

Improved Lower Bounds for Ginzburg-Landau Energies via Mass Displacement

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Abstract

We prove some improved estimates for the Ginzburg-Landau energy (with or without magnetic field) in two dimensions, relating the asymptotic energy of an arbitrary configuration to its vortices and their degrees, with possibly unbounded numbers of vortices. The method is based on a localisation of the “ball construction method” combined with a mass displacement idea which allows to compensate for negative errors in the ball construction estimates by energy “displaced” from close by. Under good conditions, our main estimate allows to get a lower bound on the energy which includes a finite order “renormalized energy” of vortex interaction, up to the best possible precision i.e. with only a $o(1)$ error per vortex, and is complemented by local compactness results on the vortices. This is used crucially in our forthcoming paper [SS5]. It can also serve to provide lower bounds for weighted Ginzburg-Landau energies.

keywords: Ginzburg-Landau, vortices, vortex balls construction, renormalized energy, mass displacement.

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Introduction

We are interested in proving lower bounds and compactness results for Ginzburg-Landau type energies of the form

$$G_\varepsilon(u, A) = \frac{1}{2} \int_{\Omega_\varepsilon} |\nabla_A u|^2 + (\operatorname{curl} A)^2 + \frac{(1 - |u|^2)^2}{2\varepsilon^2}$$

where ε is a small parameter, u is a complex-valued function called “order parameter”, A is \mathbb{R}^2 valued and is the vector potential of the magnetic field $h := \operatorname{curl} A$ and $\nabla_A = \nabla - iA$. Here the domain of integration Ω_ε is a smooth bounded domain in \mathbb{R}^2 , which can depend

on ε . We are interested in particular in the case where Ω_ε gets large as $\varepsilon \rightarrow 0$. Note that one may set $A \equiv 0$ to recover the simpler Ginzburg-Landau energy

$$E_\varepsilon(u) = \frac{1}{2} \int_{\Omega_\varepsilon} |\nabla u|^2 + \frac{(1 - |u|^2)^2}{2\varepsilon^2}$$

without magnetic field. Our results apply to this energy functional by making this trivial choice of A .

The Ginzburg-Landau energy is a famous model for superconductivity. In this model the order-parameter u often has quantized vortices, which are the zeroes of u with nonzero topological degree. Obtaining ansatz-free lower bounds for G_ε in terms of the vortices of u has proven to be crucial in studying the asymptotics of minimizers of G_ε , in particular via Γ -convergence methods.

The first study establishing lower bounds for Ginzburg-Landau was the work of Bethuel-Brezis-Hélein [BBH] for solutions to the Ginzburg-Landau equations without magnetic field with energy E_ε bounded by $C|\log \varepsilon|$. Such an energy bound ensures that the total number of vortices remains bounded as $\varepsilon \rightarrow 0$. This was later improved and extended in two different directions by Han-Shafrir [HS] and Almeida-Bethuel [AB] for arbitrary configurations, still with a number of vortices that remains bounded. The main limitation of such estimates is that the error terms blow up as the number of vortices gets large. Then, Jerrard [Je] and Sandier [Sa] introduced the “ball construction method”, which provides lower bounds in terms of vortices for arbitrary configurations, allowing unbounded numbers of vortices and much larger energies. This is crucial for many applications, since energy minimizers of the functional with applied magnetic field do not always satisfy a $C|\log \varepsilon|$ bound on their energy. Subsequent refinements of the ball construction method were given (see for example [SS4] Chap. 4 for a recent result). The lower bound provided by the ball construction method also provides a crucial compactness result on the vorticity (roughly the sum of Dirac masses at the vortex centers, weighted by their degrees), these are the so-called “Jacobian estimates”, see Jerrard-Soner [JS] and [SS4] Chap. 6 and references therein. They say roughly that the vorticity is controlled by $\frac{1}{|\log \varepsilon|}$ times the energy. For other subsequent works refining those results in a slightly different direction, see also [SS3, JSp, ST].

In a way our objective here can be seen as obtaining next order terms (order 1 as opposed to order $|\log \varepsilon|$) in such estimates, both energy estimates and compactness results.

For a given (u, A) , let us define the energy density

$$e_\varepsilon(u, A) = \frac{1}{2} \left(|\nabla_A u|^2 + (\operatorname{curl} A)^2 + \frac{(1 - |u|^2)^2}{2\varepsilon^2} \right).$$

If (u, A) is clear from the context and defined on a set E , we will often use the abbreviation $e_\varepsilon(E)$ for $\int_E e_\varepsilon(u, A)$, and e_ε for the density $e_\varepsilon(u, A)$. We then introduce the measure

$$f_\varepsilon := e_\varepsilon - \pi |\log \varepsilon| \sum_B d_B \delta_{a_B}$$

where the a_B 's are the centers of the vortex balls constructed via Jerrard's and Sandier's ball construction, the d_B 's are the degrees of the balls and δ is the Dirac mass. Calculating $\int f_\varepsilon$ corresponds to subtracting off the cost of all vortices from the total energy: what remains should then correspond to the interaction energy between the vortices, which we can call "renormalized energy" by analogy with [BBH]. In order to obtain next order estimates of the energy G_ε , we show here lower bounds on the energy $\int f_\varepsilon$, as well as coerciveness properties of f_ε , which say, roughly, that f_ε , or in other words, the renormalized energy, suffices to control the vorticity. (This is again to be compared with the previous ball construction and Jacobian estimate, where the vorticity is controlled by $e_\varepsilon/|\log \varepsilon|$).

The motivation for this is our joint paper [SS5] where we establish a "next order" Γ -convergence result for the Ginzburg-Landau energy with applied magnetic field, and derive a limiting interaction energy between points in the plane, thus making the link to the question of the famous Abrikosov lattice (the Abrikosov lattice is a hexagonal lattice of vortices in superconductors observed in experiments and predicted by Abrikosov). More precisely, we show in [SS5] an asymptotic expansion for the minimal energy of the form

$$\min G_\varepsilon = G_\varepsilon^N + N \min W + o(N)$$

where $N \gg 1$ is the optimal number of vortices (determined by the intensity of the applied field), G_ε^N is a constant of order N^2 (the leading order estimate) and W is a renormalized energy governing the pattern formed by the vortices after blow-up at the scale \sqrt{N} . Moreover, we show that the patterns formed by the vortices of minimizers after this blow-up minimize W (almost surely, in some sense). We prove in addition that among lattice configurations (of fixed volume), W is uniquely minimized by the hexagonal lattice. The natural conjecture is that this lattice is also a minimizer among all point configurations, and if this were proved, it would completely justify the emergence of the Abrikosov hexagonal lattice.

To achieve this, with an error only $o(N)$, we needed lower bounds on the cost of vortices with a precision $o(1)$ per vortex (with still a possibly infinite number of vortices), which is finer than was available in the literature. We also needed to control the (local number of) vortices by the renormalized energy. In fact the energy density we end up having to analyze in [SS5] is exactly f_ε , and we need to be able to control the vortices through it.

The other problem we need to overcome for [SS5] is that f_ε is obviously not positive or even bounded below, and this prevents from applying standard lower semi-continuity ideas, and the abstract scheme for Γ -convergence of 2-scale energies which we introduce in [SS5]. This reflects the fact that the energy e_ε is not exactly where the vortices are, as we will explain below. The remedy which we implement here, is that we can "deform" f_ε into an energy density g_ε which is bounded below and enjoys nice coerciveness properties. To accomplish this we show that we can transport the positive mass in f_ε into the support of the negative mass in f_ε , with mass travelling at most at fixed finite distances (say distance 1), and so that the result of the operation, g_ε , is bounded below. This is done by using the following rather elementary transport lemma:

Lemma 0.1. *Assume f is a finite Radon measure on a compact set A , that Ω is open and that for any positive Lipschitz function ξ in $\text{Lip}_\Omega(A)$, i.e. vanishing on $\Omega \setminus A$,*

$$\int \xi df \geq -C_0 |\nabla \xi|_{L^\infty(A)}.$$

Then there exists a Radon measure g on A such that $0 \leq g \leq f_+$ and such that

$$\|f - g\|_{\text{Lip}_\Omega(A)^*} \leq C_0.$$

Thus what is needed is a control on the negative part of f_ε , which will be provided by the ball construction lower bounds and additional improvements of it.

The norm $\|f_\varepsilon - g_\varepsilon\|_{\text{Lip}_\Omega(\Omega)^*}$ will measure how far mass has been displaced in the process. This control appears in Theorem 1 below and more particularly Corollary 1.1. Since $\int g_\varepsilon$ will be close to $\int f_\varepsilon$, it also can be seen as a renormalized energy. Since g_ε is bounded below, we can then hope that it enjoys nice coerciveness properties, we can in fact obtain the desired compactness results which allow to control the vorticity locally by g_ε . This will be the object of Theorem 2 below.

Finally, let us point out that our results can in principle serve to obtain lower bounds for weighted Ginzburg-Landau energies, see Remark 1.4.

Let us now describe a little bit the method that we use, which will allow to control the negative part of f_ε as needed.

The best vortex ball construction lower bound on e_ε available (such as that in [SS4] Chap. 4) is of the following type: given $(u_\varepsilon, A_\varepsilon)$ and any (small) number r , there exists a family of disjoint closed balls \mathcal{B} covering all the zeros of u_ε , the sum of the radii of the balls being bounded above by r , and such that

$$(0.1) \quad \int_{\cup_{B \in \mathcal{B}} B} e_\varepsilon(u_\varepsilon, A_\varepsilon) \geq \pi D \left(\log \frac{r}{\varepsilon D} - C \right),$$

where $D = \sum_{B \in \mathcal{B}} |d_B|$ with $d_B = \deg(u_\varepsilon, \partial B)$ if $B \subset \Omega$ and 0 otherwise. We shall reprove here in Proposition 2.1 a version of this result using Jerrard's ball construction.

This above estimate says that a vortex of degree d costs an energy at least $\simeq \pi |d| |\log \varepsilon|$, but this is only really true when the vortex is well isolated from other vortices and from the boundary, and if there are not too many of them locally, as the factor r/D in the logarithm above somewhat reflects: an ideal lower bound would be

$$e_\varepsilon(B) \geq \pi |d_B| \left(\log \frac{r}{\varepsilon} - C \right),$$

and compared to this, the lower bound above contains a negative error $-\pi D \log D$ which tends to $-\infty$ if the total number of vortices becomes large when $\varepsilon \rightarrow 0$. In truth, this ideal lower bound cannot hold in general as can be seen in the case of n vortices of degree 1 all positioned regularly near the boundary of the domain, a case where (0.1) is optimal.

Moreover the energy density e_ε is not localized exactly where the vortices are: vortices can be viewed as points, while their energy is spread over annular regions around these

points. The ball construction lower bounds such as (0.1) capture well the energy which lies very near the vortices, but some energy is missing from it, in particular when vortices accumulate locally around a point. The missing energy in that case can be recovered by the method of “lower bounds on annuli” which we introduced in [SS1] and re-used in [SS4], Chap. 9. It is based on the following: let $B(x_0, r_1) \setminus B(x_0, r_0)$ be an annulus which contains no zeros of u , roughly speaking we have

$$e_\varepsilon(B(x_0, r_1) \setminus B(x_0, r_0)) \geq \pi D^2 \log \frac{r_1}{r_0}$$

where $D = \deg(u, \partial B(x_0, r_1)) = \deg(u, \partial B(x_0, r_0))$. In other words, if a fixed size ball in the domain contains some large degree D of vorticity, then there is an energy of order D^2 lying not in that ball, but in a thick enough annulus around that ball. This energy of order D^2 should suffice to “neutralize” the error term $-\pi D \log D$ found above through the ball construction. However, it lies at a certain (finite) distance from the center of the vortices. The main technique is then to combine in a systematic way the ball construction lower bounds and the “lower bounds on annuli”, in order to recover enough energy.

