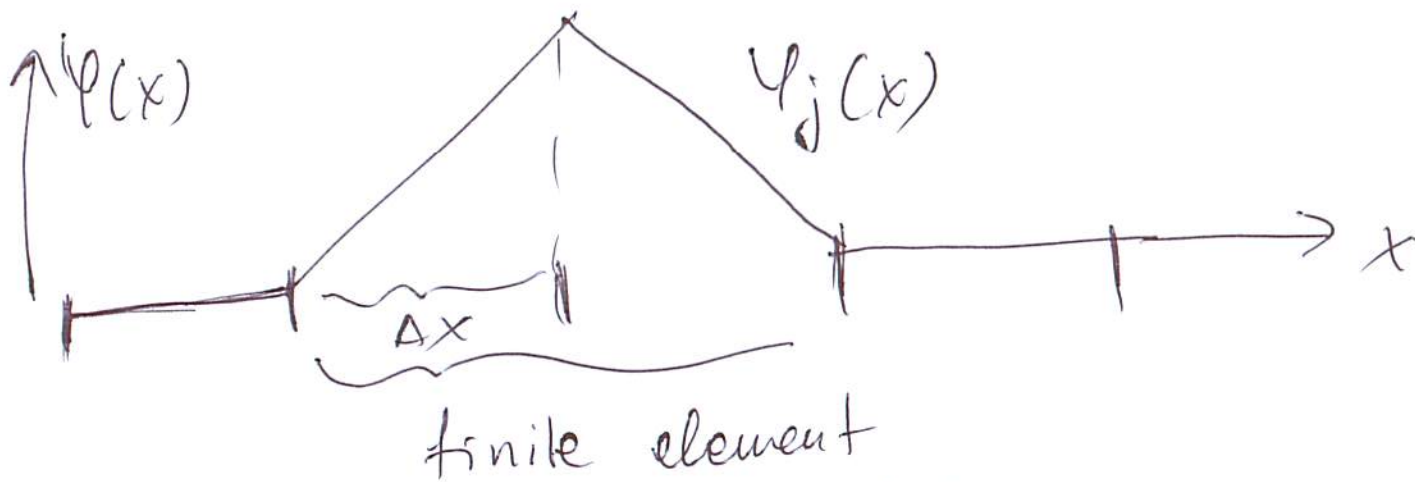
We have now converted the PDE into a system of ODEs for $a_k(t)$
→ spatial discretization

Depending on the choice of basis functions, we can be doing spectral, finite-element, etc.

Finite - Element Method

(8)

The main difference with (pseudo) spectral is that now we choose a localized basis function set; e.g.



$$\int \psi_j \psi_{j+1} dx = \int_0^{\Delta x} \frac{x}{\Delta x} \frac{\Delta x - x}{\Delta x} dx = \frac{\Delta x}{6}$$

$$\int \psi_j^2 dx = \frac{2\Delta x}{3} \quad \text{and we also need}$$

$$-\int \frac{\partial \psi_{j-1}}{\partial x} \psi_j dx = \frac{1}{2} \quad (9)$$

can be done using integration by parts
if ψ is not differentiable

So for $\frac{\partial \psi}{\partial t} + c \frac{\partial \psi}{\partial x} = 0$ this gives
the Galerkin finite-element discretisation

$$\frac{\dot{a}_{j+1} + 4\dot{a}_j + \dot{a}_{j-1}}{6} + c \left(\frac{a_{j+1} - a_{j-1}}{2\Delta x} \right) = 0$$

sparse matrix

Note that $a_j \equiv \psi_j$ for this basis

Denote

$$M = \frac{1}{6} \begin{bmatrix} 4 & & & & \\ & 1 & & & \\ & & -4 & & \\ & & & 1 & \\ & & & & 4 \end{bmatrix} \Rightarrow$$

(10)

$$\dot{a} = \dot{\psi} = -M^{-1} (c \Delta a) = -c (M^{-1} \Delta) \psi$$

centered finite difference

$$\Rightarrow \boxed{\dot{\psi} = -c (M^{-1} \Delta) \psi}$$

Finite-difference Approximation of $\frac{\partial}{\partial x} \rightarrow$
called "compact finite difference"
in the literature.