Let us finally emphasize a technical difficulty. Since we want a local control on the vortices, the lower bound (0.1) provided by the ball construction is not quite sufficient because it cannot be localized in general, i.e. we cannot deduce a bound for $\int_B e_\varepsilon$ for each $B \in \mathcal{B}$. It is only possible when a matching upper bound on the total in (0.1) is known, see Proposition 2.1 for more details.

The idea to remedy this difficulty is to “localize” the construction, i.e. split the domain into pieces on which one expects to have a bounded vorticity, then apply the ball construction on each piece, and paste together the constructions and lower bounds obtained this way, whose error terms will now be bounded below by a constant. However, this is not completely easy: one needs to localize the construction and still get a global covering of the vortices by balls while preserving the disjointness of the balls. In applications, trying to split the domain into pieces where the vorticity is expected to be bounded leads us to splitting the domain into very small (as $\varepsilon \rightarrow 0$) pieces. Equivalently after rescaling one can consider very large domains cut into bounded size pieces. In other words, in order to be able to treat the case where the vortex density becomes large, we need to be able to treat the case of unbounded domains as $\varepsilon \rightarrow 0$.

This is precisely what we do in this paper: we consider possibly large domains. This way we may in practice rescale our domains as much as needed until the local density of vortices remains bounded as $\varepsilon \rightarrow 0$. We consider vortex ball constructions obtained over coverings of Ω_ε by domains of fixed size, and we work at pasting together these lower bounds while combining them with the method of lower bounds on annuli, as explained above, and finally retrieving “finite numbers of vortices” estimates (of [BBH] type) which bound from below the energy f_ε or g_ε by the exact renormalized energy of [BBH] type up to only $o(1)$ errors.

1 Statement of the main results

In this paper we will deal with families $(u_\varepsilon, A_\varepsilon)_\varepsilon$ defined on domains $\{\Omega_\varepsilon\}_\varepsilon$ in \mathbb{R}^2 which become large as $\varepsilon \rightarrow 0$. The example we have in mind is $\Omega_\varepsilon = \lambda_\varepsilon \Omega$ where Ω is a fixed bounded smooth domain and $\lambda_\varepsilon \rightarrow +\infty$ as $\varepsilon \rightarrow 0$, but we don't need to make any particular hypothesis on $\{\Omega_\varepsilon\}_\varepsilon$, which could even be a fixed bounded domain.

Next we introduce some notation.

For $E \subset \mathbb{R}^2$ we let

$$\widehat{E} = \{x \in \Omega_\varepsilon, \text{dist}(x, E) \leq 1\}.$$

We also define, for any real-valued or vector-valued function f in Ω_ε ,

$$\widehat{f}(x) = \sup\{|f(y)|, y \in B(x, 1) \cap \Omega_\varepsilon\}.$$

Note that both \widehat{f} and \widehat{E} depend on ε , but the value of ε will be clear from the context. The choice of 1 in the definitions is arbitrary but constrains the choice of other constants below.

In all the paper, f_+ and f_- will denote the positive and negative parts of a function or measure, both being positive functions or measures, and $\|f\|$ is the total variation of f . If f and g are two measures then $f \leq g$ means that $g - f$ is a positive measure.

Given a family $\{(u_\varepsilon, A_\varepsilon)\}_\varepsilon$, where $u_\varepsilon : \Omega_\varepsilon \rightarrow \mathbb{C}$ and $A_\varepsilon : \Omega_\varepsilon \rightarrow \mathbb{R}^2$ we define the *currents* and *vorticities* to be

$$j_\varepsilon = (iu_\varepsilon, \nabla_{A_\varepsilon} u_\varepsilon), \quad \mu_\varepsilon = \text{curl } j_\varepsilon + h_\varepsilon,$$

where $(a, b) = \frac{1}{2}(a\bar{b} + \bar{a}b)$ and $h_\varepsilon = \text{curl } A_\varepsilon$ is the *induced magnetic field*.

We denote by $\text{Lip}_\Omega(A)$ the set of Lipschitz functions on A which are 0 on $\Omega \setminus A$, and let $\|f\|_{\text{Lip}_\Omega(A)^*} = \sup \int \xi df$, the supremum being taken over functions $\xi \in \text{Lip}_\Omega(A)$ such that $|\nabla \xi|_{L^\infty(A)} \leq 1$.

We say a family $\{f_\alpha\}_\alpha$ is subordinate to a cover $\{A_\alpha\}_\alpha$ if $\text{Supp}(f_\alpha) \subset A_\alpha$ for every α .

Despite the slightly confusing notation, the covering A_α will have nothing to do with the magnetic gauge A_ε . Also, the densities f_α and g_α , as well as n_α and ν_α will implicitly depend on ε , and should be really $f_{\varepsilon, \alpha}$ and $g_{\varepsilon, \alpha}$, etc, but for simplicity we do not indicate this dependence.

Theorem 1. *Let $\{\Omega_\varepsilon\}_{\varepsilon > 0}$ be a family of bounded open sets in \mathbb{R}^2 . Assume that $\{(u_\varepsilon, A_\varepsilon)\}_\varepsilon$, where $(u_\varepsilon, A_\varepsilon)$ is defined over Ω_ε , satisfies for some $0 < \beta < 1$ small enough*

$$(1.1) \quad G_\varepsilon(u_\varepsilon, A_\varepsilon) \leq \varepsilon^{-\beta}.$$

Then the following holds, for ε small enough:

1. (Vortices) *There exists a measure ν_ε , depending only on u_ε (and not on A_ε) of the form $2\pi \sum_i d_i \delta_{a_i}$ for some points $a_i \in \Omega_\varepsilon$ and some integers d_i such that, C denoting a generic constant independent of ε ,*

$$(1.2) \quad \|\mu_\varepsilon - \nu_\varepsilon\|_{(C_0^{0,1}(\Omega_\varepsilon))^*} \leq C\sqrt{\varepsilon}G_\varepsilon(u_\varepsilon, A_\varepsilon),$$

and for any measurable set E

$$|\nu_\varepsilon|(E) \leq C \frac{e_\varepsilon(\widehat{E})}{|\log \varepsilon|}.$$

2. (Covering) There exists a cover $\{A_\alpha\}_\alpha$ of Ω_ε by open sets with diameter and overlap number bounded by a universal constant, and measures $\{f_\alpha\}_\alpha, \{\nu_\alpha\}_\alpha$ subordinate to this cover such that, letting $f_\varepsilon := e_\varepsilon - \frac{1}{2}|\log \varepsilon|\nu_\varepsilon$,

$$f_\varepsilon \geq \sum_\alpha f_\alpha, \quad \nu_\varepsilon = \sum_\alpha \nu_\alpha, \quad \nu_{\alpha_1} \perp \nu_{\alpha_2} \text{ for } \alpha_1 \neq \alpha_2.$$

3. (Energy transport) Letting $n_\alpha := \|\nu_\alpha\|/2\pi$, for each α the following holds: If $\text{dist}(A_\alpha, \Omega_\varepsilon^c) > \varepsilon$ there exists a measure $g_\alpha \geq -C$ such that either

$$(1.3) \quad \|f_\alpha - g_\alpha\|_{\text{Lip}_\Omega(A_\alpha)^*} \leq Cn_\alpha(1 + \beta|\log \varepsilon|) \quad \text{and} \quad g_\alpha(A_\alpha) \geq cn_\alpha|\log \varepsilon|,$$

or

$$(1.4) \quad \|f_\alpha - g_\alpha\|_{\text{Lip}_\Omega(A_\alpha)^*} \leq Cn_\alpha(1 + \log n_\alpha) \quad \text{and} \quad g_\alpha(A_\alpha) \geq cn_\alpha^2 - Cn_\alpha,$$

where and $c, C > 0$ are positive universal constants.

If $\text{dist}(A_\alpha, \Omega_\varepsilon^c) \leq \varepsilon$ there exists $g_\alpha \geq 0$ such that for any function ξ

$$(1.5) \quad \int \xi d(f_\alpha - g_\alpha) \leq Cn_\alpha(|\nabla \xi|_{L^\infty(A_\alpha)} + \beta|\log \varepsilon|\|\xi\|_{L^\infty(A_\alpha)}).$$

4. (Properties of g_ε) Letting $g_\varepsilon = f_\varepsilon + \sum_\alpha (g_\alpha - f_\alpha)$ it holds that

$$(1.6) \quad -C \leq g_\varepsilon \leq e_\varepsilon + \frac{1}{2}|\log \varepsilon|(\nu_\varepsilon)_-,$$

and for any measurable set $E \subset \Omega_\varepsilon$,

$$(1.7) \quad (g_\varepsilon)_-(E) \leq C \frac{e_\varepsilon(\widehat{E})}{|\log \varepsilon|}, \quad (g_\varepsilon)_+(E) \leq C e_\varepsilon(\widehat{E}).$$

Moreover, assuming $|u_\varepsilon| \leq 1$ in Ω_ε and that $E + B(0, C) \subset \Omega_\varepsilon$, for some $C > 0$ large enough, then for every $p < 2$,

$$(1.8) \quad \int_E |j_\varepsilon|^p \leq C_p ((g_\varepsilon)_+(E + B(0, C)) + |E|).$$

The third item admits, or rather implies the following form, from which the covering $\{A_\alpha\}_\alpha$ is hidden.

Corollary 1.1. *Under the hypothesis above and using the same notation, for every $0 < \eta \leq 1$ we have if $\varepsilon > 0$ is small enough: First, for every Lipschitz function ξ vanishing on $\partial\Omega_\varepsilon$*

$$(1.9) \quad \int_{\Omega_\varepsilon} \xi d(g_\varepsilon - f_\varepsilon) \leq C \int_{\Omega_\varepsilon} \widehat{\nabla} \xi \left[d|\nu_\varepsilon| + (\beta + \eta) d(g_\varepsilon)_+ + \frac{|\log \eta|^2}{\eta} dx \right] + C\beta \int_{\widehat{\partial\Omega_\varepsilon}} \widehat{\xi} e_\varepsilon.$$

Second, if $d(E, \partial\Omega_\varepsilon) > C$ then

$$(1.10) \quad |\nu_\varepsilon|(E) \leq C \left(\eta(g_\varepsilon)_+(\widehat{E}) + \frac{1}{\eta} |\widehat{E}| + \frac{e_\varepsilon(\widehat{E} \cap \widehat{\partial\Omega_\varepsilon})}{|\log \varepsilon|} \right).$$

The point in introducing the extra parameter η is that we want to be able to use only a small η -fraction of the “remaining” energy g_ε to control the error $f_\varepsilon - g_\varepsilon$ between the original energy and the displaced one. This corollary is obtained by simply summing the relations (1.5) and controlling n_α and $n_\alpha \log n_\alpha$ by a small fraction of n_α^2 through the elementary relations

$$x \log x \leq \eta x^2 + C \frac{\log^2 \eta}{\eta} \quad 2x \leq \eta x^2 + \frac{1}{\eta}$$

and then controlling n_α^2 by $g_\alpha(A_\alpha)$ via (1.3) or (1.4).