It turns out $M^{-1} \Delta$ is a
fourth-order approximation of $\frac{\partial}{\partial x}$!

So this finite-element method is fourth order in space.

(11)

An alternative approach to minimizing the residual $R(\psi) = \frac{\partial \psi}{\partial t} + \frac{\partial \psi}{\partial x}$ is

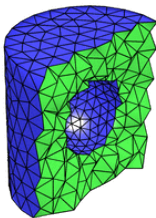
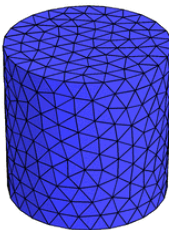
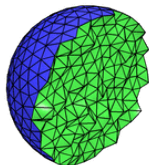
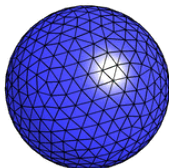
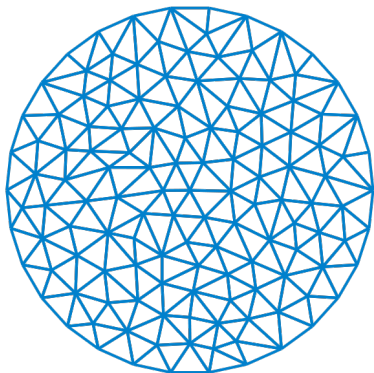
to replace the time derivative right-away with a difference:

$$\left\{ \begin{array}{l} \frac{\psi^{n+1} - \psi^n}{\Delta t} + c \frac{\partial \psi}{\partial x} = 0 \dots (*) \\ \text{where } \psi^n = \sum a_k^n \psi_k(x) \end{array} \right.$$

Now minimize the residual in $(*)$ to get $\psi^{n+1} \leftarrow \psi^n$, i.e., $a^{n+1} \leftarrow a^n$

Irregular (Simplicial) Meshes

Any polygon can be triangulated into arbitrarily many **disjoint triangles**.
Similarly **tetrahedral meshes** in 3D.



Basis functions on triangles

- For irregular grids the x and y directions are no longer separable.
- But the idea of using basis functions $\phi_{i,j}$, a **reference triangle**, and **piecewise polynomial interpolants** still applies.
- For a linear function we need 3 coefficients (x, y, const), for quadratic 6 ($x, y, x^2, y^2, xy, \text{const}$):

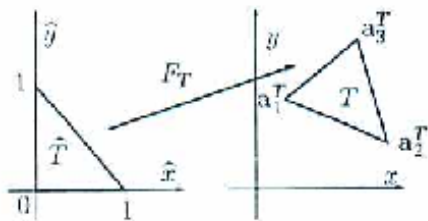


Fig. 8.8. Local interpolation nodes on \hat{T} for $k=0$ (left), $k=1$ (center), $k=2$ (right)

Piecewise constant / linear basis functions

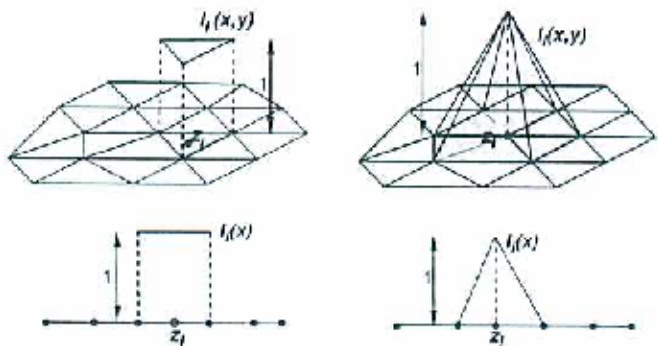


Fig. 8.7. Characteristic piecewise Lagrange polynomial, in two and one space dimensions. Left, $k = 0$; right, $k = 1$