Remark 1.1. *If we let $\eta = 1$ and choose E to be at distance at least 1 from $\partial\Omega$ then (1.9) and (1.10) reduce to*

$$(1.11) \quad \int \xi d(f_\varepsilon - g_\varepsilon) \leq C \int_{\Omega_\varepsilon} \widehat{\nabla} \xi [d(g_\varepsilon)_+ + d|\nu_\varepsilon|]$$

and

$$|\nu_\varepsilon|(E) \leq C \left((g_\varepsilon)_+(\widehat{E}) + |\widehat{E}| \right).$$

If one takes $\xi = \chi_R$ to be a positive cut-off function supported in $B(0, R)$ and $\equiv 1$ in $B(0, R - 1)$ then the right-hand side in (1.11) scales like a boundary term i.e. like R as R gets large, while the left-hand side scales like an interior term.

Remark 1.2. *Assume we have proved the above Theorem and Corollary. Then, given $\{(u_\varepsilon, A_\varepsilon)\}_\varepsilon$ and $\{\Omega_\varepsilon\}_\varepsilon$ satisfying the hypothesis, we may consider for some fixed $\sigma > 0$ the rescaled quantities $\tilde{\varepsilon} = \varepsilon/\sigma$, $\tilde{x} = x/\sigma$ and let*

$$\tilde{u}_\varepsilon(\tilde{x}) = u_\varepsilon(x), \quad \tilde{A}_\varepsilon(\tilde{x}) = \sigma A_\varepsilon(x), \quad \tilde{\Omega}_\varepsilon = \Omega_\varepsilon/\sigma.$$

Then, letting $h = \text{curl } A$ and $\tilde{h} = \text{curl } \tilde{A}$, we have

$$e_\varepsilon^\sigma(u, A) := \sigma^2 \left(\frac{1}{2} |\nabla_A u|^2 + \frac{\sigma^2}{2} h^2 + \frac{1}{4\varepsilon^2} (1 - |u|^2)^2 \right) = \frac{1}{2} |\nabla_{\tilde{A}} \tilde{u}|^2 + \frac{1}{2} \tilde{h}^2 + \frac{1}{4\tilde{\varepsilon}^2} (1 - |\tilde{u}|^2)^2.$$

We may then apply the Theorem to the tilded quantities, yielding a measure \tilde{g}_ε . Then if we let $g_\varepsilon(x) = \tilde{g}_\varepsilon(\tilde{x})$, the measure g_ε will satisfy the properties stated in Theorem 1 and Corollary 1.1, with e_ε replaced by e_ε^σ (and with a different C) provided we modify the definition of \widehat{E} to

$$\widehat{E} = \{x \mid \text{dist}(\tilde{x}, \tilde{E}) < 1\} = \{x \mid \text{dist}(x, E) < \sigma\},$$

(note that we can keep the original definition provided $\sigma \leq 1$).

Then we may add to both e_ε and g_ε the quantity $\left(\frac{1}{2} - \frac{\sigma^2}{2}\right) h_\varepsilon^2$ and obtain in this manner a new g_ε satisfying the listed properties and — for the particular choice $\sigma^2 = \frac{1}{2}$ — the lower bound

$$(1.12) \quad g_\varepsilon \geq \frac{h_\varepsilon^2}{4} - C.$$

We will then usually assume when applying Theorem 1 that this lower bound holds as well as the other conclusions of the theorem.

The next result shows how g_ε has the desired coerciveness properties, and behaves like the renormalized energy. Indeed, under the assumption that the family $\{g_\varepsilon\}_\varepsilon$ is bounded on compact sets (recall that the domains become increasingly large as $\varepsilon \rightarrow 0$) we have compactness results for the vorticities and currents, and lower bounds on $\int g_\varepsilon$ (hence $\int f_\varepsilon$ via (1.9)) in terms of the renormalized energy W .

Before stating that result, we introduce some additional notation. We denote by $\{\mathbf{U}_R\}_{R>0}$ a family of sets in \mathbb{R}^2 such that for some constant $C > 0$ independent of R

$$(1.13) \quad \mathbf{U}_R + B(0, 1) \subset \mathbf{U}_{R+C} \text{ and } \mathbf{U}_{R+1} \subset \mathbf{U}_R + B(0, C).$$

For example $\{\mathbf{U}_R\}_{R>0}$ can be the family $\{B_R\}_{R>0}$ of balls centered at 0 of radius R .

Then we use the notation $\chi_{\mathbf{U}_R}$ for cutoff functions satisfying, for some C independent of R ,

$$(1.14) \quad |\nabla \chi_{\mathbf{U}_R}| \leq C \quad \text{Supp}(\chi_{\mathbf{U}_R}) \subset \mathbf{U}_R \quad \chi_{\mathbf{U}_R}(x) = 1 \text{ if } \text{dist}(x, \mathbf{U}_R^c) \geq 1.$$

Finally, given a vector field $j : \mathbb{R}^2 \rightarrow \mathbb{R}^2$ such that $\text{curl } j = 2\pi \sum_{p \in \Lambda} \delta_p + h$ with Λ , where h is in L^2_{loc} and Λ a discrete set, we define the *renormalized energy* of j by

$$W(j) = \limsup_{R \rightarrow \infty} \frac{W(j, \chi_{B_R})}{|B_R|},$$

where for any χ

$$(1.15) \quad W(j, \chi) = \liminf_{\eta \rightarrow 0} \left(\frac{1}{2} \int_{\mathbb{R}^2 \setminus \cup_{p \in \Lambda} B(p, \eta)} \chi |j|^2 + \pi \log \eta \sum_{p \in \Lambda} \chi(p) \right).$$

Various results on W , in particular on its minimizers, are proved in [SS5]. Note in particular that if we assume $\text{div } j = 0$, then the \liminf in (1.15) is in fact a limit, because in this case $j = \nabla^\perp H$ with $\Delta H = 2\pi \delta_p + h$ in a neighbourhood of p , and thus $H = \log |\cdot - p| + f$ with $f \in H^1$ in this neighbourhood.

Theorem 2. *Under the hypothesis of Theorem 1, and assuming $|u_\varepsilon| \leq 1$ in Ω_ε we have the following.*

1. *Assume that $\text{dist}(0, \partial\Omega_\varepsilon) \rightarrow +\infty$ as $\varepsilon \rightarrow 0$ and that, for any $R > 0$,*

$$(1.16) \quad \limsup_{\varepsilon \rightarrow 0} \int g_\varepsilon(\mathbf{U}_R) dx < +\infty,$$

where $\{\mathbf{U}_R\}_R$ satisfies (1.13).

Then, up to extraction of a subsequence, the vorticities $\{\mu_\varepsilon\}_\varepsilon$ converge in $W_{\text{loc}}^{-1,p}(\mathbb{R}^2)$ to a measure ν of the form $2\pi \sum_{p \in \Lambda} \delta_p$, where Λ is a discrete subset of \mathbb{R}^2 , the currents $\{j_\varepsilon\}_\varepsilon$ converge weakly in $L_{\text{loc}}^p(\mathbb{R}^2, \mathbb{R}^2)$ for any $p < 2$ to j , and the induced fields $\{h_\varepsilon\}_\varepsilon$ converge weakly in $L_{\text{loc}}^2(\mathbb{R}^2)$ to h which are such that

$$\text{curl } j = \nu - h \quad \text{in } \mathbb{R}^2.$$

2. *If we replace the assumption (1.16) by the stronger assumption*

$$(1.17) \quad \limsup_{\varepsilon \rightarrow 0} \int g_\varepsilon(\mathbf{U}_R) < CR^2,$$

where C is independent of R , then the limit j of the currents satisfies, for any $p < 2$,

$$(1.18) \quad \limsup_{R \rightarrow +\infty} \int_{\mathbf{U}_R} |j|^p dx < +\infty.$$

Moreover for every family $\chi_{\mathbf{U}_R}$ satisfying (1.14) we have

$$(1.19) \quad \liminf_{\varepsilon \rightarrow 0} \int_{\mathbb{R}^2} \frac{\chi_{\mathbf{U}_R}}{|\mathbf{U}_R|} dg_\varepsilon \geq \left(\frac{W(j, \chi_{\mathbf{U}_R})}{|\mathbf{U}_R|} + \frac{1}{2} \int_{\mathbf{U}_R} h^2 + \frac{\gamma}{2\pi} \int_{\mathbf{U}_R} h \right) + o_R(1),$$

where γ is a constant defined below and $o_R(1)$ is function tending to 0 as $R \rightarrow +\infty$.

Remark 1.3. *The constant γ in (1.19) was introduced in [BBH] and may be defined by*

$$\gamma = \lim_{R \rightarrow \infty} \left(\frac{1}{2} \int_{B_R} |\nabla u_0|^2 + \frac{(1 - |u_0|^2)^2}{2} - \pi \log R \right),$$

where $u_0(r, \theta) = f(r)e^{i\theta}$ is the unique (up to translation and rotation) radially symmetric degree-one vortex (see [BBH, Mi]).

Remark 1.4. *Lower bounds immediately follow from this theorem. Indeed f_ε is the energy density minus the energetic cost of a vortex, and $f_\varepsilon - g_\varepsilon$ is controlled by Theorem 1, see also Remark 1.1. This, combined with the lower bound (1.19) shows that in good cases the averages over large balls of f_ε are bounded below by W plus explicit constants, which proves a sharp lower bound for the energy with a $o(1)$ order error, à la Bethuel-Brezis-Hélein [BBH].*

The bound (1.9) may also be interpreted as a lower bound for the Ginzburg-Landau energy with weight. Assuming a fixed domain Ω and $G_\varepsilon(u_\varepsilon, A_\varepsilon) < C|\log \varepsilon|$ for instance, and that $\mu_\varepsilon \rightarrow 2\pi \sum_{i=1}^n \delta_{a_i}$, where $a_i \in \Omega$, then by blowing up by a factor independent of ε we may assume the points are at distance 2, say, from the boundary and then if ξ is a fixed positive weight we may multiply it by a cutoff $0 \leq \chi \leq 1$ equal to zero on $\widehat{\partial\Omega}$ and equal to 1 at each a_i . Then (1.9) becomes

$$\int_{\Omega} \xi e_\varepsilon \geq \pi |\log \varepsilon| \sum_{i=1}^n \xi(a_i) + \int \chi \xi dg_\varepsilon - C \int \widehat{\nabla(\chi\xi)} \left[d|\nu_\varepsilon| + (\beta + \eta) d(g_\varepsilon)_+ + \frac{|\log \eta|^2}{\eta} dx \right].$$

Typically, there will be an upper bound for the energy which implies that $(g_\varepsilon)_+(\Omega) < C$ and since also $g_\varepsilon \geq -C$, the integrals on the right-hand side may be bounded below by a constant independent of ε .

The paper is organized as follows: In Section 2 we state without proof the result on lower bounds via Jerrard's ball construction (the proof is postponed to Section 5) which we adapt for our purposes, and explain how we use it on a covering of Ω_ε by a collection U_α of balls of finite size. In Section 3, we present the tool used to transport the negative part of f_ε to absorb it into the positive part, and deduce Theorem 1. In Section 4, we prove Theorem 2. Finally in Section 5, we prove the ball-construction lower bound.

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2 Use of the ball construction and coverings of the domain

The first step consists in performing a ball construction in Ω_ε in order to obtain lower bounds. This follows essentially the method of Jerrard [Je], the difficulty being that we are not allowed more than an error of order one per vortex. This is hopeless if the total number of vortices diverges when $\varepsilon \rightarrow 0$, hence we need to localize the construction in pieces of Ω_ε small enough for the number of vortices in each piece to remain bounded as $\varepsilon \rightarrow 0$.

2.1 The ball construction lower bound

We start by stating the result of Jerrard's ball construction in a version adapted to our situation, in particular including the magnetic field. The proof is postponed to Section 5. In all what follows, if \mathcal{B} is a collection of balls, $r(\mathcal{B})$ denotes the sum of the radii of the balls in the collection. In all the sequel we will sometimes abuse notation by writing \mathcal{B} for $\cup_{B \in \mathcal{B}} B$, i.e. identify the collection of balls and the set it covers.

Proposition 2.1. *There exists $\varepsilon_0, C > 0$ such that if $U \subset \mathbb{R}^2$, $\varepsilon \in (0, \varepsilon_0)$, and $(u_\varepsilon, A_\varepsilon)$ defined on U is such that*

$$G_\varepsilon(u_\varepsilon, A_\varepsilon) \leq \varepsilon^{-\beta},$$

where $\beta \in (0, 1)$, the following holds.

For every $r \in (C\varepsilon^{1-\beta}, \frac{1}{2})$, there exists a collection of disjoint closed balls \mathcal{B} depending only on u_ε (and not on A_ε) such that, letting $U_\varepsilon = \{x \mid d(x, U^c) > \varepsilon\}$,

1. $\{x \in U_\varepsilon \mid |u_\varepsilon(x)| < \frac{1}{2}\} \subset \mathcal{B}$.

2. $r(\mathcal{B}) \leq r$.

3. For any $2 \leq \bar{C} \leq (r/\varepsilon)^{\frac{1}{2}}$ it holds that either

$$e_\varepsilon(\mathcal{B} \cap U) \geq \bar{C} \log \frac{r}{\varepsilon}$$

or

$$\forall B \in \mathcal{B} \text{ such that } B \subset U_\varepsilon, \quad e_\varepsilon(B) \geq \pi |d_B| \left(\log \frac{r}{\varepsilon \bar{C}} - C \right),$$

where $d_B = \deg(u_\varepsilon, \partial B)$.

A natural choice of \bar{C} above is πD , where $D = \sum_{B \in \mathcal{B}} |d_B|$ and we have let $d_B = 0$ if $B \not\subset U_\varepsilon$. With this choice we find in all cases

$$e_\varepsilon(\mathcal{B} \cap U) \geq \pi D \left(\log \frac{r}{\varepsilon D} - C \right)$$

i.e. we recover the same lower bound as in [SS4], Theorem 4.1, mentioned in the introduction as (0.1). The reason why we don't simply use that theorem directly is that we need to keep the dichotomy above, and thus a lower bound localized in each ball.

2.2 Localizing the ball construction

For any $\varepsilon > 0$ we construct an open cover $\{U_\alpha\}_\alpha$ of Ω_ε as follows: We consider the collection \mathcal{B} of balls of radius ℓ_0 — where $\ell_0 \in (0, \frac{1}{8})$ is to be chosen below, small enough but independent of ε — centered at the points of $\ell_0 \mathbb{Z}^2$. The cover consists of the open sets $\Omega_\varepsilon \cap B$, for $B \in \mathcal{B}$.

This cover depends on ε , but the maximal number of *neighbours* of a given α — defined as the indices β such that $U_\alpha \cap U_\beta \neq \emptyset$ — is bounded independently of ε by an integer we denote by m (in fact $m = 9$). Note that m also bounds the *overlap number* of the cover, i.e. the maximal number of U_α 's to which a given x can belong. There is also $\ell > 0$ independent of ε which is a Lebesgue number of the cover, i.e. such that for every $x \in \Omega_\varepsilon$, there exists α such that $B(x, \ell) \cap \Omega_\varepsilon \subset U_\alpha$ or, equivalently, $\text{dist}(x, \Omega_\varepsilon \cap U_\alpha^c) \geq \ell$.

Assuming $\beta < 1/4$, and applying Proposition 2.1 to $(u_\varepsilon, A_\varepsilon)$ in U_α for every α we obtain, since $\sqrt{\varepsilon} > C\varepsilon^{1-\beta}$ if ε is small enough, a collection $\mathcal{B}_\varepsilon^{\alpha, r}$ for every $\sqrt{\varepsilon} \leq r \leq 1/2$.

If ρ is chosen small enough depending on ℓ and m only, thus less than a universal constant, we may extract from $\cup_\alpha \mathcal{B}_\varepsilon^{\alpha, \rho}$ a subcollection \mathcal{B}_ε such that any two balls B, B' in \mathcal{B}_ε satisfy $\Omega_\varepsilon \cap B \cap B' = \emptyset$. We will say \mathcal{B}_ε is *disjoint in Ω_ε* :

Proposition 2.2. *Assume $\rho \leq \ell/(8m)$. Then, writing in short $\mathcal{B}_\varepsilon^\alpha$ instead of $\mathcal{B}_\varepsilon^{\alpha,\rho}$, there exists a subcollection of $\cup_\alpha \mathcal{B}_\varepsilon^\alpha$ — call it \mathcal{B}_ε — which is disjoint in Ω_ε and such that*

$$(2.1) \quad \{|u_\varepsilon| \leq 1/2\} \cap \{x \mid \text{dist}(x, \Omega_\varepsilon^c) > \varepsilon\} \subset \cup_{B \in \mathcal{B}_\varepsilon} B.$$

Moreover, for every $B \in \mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha$ we have $B \cap \Omega_\varepsilon = B \cap U_\alpha$ and

$$\text{dist}(B, \Omega_\varepsilon^c) > \varepsilon \iff \text{dist}(B, U_\alpha^c) > \varepsilon.$$

Proof. Assume $C = \Omega_\varepsilon \cap (B_1 \cup \dots \cup B_k)$ is a connected component of $\Omega_\varepsilon \cap (\cup_\alpha \mathcal{B}_\varepsilon^\alpha)$. Reordering if necessary, we may assume that $B_i \cap (B_1 \cup \dots \cup B_{i-1}) \neq \emptyset$ for every $1 \leq i \leq k$. There exists $x \in \Omega_\varepsilon \cap B_1$ and α such that $\text{dist}(x, \Omega_\varepsilon \cap U_\alpha^c) \geq \ell$. Then $\text{dist}(B_1, \Omega_\varepsilon \cap U_\alpha^c) > 3\ell/4$. Assume

$$\text{dist}(B_1 \cup \dots \cup B_{i-1}, \Omega_\varepsilon \cap U_\alpha^c) \geq \frac{3\ell}{4}.$$

Then $\text{dist}(B_i, \Omega_\varepsilon \cap U_\alpha^c) > \ell/2$ hence for every $1 \leq j \leq i$ the ball B_j belongs to $\mathcal{B}_\varepsilon^\beta$, where β is a neighbour of α . It follows that $r_1 + \dots + r_i \leq m\rho \leq \ell/8$, where r_i is the radius of B_i , and we deduce that $B_1 \cup \dots \cup B_i \subset B(x, \ell/4)$ and then

$$\text{dist}(B_1 \cup \dots \cup B_i, \Omega_\varepsilon \cap U_\alpha^c) \geq \frac{3\ell}{4}.$$

We have thus proved by induction that $C \subset U_\alpha$ and even that $\text{dist}(C, \Omega_\varepsilon \cap U_\alpha^c) \geq 3\ell/4$ for every i .

We delete from $\{B_1, \dots, B_k\}$ the balls which do not belong to $\mathcal{B}_\varepsilon^\alpha$ and call C' the union of the remaining balls. If y belongs to

$$C \cap \{|u_\varepsilon| \leq 1/2\} \cap \{x \mid \text{dist}(x, \Omega_\varepsilon^c) > \varepsilon\}$$

then, since $\text{dist}(C, \Omega_\varepsilon \cap U_\alpha^c) \geq 3\ell/4$ and $\text{dist}(y, \Omega_\varepsilon^c) > \varepsilon$, provided $\varepsilon < 3\ell/4$ we have that $\text{dist}(y, U_\alpha^c) > \varepsilon$ hence y belongs to some ball $B \in \mathcal{B}_\varepsilon^\alpha$ (since $\mathcal{B}_\varepsilon^\alpha$ covers the set $\{|u_\varepsilon| \leq \frac{1}{2}\} \cap \{\text{dist}(x, U_\alpha^c) > \varepsilon\}$), thus $y \in C'$. The balls in C' are disjoint in Ω_ε since they belong to the collection $\mathcal{B}_\varepsilon^\alpha$ which is itself disjoint in Ω_ε .

Performing this operation on each connected component of $\Omega_\varepsilon \cap (\cup_\alpha \mathcal{B}_\varepsilon^\alpha)$ we thus obtain a collection \mathcal{B}_ε which covers $\{|u_\varepsilon| \leq 1/2\} \cap \{x \mid \text{dist}(x, \Omega_\varepsilon^c) > \varepsilon\}$ and is disjoint in Ω_ε . Moreover, if $B \in \mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha$ then $\text{dist}(B, \Omega_\varepsilon \cap U_\alpha^c) \geq 3\ell/4$ hence $B \cap \Omega_\varepsilon = B \cap U_\alpha$ and

$$\text{dist}(B, \Omega_\varepsilon^c) > \varepsilon \iff \text{dist}(B, U_\alpha^c) > \varepsilon.$$

□

The value ρ will be fixed to some value smaller than $\ell/8m$ and independent of ε , to be specified below. The above proposition provides us for any $\varepsilon > 0$ small enough with collections of balls \mathcal{B}_ε and $\mathcal{B}_\varepsilon^\alpha$. We will also need the following

Definition 1. For any $\sqrt{\varepsilon} \leq r \leq \rho$, and any $B \in \mathcal{B}_\varepsilon^\alpha$, we let $\mathcal{B}_\varepsilon^{B,r}$ be the collection of balls in $\mathcal{B}_\varepsilon^{\alpha,r}$ which are included in B . Then we let

$$\mathcal{B}_\varepsilon^r = \cup_{B \in \mathcal{B}_\varepsilon} \mathcal{B}_\varepsilon^{B,r}.$$

It is disjoint in Ω_ε and covers the set $\{|u_\varepsilon| \leq 1/2\} \cap \{x \mid \text{dist}(x, \Omega_\varepsilon^c) > \varepsilon\}$ and of course if $B \in \mathcal{B}_\varepsilon^r \cap \mathcal{B}_\varepsilon^{\alpha,r}$, then $B \cap \Omega_\varepsilon = B \cap U_\alpha$ and

$$\text{dist}(B, \Omega_\varepsilon^c) > \varepsilon \iff \text{dist}(B, U_\alpha^c) > \varepsilon.$$

In other words, the disjoint collection \mathcal{B}_ε permits us to construct disjoint collections of smaller radius by discarding from $\mathcal{B}_\varepsilon^{\alpha,r}$ those balls which are inside a ball discarded from $\mathcal{B}_\varepsilon^{\alpha,\rho}$. The collection $\mathcal{B}_\varepsilon^{\sqrt{\varepsilon}}$ should be seen as the collection of “small balls” and \mathcal{B}_ε (obtained from $\mathcal{B}_\varepsilon^{\alpha,\rho}$) as the collection of “large balls”. We will sometimes also use the collection of the intermediate size balls $\mathcal{B}_\varepsilon^r$ with $\sqrt{\varepsilon} \leq r \leq \rho$.

Finally we let

$$(2.2) \quad \nu_\varepsilon = \sum_{\substack{B \in \mathcal{B}_\varepsilon^{\sqrt{\varepsilon}} \\ \text{dist}(B, \Omega_\varepsilon^c) > \varepsilon}} 2\pi d_B \delta_{a_B}, \quad |\nu_\varepsilon| = \sum_{\substack{B \in \mathcal{B}_\varepsilon^{\sqrt{\varepsilon}} \\ \text{dist}(B, \Omega_\varepsilon^c) > \varepsilon}} 2\pi |d_B| \delta_{a_B},$$

where a_B is the center of B , and d_B denotes the winding number of $u_\varepsilon/|u_\varepsilon|$ restricted to ∂B . This is the ν_ε given by the conclusion of the theorem. Note that since the balls only depend on u_ε (and not on A_ε), ν_ε satisfies the same. If B is any ball which does not cross the boundary of balls in $\mathcal{B}_\varepsilon^{\sqrt{\varepsilon}}$ and $\text{dist}(B, \Omega_\varepsilon^c) > \varepsilon$ then $\nu_\varepsilon(B) = 2\pi d_B$. From the Jacobian estimate (see [JS] or the version in [SS4], Theorem 6.1) we have that (1.2) is satisfied. We also have (recall that $\|f\|$ denotes the total variation of a measure)

Lemma 2.1. *There exists $\varepsilon_0 > 0$ such that if $\beta < 1/4$ in (1.1) and $\varepsilon < \varepsilon_0$ then*

$$|\nu_\varepsilon|(E) \leq 16 \frac{e_\varepsilon(\Omega_\varepsilon \cap \widehat{E})}{|\log \varepsilon|}$$

for any measurable set E , so that choosing $E = \Omega_\varepsilon$ and taking logarithms,

$$(2.3) \quad \log \|\nu_\varepsilon\| \leq \beta |\log \varepsilon| + C.$$

Proof. We use the properties of $\mathcal{B}_\varepsilon^{\alpha,\sqrt{\varepsilon}}$. Letting $\overline{C} = (\sqrt{\varepsilon}/\varepsilon)^{\frac{1}{2}} = \varepsilon^{-\frac{1}{4}}$, it is impossible when ε is small enough that $e_\varepsilon(\Omega_\varepsilon \cap \mathcal{B}_\varepsilon^{\alpha,\sqrt{\varepsilon}}) \geq \overline{C} \log(\sqrt{\varepsilon}/\varepsilon)$ since we assumed that $e_\varepsilon(\Omega_\varepsilon) \leq \varepsilon^{-\beta}$. Thus Proposition 2.1 implies that, for every $B \in \mathcal{B}_\varepsilon^{\alpha,\sqrt{\varepsilon}}$ such that $\text{dist}(B, U_\alpha^c) > \varepsilon$,

$$e_\varepsilon(B) \geq \pi |d_B| \left(\log \varepsilon^{-\frac{1}{4}} - C \right) \geq \frac{\pi}{8} |d_B| |\log \varepsilon|,$$

if ε is small enough. If, moreover, $B \in \mathcal{B}_\varepsilon^{\sqrt{\varepsilon}}$, then from Definition 1 we have

$$\text{dist}(B, U_\alpha^c) > \varepsilon \iff \text{dist}(B, \Omega_\varepsilon^c) > \varepsilon$$

Hence for any set E , using (2.2) and the fact that balls in $\mathcal{B}_\varepsilon^{\sqrt{\varepsilon}}$ have radius smaller than $1/2$ if ε is small enough,

$$|\nu_\varepsilon|(E) \leq \sum_{\substack{B \in \mathcal{B}_\varepsilon^{\sqrt{\varepsilon}} \\ \text{dist}(B, \Omega_\varepsilon^c) > \varepsilon \\ B \cap E \neq \emptyset}} |\nu_\varepsilon|(B) \leq 16 \frac{e_\varepsilon(\Omega_\varepsilon \cap \widehat{E})}{|\log \varepsilon|}.$$

□

Definition 2. For any α we let ν_α denote the restriction of ν_ε to the balls in $\mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha$ and $n_\alpha = \frac{\|\nu_\alpha\|}{2\pi}$, so that

$$\nu_\varepsilon = \sum_\alpha \nu_\alpha, \quad n_\alpha = \sum_{B \in \mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha} \frac{|\nu_\varepsilon|(B)}{2\pi}, \quad \|\nu_\varepsilon\| = 2\pi \sum_\alpha n_\alpha.$$

We also define

$$(2.4) \quad \bar{C}_\alpha = \begin{cases} \max\left(Mn_\alpha, \frac{3e_\alpha}{|\log \varepsilon|}\right) & \text{if } n_\alpha \neq 0, \\ 2 & \text{otherwise,} \end{cases}$$

where M is a large universal constant to be chosen later and

$$e_\alpha = \sum_{B \in \mathcal{B}_\varepsilon^\alpha} e_\varepsilon(B \cap U_\alpha).$$

Note that n_α is the sum of the absolute values of the degrees of the small balls included in the large balls of $\mathcal{B}_\varepsilon^\alpha$.

We have the following

Proposition 2.3. *There exists $\varepsilon_0, C_0 > 0$ such that if $\beta < 1/4$ in (1.1) and $\varepsilon < \varepsilon_0$, $\varepsilon^{\frac{1}{2}} < r < \rho$ then $2 \leq \bar{C}_\alpha \leq (r/\varepsilon)^{\frac{1}{2}}$ and for any $B \in \mathcal{B}_\varepsilon^r \cap \mathcal{B}_\varepsilon^{\alpha,r}$ such that $\text{dist}(B, \Omega_\varepsilon^c) > \varepsilon$ we have*

$$(2.5) \quad e_\varepsilon(B) \geq 2\pi |d_B| \Lambda_\varepsilon^{\alpha,r}, \quad \text{where} \quad \Lambda_\varepsilon^{\alpha,r} = \frac{1}{2} \left(\log \frac{r}{\varepsilon \bar{C}_\alpha} - C_0 \right).$$

Moreover, $0 \leq \Lambda_\varepsilon^{\alpha,r} \leq \frac{1}{2} |\log \varepsilon|$ and

$$(2.6) \quad 0 \leq \frac{1}{2} |\log \varepsilon| - \Lambda_\varepsilon^{\alpha,r} \leq \frac{1}{2} (\beta |\log \varepsilon| + |\log r| + C_0).$$

Proof. From the definition (2.4), from (1.1) and Lemma 2.1 we have for ε small enough that $2 \leq \bar{C}_\alpha \leq \varepsilon^{-\beta}$. It follows that if $\varepsilon^{\frac{1}{2}} < r < 1$ then $2 \leq \bar{C}_\alpha \leq (r/\varepsilon)^{\frac{1}{2}}$, since $\beta < 1/4$. Also, from the definition of \bar{C}_α it is impossible that $e_\varepsilon(\mathcal{B}_\varepsilon^{\alpha,r} \cap U_\alpha) \geq \bar{C}_\alpha \log(r/\varepsilon)$ since for $\sqrt{\varepsilon} \leq r \leq \rho$ we have $\bar{C}_\alpha \geq 3e_\varepsilon(\mathcal{B}_\varepsilon^{\alpha,r})/|\log \varepsilon|$.

Then from Proposition 2.1, letting $\overline{C} = \overline{C}_\alpha$, we deduce (2.5) for any $B \in \mathcal{B}_\varepsilon^{\alpha,r}$ such that $\text{dist}(B, U_\alpha^c) > \varepsilon$, which is equivalent to $\text{dist}(B, \Omega_\varepsilon^c) > \varepsilon$ if $B \in \mathcal{B}_\varepsilon^r \cap \mathcal{B}_\varepsilon^{\alpha,r}$.

Finally, $r/(\varepsilon\overline{C}_\alpha) \geq \varepsilon^{-\frac{1}{4}}$ using $\overline{C}_\alpha \leq (r/\varepsilon)^{\frac{1}{2}}$ and $r \geq \sqrt{\varepsilon}$, which easily implies that $\Lambda_\varepsilon^{\alpha,r} > 0$ if ε is small enough, and $\Lambda_\varepsilon^{\alpha,r} \leq \frac{1}{2}|\log \varepsilon|$ is clear from the definition. Inequality (2.6) follows from

$$\frac{1}{2}|\log \varepsilon| - \Lambda_\varepsilon^{\alpha,r} = \frac{1}{2} \left(\log \frac{\overline{C}_\alpha}{r} + C_0 \right)$$

since $\overline{C}_\alpha \leq \varepsilon^{-\beta}$. □

3 Mass Transport

We proceed to the displacement of the negative part of

$$f_\varepsilon = e_\varepsilon - \frac{1}{2}|\log \varepsilon|\nu_\varepsilon.$$

3.1 Mass transport abstract lemmas

For the displacements we will use the following two lemmas. The first, more sophisticated one, was already stated in the introduction and uses optimal transportation for the 1-Wasserstein distance (or minimal connection cost).

Lemma 3.1. *Assume f is a finite Radon measure on a compact set A , that Ω is open and that for any positive Lipschitz function ξ in $\text{Lip}_\Omega(A)$, i.e. vanishing on $\Omega \setminus A$,*

$$\int \xi df \geq -C_0|\nabla \xi|_{L^\infty(A)}.$$

Then there exists a Radon measure g on A such that $0 \leq g \leq f_+$ and such that

$$\|f - g\|_{\text{Lip}_\Omega(A)^*} \leq C_0.$$

Proof. The proof uses convex analysis. Let $X = C(A)$ denotes the space of continuous functions and for $\xi \in X$ let

$$\varphi(\xi) = \int \xi_+ df_+ \quad \text{and} \quad \psi(\xi) = \begin{cases} +\infty & \text{if } |\nabla \xi|_{L^\infty(A)} > 1 \text{ or } \xi \notin \text{Lip}_\Omega(A) \\ -\int \xi df & \text{otherwise} \end{cases}.$$

Then ψ is lower semicontinuous because $\{\xi \in \text{Lip}_\Omega(A) \mid |\nabla \xi|_{L^\infty} \leq 1\}$ is closed under uniform convergence, and φ is continuous. Moreover both functions are convex, and finite for $\xi = 0$. Then the theorem of Fenchel-Rockafellar (see for instance [ET]) yields

$$\inf_X (\varphi + \psi) = \max_{\mu \in X^*} (-\varphi^*(-\mu) - \psi^*(\mu)),$$

where X^* is the dual of X , i.e. the Radon measures on A and

$$\varphi^*(\mu) = \sup_{\xi \in X} \int \xi d\mu - \int \xi_+ df_+ = \begin{cases} 0 & \text{if } 0 \leq \mu \leq f_+, \\ +\infty & \text{otherwise} \end{cases},$$

$$\psi^*(\mu) = \sup_{\substack{\xi \in \text{Lip}_\Omega \\ |\nabla \xi|_\infty \leq 1}} \int \xi d\mu + \int \xi df = \|\mu + f\|_{\text{Lip}_\Omega^*}.$$

We deduce that

$$\inf_{\substack{\xi \in \text{Lip}_\Omega \\ |\nabla \xi|_{L^\infty} \leq 1}} \int \xi_+ df_+ - \int \xi df = \max_{0 \leq -\mu \leq f_+} (-\|\mu + f\|_{\text{Lip}_\Omega^*})$$

and then the existence of a Radon measure g such that $-g$ maximizes the right-hand side, ie such that $0 \leq g \leq f_+$ and

$$-\|f - g\|_{\text{Lip}_\Omega^*} = \inf_{\substack{\xi \in \text{Lip}_\Omega \\ |\nabla \xi|_{L^\infty} \leq 1}} \int \xi_+ df_+ - \int \xi df.$$

But

$$\begin{aligned} \inf_{\substack{\xi \in \text{Lip}_\Omega \\ |\nabla \xi|_{L^\infty} \leq 1}} \int \xi_+ df_+ - \int \xi df &= - \sup_{\substack{\xi \in \text{Lip}_\Omega \\ |\nabla \xi|_{L^\infty} \leq 1}} \left(\int \xi df - \int \xi_+ df_+ \right) \\ &= - \sup_{\substack{\xi \in \text{Lip}_\Omega \\ |\nabla \xi|_{L^\infty} \leq 1}} \left(\int \xi_+ d(f - f_+) - \int \xi_- df \right) \\ &= - \sup_{\substack{\xi \in \text{Lip}_\Omega \\ |\nabla \xi|_{L^\infty} \leq 1}} \left(- \int \xi_- df \right) = \inf_{\substack{\xi \in \text{Lip}_\Omega \\ |\nabla \xi|_{L^\infty} \leq 1}} \int \xi_- df. \end{aligned}$$

The assumption of the lemma implies that this last right-hand side is $\geq -C_0$ therefore

$$\|f - g\|_{\text{Lip}_\Omega(A)^*} \leq C_0.$$

□

The second, less sophisticated, displacement result is

Lemma 3.2. *Assume f is a finite Radon measure supported in Ω and such that $f(\Omega) \geq 0$. Then there exists $0 \leq g \leq f_+$ such that for any Lipschitz function ξ*

$$\int_{\Omega} \xi d(f - g) \leq 2 \text{diam}(\Omega) |\nabla \xi|_{L^\infty(\Omega)} f_-(\Omega).$$

Proof. This follows from the previous Lemma but can be proved directly by letting (assuming $f \neq 0$, otherwise $g = 0$ is the answer),

$$g = f_+ \left(1 - \frac{f_-(\Omega)}{f_+(\Omega)} \right).$$

Then g is positive because $f(\Omega) \geq 0$ implies $f_-(\Omega) \leq f_+(\Omega)$ and

$$\int \xi d(f - g) = \int \xi d \left(f_+ \frac{f_-(\Omega)}{f_+(\Omega)} - f_- \right) = \int (\xi - \bar{\xi}) d \left(f_+ \frac{f_-(\Omega)}{f_+(\Omega)} - f_- \right),$$

where $\bar{\xi}$ is the average of ξ over Ω , and the right-hand side is clearly bounded above by $2 \operatorname{diam}(\Omega) |\nabla \xi|_\infty f_-(\Omega)$. \square

3.2 Mass displacement in the balls

Definition 3. For $B \in \mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha$. We let

$$f_\varepsilon^B = (e_\varepsilon - \Lambda_\varepsilon^\alpha \nu_\varepsilon) \mathbf{1}_{B \cap \Omega_\varepsilon}.$$

where $\Lambda_\varepsilon^{\alpha,r}$ is defined in (2.5) and we have set $\Lambda_\varepsilon^\alpha = \Lambda_\varepsilon^{\alpha,\rho}$.

This corresponds to the excess energy in the balls i.e. the energy remaining after subtracting off the expected value from the ball construction. There is a difference of order $|\nu_\varepsilon|(B) \log \bar{C}_\alpha$ between $f_\varepsilon(B)$ and $f_\varepsilon^B(B)$ which will be dealt with later.

Proposition 3.1. *There exists $\varepsilon_0, C > 0$ such that for any $\varepsilon < \varepsilon_0$, and any $B \in \mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha$, there exists a positive measure g_ε^B defined in $B \cap \Omega_\varepsilon$ and such that*

$$(3.1) \quad g_\varepsilon^B \leq e_\varepsilon + \Lambda_\varepsilon^\alpha (\nu_\varepsilon)_- \quad \text{and} \quad \int_{B \cap \Omega_\varepsilon} \xi d(f_\varepsilon^B - g_\varepsilon^B) \leq C |\nabla \xi|_{L^\infty(B \cap \Omega_\varepsilon)} |\nu_\varepsilon|(B),$$

for any Lipschitz function ξ vanishing on $\Omega_\varepsilon \setminus B$.

Proof. To prove the existence of g_ε^B , in view of Lemma 3.1 and since $(f_\varepsilon^B)_+ = e_\varepsilon + \Lambda_\varepsilon^\alpha (\nu_\varepsilon)_-$ on B it suffices to prove that for any positive function ξ defined on B and vanishing on $B \setminus \Omega_\varepsilon$ we have

$$(3.2) \quad \int \xi d f_\varepsilon^B \geq -C |\nabla \xi|_{L^\infty(B)} |\nu_\varepsilon|(B).$$

We turn to the proof of (3.2). Let $B \in \mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha$ and ξ be as above. Then

$$(3.3) \quad \int \xi d f_\varepsilon^B = \int_0^{+\infty} f_\varepsilon^B(E_t \cap B) dt,$$

where we have set $E_t = \{x \in B \mid \xi(x) \geq t\}$ and $f_\varepsilon^B(A) = \int_A f_\varepsilon^B$.

We will divide the integral (3.3) into $\int_0^{t_\varepsilon} + \int_{t_\varepsilon}^{+\infty}$, with $t_\varepsilon = \varepsilon|\nabla\xi|_{L^\infty}$. The first integral is straightforward to bound from below. Indeed $(f_\varepsilon^B)_-(B) \leq C|\log \varepsilon||\nu_\varepsilon|(B)$ hence

$$(3.4) \quad \int_0^{t_\varepsilon} f_\varepsilon^B(E_t) dt \geq -C\varepsilon|\log \varepsilon||\nabla\xi|_{L^\infty}|\nu_\varepsilon|(B) \geq -C|\nabla\xi|_{L^\infty}|\nu_\varepsilon|(B).$$

On the other hand, if $t > t_\varepsilon$, and this motivated our choice of t_ε , then since $\xi = 0$ in $B \setminus \Omega_\varepsilon$ we have $\text{dist}(E_t, \Omega_\varepsilon^c) > \varepsilon$. Let then $t > t_\varepsilon$, and $a \in E_t$ be a point in the support of ν_ε . Then for any $r \in [\sqrt{\varepsilon}, \rho]$, there exists a ball $B_{a,r} \in \mathcal{B}_\varepsilon^r$ containing a . Since $\{\mathcal{B}_\varepsilon^r\}$ is monotonic with respect to r , $B_{a,r} \subset B$. We call

$$r(a, t) = \sup\{r \in [\sqrt{\varepsilon}, \rho], B_{a,r} \subset E_t\}$$

if this set is nonempty, and 0 otherwise. We then let

$$B_a^t = B_{a,r(a,t)}.$$

If $0 < r(a, t) < \rho$ then $r(a, t)$ bounds from above the distance of a to the complement of E_t . In particular

$$(3.5) \quad \xi(a) - t \leq r(a, t)|\nabla\xi|_{L^\infty}.$$

Indeed for any $r(a, t) < s < \rho$ we have $B_{a,s} \subset B$ and $B_{a,s} \cap (E_t)^c \neq \emptyset$ hence there exists $b \in B_{a,s} \cap \partial E_t$. Then $\xi(a) - \xi(b) \leq s|\nabla\xi|_{L^\infty}$ and since $\partial E_t \subset \{\xi = t\}$ we deduce $\xi(a) - t \leq s|\nabla\xi|_{L^\infty}$, proving (3.5) by making s tend to $r(a, t)$ from above.

A second fact is that if $r(a, t) = 0$, then $\overline{B_{a,\sqrt{\varepsilon}}}$ intersects $B \setminus E_t$ and as above we deduce

$$(3.6) \quad \xi(a) - t \leq \sqrt{\varepsilon}|\nabla\xi|_{L^\infty(B)}.$$

The third fact is that the collection $\{B_a^t\}_a$, where a ranges over E_t and the a 's for which $r(a, t) = 0$ have been excluded, is disjoint. Indeed take $a, b \in E_t$ and assume that $r(a, t) \geq r(b, t)$. Then, since $\mathcal{B}_{r(a,t)}$ is disjoint, the balls $B_{a,r(a,t)}$ and $B_{b,r(a,t)}$ are either equal or disjoint. If they are disjoint we note that $r(a, t) \geq r(b, t)$ implies that $B_{b,r(b,t)} \subset B_{b,r(a,t)}$ and therefore $B_b^t = B_{b,r(b,t)}$ and $B_a^t = B_{a,r(a,t)}$ are disjoint. If they are equal, then $B_{b,r(a,t)} \subset E_t$ and therefore $r(b, t) \geq r(a, t)$, which implies $r(b, t) = r(a, t)$ and then $B_b^t = B_a^t$.

Now, for any $B' \in \{B_a^t\}_a$ we have $B' \subset E_t$ and $\text{dist}(E_t, \Omega_\varepsilon^c) > \varepsilon$ hence $\text{dist}(B', \Omega_\varepsilon^c) > \varepsilon$ and from Proposition 2.3, we have, since $\Lambda_\varepsilon^{\alpha,r} = \Lambda_\varepsilon^{\alpha,\rho} - \frac{1}{2} \log \frac{\rho}{r}$,

$$e_\varepsilon(B') \geq |\nu_\varepsilon(B')| \left(\Lambda_\varepsilon^\alpha - \frac{1}{2} \log \frac{\rho}{r} \right)_+,$$

where r is the common value of $r(a, t)$ for a 's in B' which are in the support of ν_ε . We may rewrite the above as

$$e_\varepsilon(B') \geq \left| \sum_{a \in B' \cap \text{Supp } \nu_\varepsilon} \nu_\varepsilon(a) \left(\Lambda_\varepsilon^\alpha - \frac{1}{2} \log \frac{\rho}{r(a,t)} \right)_+ \right|,$$

and summing over $B' \in \{B_a^t\}_a$ we deduce

$$e_\varepsilon(E_t \cap B) \geq \left| \sum_{a \in \mathcal{P}_t} \left(\Lambda_\varepsilon^\alpha - \frac{1}{2} \log \frac{\rho}{r(a,t)} \right)_+ \nu_\varepsilon(a) \right|,$$

where \mathcal{P}_t is the set of points in $E_t \cap \text{Supp } \nu_\varepsilon$ such that $r(a,t) > 0$. We will let \mathcal{Q}_t be the set of points in $E_t \cap \text{Supp } \nu_\varepsilon$ such that $r(a,t) = 0$.

Since $\nu_\varepsilon(E_t) = \nu_\varepsilon(\mathcal{P}_t) + \nu_\varepsilon(\mathcal{Q}_t)$, subtracting from the above $\Lambda_\varepsilon^\alpha \nu_\varepsilon(E_t)$ we find

$$f_\varepsilon^B(E_t) \geq - \sum_{a \in \mathcal{Q}_t} |\nu_\varepsilon|(a) \Lambda_\varepsilon^\alpha - \frac{1}{2} \sum_{a \in \mathcal{P}_t} |\nu_\varepsilon|(a) \log \frac{\rho}{r(a,t)}.$$

From (3.6), a given $a \in \text{Supp } \nu_\varepsilon \cap B$ can belong to \mathcal{Q}_t only if $|t - \xi(a)| \leq \sqrt{\varepsilon} |\nabla \xi|_{L^\infty}$. Therefore integrating the above with respect to t yields, using the fact that $t \leq \xi(a)$ if $a \in E_t$, that

$$\int_{t_\varepsilon}^\infty f_\varepsilon^B(E_t) dt \geq - \sum_{a \in \text{Supp } \nu_\varepsilon \cap B} |\nu_\varepsilon|(a) \left(\int_{\xi(a) - \sqrt{\varepsilon} |\nabla \xi|_{L^\infty}}^{\xi(a) + \sqrt{\varepsilon} |\nabla \xi|_{L^\infty}} \Lambda_\varepsilon^\alpha dt + \frac{1}{2} \int_0^{\xi(a)} \left(\log \frac{\rho}{r(a,t)} \right)_+ dt \right)$$

hence

$$\int_{t_\varepsilon}^\infty f_\varepsilon^B(E_t) dt \geq -2\Lambda_\varepsilon^\alpha \sqrt{\varepsilon} |\nabla \xi|_{L^\infty} |\nu_\varepsilon|(B) - \frac{1}{2} \sum_{a \in \text{Supp } \nu_\varepsilon \cap B} |\nu_\varepsilon|(a) \int_0^{\xi(a)} \left(\log \frac{\rho}{r(a,t)} \right)_+ dt.$$

We now note that — since $\Lambda_\varepsilon^\alpha \leq \frac{1}{2} |\log \varepsilon| - \sqrt{\varepsilon} \Lambda_\varepsilon^\alpha$ is bounded independently of $\varepsilon \leq 1$ and, using the inequality (3.5), we get

$$\int_0^{\xi(a)} \left(\log \frac{\rho}{r(a,t)} \right)_+ dt \leq \int_0^{\xi(a)} \left(\log \frac{\rho |\nabla \xi|_{L^\infty}}{\xi(a) - t} \right)_+ dt = \int_{\xi(a) - \rho |\nabla \xi|_{L^\infty}}^{\xi(a)} \log \frac{\rho |\nabla \xi|_{L^\infty}}{\xi(a) - t} dt,$$

and the rightmost integral is equal, by change of variables $u = \frac{\xi(a) - t}{\rho |\nabla \xi|_{L^\infty}}$, to $\rho |\nabla \xi|_{L^\infty}$. Therefore

$$\int_{t_\varepsilon}^{+\infty} f_\varepsilon^B(E_t) dt \geq -C |\nu_\varepsilon|(B) |\nabla \xi|_{L^\infty}.$$

In view of (3.3), adding (3.4) yields the result. \square

Remark 3.1. Note that in the proof of (3.2), the final radius ρ may be replaced by any $r \in (\sqrt{\varepsilon}, \rho)$. This yields the following result: Assume that $r \in (\sqrt{\varepsilon}, \rho)$ and that $B \in \mathcal{B}_\varepsilon^r$ is included in some ball in $\mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha$. Then, for any positive function ξ vanishing on $B \setminus \Omega_\varepsilon$,

$$(3.7) \quad \int_B (e_\varepsilon - \Lambda_\varepsilon^{\alpha,r} \nu_\varepsilon) \xi \geq -C |\nabla \xi|_{L^\infty(B)} |\nu_\varepsilon|(B).$$

We record the following lower bounds:

Proposition 3.2. *For ε small enough and $B \in \mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha$:*

$$(3.8) \quad e_\varepsilon(\Omega_\varepsilon \cap B) \geq \left(\frac{1}{8} |\log \varepsilon| - C \right) |\nu_\varepsilon|(B).$$

For ε small enough and $B \in \mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha$ such that $\text{dist}(B, \Omega_\varepsilon^c) > \varepsilon$, we have

$$(3.9) \quad g_\varepsilon^B(\Omega_\varepsilon \cap B) \geq \left(\frac{1}{8} |\log \varepsilon| - C \right) |\nu_\varepsilon|(B) - \frac{1}{2} |\log \varepsilon| |\nu_\varepsilon(B)|.$$

If in addition $d_B < 0$, then

$$(3.10) \quad g_\varepsilon^B(\Omega_\varepsilon \cap B) - \left(\frac{1}{2} |\log \varepsilon| - \Lambda_\varepsilon^\alpha \right) \nu_\varepsilon(B) \geq \left(\frac{1}{8} |\log \varepsilon| - C \right) |\nu_\varepsilon|(B).$$

The meaning of this lower bound is that $e_\varepsilon(B)$ is not only bounded below by $\Lambda_\varepsilon^\alpha |\nu_\varepsilon(B)|$, which to leading order is $\frac{1}{2} |\log \varepsilon| |\nu_\varepsilon(B)|$ — this is the positivity of g_ε^B in the above proposition — but also by some constant times $|\log \varepsilon| |\nu_\varepsilon|(B)$, even though the constant is no longer guaranteed to be the (optimal) value $1/2$. This information is valuable in the case where $|\nu_\varepsilon(B)|$ is much smaller than $|\nu_\varepsilon|(B)$. The precise value of the constants is unimportant.

Proof. As we noticed, $\bar{C}_\alpha < (\sqrt{\varepsilon}/\varepsilon)^{\frac{1}{2}}$ implies $\sqrt{\varepsilon}/(\varepsilon \bar{C}_\alpha) \geq \varepsilon^{-\frac{1}{4}}$ thus, using Proposition 2.3,

$$e_\varepsilon(B \cap \Omega_\varepsilon) \geq \sum_{\substack{B' \in \mathcal{B}_\varepsilon^{\sqrt{\varepsilon}} \\ B' \subset B \\ \text{dist}(B', \Omega_\varepsilon^c) > \varepsilon}} e_\varepsilon(B') \geq \sum_{\substack{B' \in \mathcal{B}_\varepsilon^{\sqrt{\varepsilon}} \\ B' \subset B \\ \text{dist}(B', \Omega_\varepsilon^c) > \varepsilon}} \pi |d_{B'}| \left(\log \varepsilon^{-\frac{1}{4}} - C \right) = |\nu_\varepsilon|(B) \left(\frac{1}{8} |\log \varepsilon| - \frac{C}{2} \right),$$

which proves the first assertion. Secondly, note that from (3.1), if $\text{dist}(B, \Omega_\varepsilon^c) > \varepsilon$, choosing ξ compactly supported in Ω_ε such that $\xi = 1$ in B , we have

$$f_\varepsilon^B(B \cap \Omega_\varepsilon) = g_\varepsilon^B(B \cap \Omega_\varepsilon).$$

Since $\Lambda_\varepsilon^\alpha \leq \frac{1}{2} |\log \varepsilon|$ we deduce (3.9) in view of

$$g_\varepsilon^B(B \cap \Omega_\varepsilon) = f_\varepsilon^B(B \cap \Omega_\varepsilon) \geq |\nu_\varepsilon|(B) \left(\frac{1}{8} |\log \varepsilon| - C \right) - \frac{1}{2} |\log \varepsilon| |\nu_\varepsilon(B)|.$$

For the last assertion, since $\nu_\varepsilon(B) = 2\pi d_B < 0$, we write

$$g_\varepsilon^B(B \cap \Omega_\varepsilon) - \left(\frac{1}{2} |\log \varepsilon| - \Lambda_\varepsilon^\alpha \right) \nu_\varepsilon(B) = e_\varepsilon(B \cap \Omega_\varepsilon) - \frac{1}{2} |\log \varepsilon| |\nu_\varepsilon(B)| \geq e_\varepsilon(B \cap \Omega_\varepsilon),$$

and this is bounded below using (3.8). \square

3.3 Mass displacement of the remainder

Proposition 3.1 will allow to replace f_ε^B by the positive g_ε^B , and we have

$$(3.11) \quad f_\varepsilon - \sum_{B \in \mathcal{B}_\varepsilon} f_\varepsilon^B = e_\varepsilon \mathbf{1}_{\mathcal{B}_\varepsilon^c} + \sum_\alpha \left(\frac{1}{2} |\log \varepsilon| - \Lambda_\varepsilon^\alpha \right) \nu_\alpha.$$

We now proceed to absorb the negative part of $f_\varepsilon - \sum f_\varepsilon^B$, which is $(\frac{1}{2} |\log \varepsilon| - \Lambda_\varepsilon^\alpha) (\nu_\alpha)_+$. This will be easy if $\bar{C}_\alpha = \frac{3e_\alpha}{|\log \varepsilon|}$ and if not, in view of (2.5), we have

$$0 \leq \frac{1}{2} |\log \varepsilon| - \Lambda_\varepsilon^\alpha \leq \frac{1}{2} \log n_\alpha + C,$$

which allows to bound the mass of the negative part by $C \sum_\alpha n_\alpha (\log n_\alpha + 1)$. Following the method in [SS1] (see also [SS4], Chap. 9), this will be balanced by a lower bound by $c[n_\alpha]^2$ for the energy on annuli surrounding U_α .

Recall that $U_\alpha = B(x_\alpha, \ell_0) \cap \Omega_\varepsilon$. We let $A_\alpha = B(x_\alpha, r_1)$, where $r_1 = 3\ell_0$. Choosing ℓ_0 small enough, we may require that

$$\text{diam}(A_\alpha) < 1 \quad \text{and} \quad \left(A_\alpha \cap \Omega_\varepsilon^c \neq \emptyset \implies A_\alpha \subset \left\{ x \mid \text{dist}(x, \partial\Omega_\varepsilon) < \frac{1}{2} \right\} \right).$$

We will denote below by m' a bound, uniform in ε for the overlap number of the $\{A_\alpha\}_\alpha$.

Now we choose ρ such that for any $\varepsilon > 0$

$$|T_\varepsilon^\alpha| \geq \ell_0, \quad \text{where } T_\varepsilon^\alpha \text{ is the set of } t \in (r_0, r_1) \text{ such that } \{|x - x_\alpha| = t\} \cap \mathcal{B}_\varepsilon = \emptyset,$$

where $r_0 = \ell_0$. Indeed, the number of U_β 's which intersect $B(x_\alpha, r_1)$ is bounded by a certain number N independent of ε and α . Choosing $\rho = \ell_0/N$, the sum of the radii of balls in $\cup_\beta \mathcal{B}_\varepsilon^\beta$ which intersect $B(x_\alpha, r_1)$ is bounded above by ℓ_0 , hence $|T_\varepsilon^\alpha| \geq (r_1 - r_0) - \ell_0 = \ell_0$.

3.3.1 Lower bounds on annuli

For any α let

$$(3.12) \quad (\tilde{g}_\varepsilon^\alpha)_+ = \frac{1}{4m'} \left(e_\varepsilon \mathbf{1}_{\mathcal{B}_\varepsilon^c} + \sum_{B \in \mathcal{B}_\varepsilon} g_\varepsilon^B \right) \mathbf{1}_{A_\alpha}, \quad (\tilde{g}_\varepsilon^\alpha)_- = \left(\frac{1}{2} |\log \varepsilon| - \Lambda_\varepsilon^\alpha \right) (\nu_\varepsilon)_+ \mathbf{1}_{\mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha},$$

and $\tilde{g}_\varepsilon^\alpha = (\tilde{g}_\varepsilon^\alpha)_+ - (\tilde{g}_\varepsilon^\alpha)_-$. We have

$$\tilde{g}_\varepsilon - \sum_\alpha \tilde{g}_\varepsilon^\alpha \geq \frac{3}{4} \left(e_\varepsilon \mathbf{1}_{\mathcal{B}_\varepsilon^c} + \sum_{B \in \mathcal{B}_\varepsilon} g_\varepsilon^B \right) + \sum_\alpha \left(\frac{1}{2} |\log \varepsilon| - \Lambda_\varepsilon^\alpha \right) (\nu_\varepsilon)_- \mathbf{1}_{\mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha}.$$

In particular

$$(\tilde{g}_\varepsilon^\alpha)_+(A_\alpha) \leq \frac{1}{3m'} \left(\tilde{g}_\varepsilon - \sum_\beta \tilde{g}_\varepsilon^\beta \right) (A_\alpha).$$

Proposition 3.3. *There exist $\varepsilon_0, C, c > 0$ such that if $\beta < 1/4$ in (1.1), then for any $\varepsilon < \varepsilon_0$ and any index α*

$$(3.13) \quad (\tilde{g}_\varepsilon^\alpha)_-(A_\alpha) \leq \pi n_\alpha (\beta |\log \varepsilon| + C).$$

If moreover $\text{dist}(A_\alpha, \Omega_\varepsilon^c) > \varepsilon$ then at least one of the following is true:

$$(3.14) \quad (\tilde{g}_\varepsilon^\alpha)_-(A_\alpha) \leq \pi n_\alpha (\beta |\log \varepsilon| + C), \quad (\tilde{g}_\varepsilon^\alpha)_+(A_\alpha) \geq c n_\alpha |\log \varepsilon|$$

or

$$(3.15) \quad (\tilde{g}_\varepsilon^\alpha)_-(A_\alpha) \leq \pi n_\alpha (\log n_\alpha + C), \quad (\tilde{g}_\varepsilon^\alpha)_+(A_\alpha) \geq c n_\alpha^2.$$

Proof. The bound (3.13) follows from (3.12), (2.6). Now assume $\text{dist}(A_\alpha, \Omega_\varepsilon^c) > \varepsilon$.

First, if $n_\alpha = 0$ then $(\tilde{g}_\varepsilon^\alpha)_- = 0$, $(\tilde{g}_\varepsilon^\alpha)_+ \geq 0$ hence (3.14) is true.

Second, if $3e_\alpha/|\log \varepsilon| \geq M n_\alpha$ then, since for $B \subset A_\alpha$, we have $g_\varepsilon^B(B) = f_\varepsilon^B(B) = e_\varepsilon(B) - \Lambda_\varepsilon^\alpha \nu_\varepsilon(B)$ and $\Lambda_\varepsilon^\alpha \leq \frac{1}{2} |\log \varepsilon|$ it follows that

$$\begin{aligned} (\tilde{g}_\varepsilon^\alpha)_+(A_\alpha) &\geq \frac{1}{4m'} \int_{A_\alpha} e_\varepsilon - \frac{1}{4m'} \Lambda_\varepsilon^\alpha \sum_{B \in \mathcal{B}_\varepsilon \cap A_\alpha} |d_B| \geq \frac{1}{4m'} \int_{U_\alpha} e_\varepsilon - \frac{1}{4m'} \Lambda_\varepsilon^\alpha \sum_{B \in \mathcal{B}_\varepsilon \cap A_\alpha} |d_B| \\ &\geq \frac{M}{12m'} n_\alpha |\log \varepsilon| - \pi n_\alpha |\log \varepsilon| \geq \left(\frac{M}{12m'} - \pi \right) n_\alpha |\log \varepsilon|. \end{aligned}$$

Together with (3.13), this implies (3.14) if M was chosen strictly greater than $12m'\pi$. The last case is that where $\overline{C}_\alpha = M n_\alpha$. Then $\frac{1}{2} |\log \varepsilon| - \Lambda_\varepsilon^\alpha = \frac{1}{2} \log n_\alpha + C$ and therefore, using (2.3),

$$(3.16) \quad (\tilde{g}_\varepsilon^\alpha)_-(A_\alpha) \leq 2\pi n_\alpha \left(\frac{1}{2} \log n_\alpha + C \right) \leq n_\alpha (\pi \beta |\log \varepsilon| + C).$$

We define

$$D_0^+ = \sum_{\substack{B \in \mathcal{B}_\varepsilon \\ B \subset B(x_\alpha, r_0) \\ d_B > 0}} d_B, \quad D_1^- = \sum_{\substack{B \in \mathcal{B}_\varepsilon \\ B \subset B(x_\alpha, r_1) \\ d_B < 0}} |d_B|,$$

and again we distinguish several cases.

First from (3.16) we will have proven (3.14) if we prove that

$$(3.17) \quad (\tilde{g}_\varepsilon^\alpha)_+(A_\alpha) \geq c n_\alpha |\log \varepsilon|,$$

for some $c > 0$. This inequality holds in the following two cases.

First case : $D_1^- > n_\alpha/20$. This means there is a significant proportion of balls with negative degrees. For each such negative ball we have from (3.10), and since $|\nu_\varepsilon|(B) \geq |\nu_\varepsilon(B)|$,

$$g_\varepsilon^B(B) \geq g_\varepsilon^B(B) - \left(\frac{1}{2} |\log \varepsilon| - \Lambda_\varepsilon^\alpha \right) \nu_\varepsilon(B) \geq \left(\frac{1}{8} |\log \varepsilon| - C \right) 2\pi |d_B|.$$

This implies that

$$(\tilde{g}_\varepsilon^\alpha)_+(A_\alpha) \geq \frac{1}{4m'} \left(\frac{1}{8} |\log \varepsilon| - C \right) 2\pi D_1^-,$$

hence (3.17) is satisfied when $D_1^- > n_\alpha/20$.

Second case : $D_0^+ \leq n_\alpha/10$ and $D_1^- \leq n_\alpha/20$. Then for each $B \in \mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha$, Proposition 3.2 yields

$$g_\varepsilon^B(B) \geq \begin{cases} \left(\frac{1}{8} |\log \varepsilon| - C \right) |\nu_\varepsilon|(B) - \frac{1}{2} |\log \varepsilon| |\nu_\varepsilon(B)| & \text{if } |d_B| > 0 \\ \left(\frac{1}{8} |\log \varepsilon| - C \right) |\nu_\varepsilon|(B) & \text{if } |d_B| < 0. \end{cases}$$

Summing with respect to B we find, since $B \in \mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha$ implies $B \subset B(x_\alpha, r_0)$, that

$$(\tilde{g}_\varepsilon^\alpha)_+(A_\alpha) \geq \frac{1}{4m'} \left(\frac{1}{8} |\log \varepsilon| - C \right) n_\alpha - \frac{1}{4m'} D_0^+ \frac{1}{2} |\log \varepsilon|,$$

which again yields (3.17) when $D_0^+ \leq n_\alpha/10$.

We are left with the third case, when $D_0^+ \geq \frac{n_\alpha}{10}$ and $D_1^- \leq \frac{n_\alpha}{20}$. In this case (3.17) and then (3.14) do not necessarily hold. We need to prove (3.15) instead, which in view of (3.16) reduces to proving

$$(\tilde{g}_\varepsilon^\alpha)_+(A_\alpha) \geq cn_\alpha^2.$$

For this we really need to use the lower bounds on annuli of the type first introduced in [SS1]. We denote

$$\mathcal{C}_\varepsilon^\alpha = B(x_\alpha, r_1) \setminus (B(x_\alpha, r_0) \cup \mathcal{B}_\varepsilon).$$

For any $t \in T_\varepsilon^\alpha$ we let $B_t = B(x_\alpha, t)$ and $\gamma_t = \partial B_t$ and recall that γ_t does not intersect \mathcal{B}_ε . If $t \in T_\varepsilon^\alpha$ then $|u_\varepsilon| \geq 1/2$ on γ_t because of (2.1) and the fact that $\text{dist}(A_\alpha, \Omega_\varepsilon^c) > \varepsilon$.

It follows (see for instance [SS4] Lemma 4.4, or (5.4) below) that for some constant $c > 0$ we have

$$(3.18) \quad \int_{\gamma_t} \left(\frac{1}{2} |\nabla_A u|^2 + \frac{(1 - |u|^2)^2}{4\varepsilon^2} \right) + \frac{1}{2} \int_{B_t} (\text{curl } A)^2 \geq c \frac{|d_\varepsilon^t|^2}{t},$$

where d_ε^t is the degree of $u_\varepsilon/|u_\varepsilon|$ on γ_t . Integrating (3.18) with respect to $t \in T_\varepsilon^\alpha$, which has measure less than 1, the left-hand side will be bounded above by $e_\varepsilon(A_\alpha)$. In view of the lower bound $d_\varepsilon^t \geq (D_0^+ - D_1^-)$, which is valid for any $t \in T_\varepsilon^\alpha$, since $|T_\varepsilon^\alpha| \geq \ell_0$, and from the assumption on D_0^+ and D_1^- we deduce that

$$e_\varepsilon(A_\alpha \setminus \mathcal{B}_\varepsilon) \geq c (D_0^+ - D_1^-)^2 \geq cn_\alpha^2.$$

Then, since $(\tilde{g}_\varepsilon^\alpha)_+ = \frac{1}{4m'} e_\varepsilon$ on $(\mathcal{B}_\varepsilon)^c$ we deduce $(\tilde{g}_\varepsilon^\alpha)_+(A_\alpha) \geq cn_\alpha^2$ and (3.15) is proved. \square

3.4 Proof of Theorem 1 and Corollary 1.1

Item 1). The estimate (1.2) was already mentioned after the definition (2.2) of ν_ε , and the bound $|\nu_\varepsilon|(E) \leq C e_\varepsilon(\widehat{E})/|\log \varepsilon|$ was proved in Lemma 2.1.

Item 2). We define

$$f_\alpha = \sum_{B \in \mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha} (f_\varepsilon^B - g_\varepsilon^B) + (\tilde{g}_\varepsilon^\alpha)_+ - (\tilde{g}_\varepsilon^\alpha)_-.$$

Then clearly f_α is supported in A_α . Moreover, using the fact (see (3.11)) that

$$f_\varepsilon - \sum_{B \in \mathcal{B}_\varepsilon} f_\varepsilon^B = e_\varepsilon \mathbf{1}_{\mathcal{B}_\varepsilon^c} - \sum_\alpha \left(\frac{1}{2} |\log \varepsilon| - \Lambda_\varepsilon^\alpha \right) \nu_\alpha$$

and since $\sum_\alpha \mathbf{1}_{A_\alpha} \leq m'$ we easily obtain

(3.19)

$$f_\varepsilon - \sum_\alpha f_\alpha = \sum_\alpha \left(\frac{1}{2} |\log \varepsilon| - \Lambda_\varepsilon^\alpha \right) (\nu_\alpha)_- + \left(e_\varepsilon \mathbf{1}_{\mathcal{B}_\varepsilon^c} + \sum_{B \in \mathcal{B}_\varepsilon} g_\varepsilon^B \right) \left(1 - \frac{1}{4m'} \sum_\alpha \mathbf{1}_{A_\alpha} \right).$$

Since $\sum_\alpha \mathbf{1}_{A_\alpha} \leq m'$ we find

$$(3.20) \quad f_\varepsilon - \sum_\alpha f_\alpha \geq \sum_\alpha \left(\frac{1}{2} |\log \varepsilon| - \Lambda_\varepsilon^\alpha \right) (\nu_\alpha)_- + \frac{3}{4} \left(e_\varepsilon \mathbf{1}_{\mathcal{B}_\varepsilon^c} + \sum_{B \in \mathcal{B}_\varepsilon} g_\varepsilon^B \right) \geq 0.$$

Item 3). We define g_α . In the case $\text{dist}(A_\alpha, \Omega_\varepsilon^c) \leq \varepsilon$ we let $g_\alpha = (\tilde{g}_\varepsilon^\alpha)_+$. Then

$$\int \xi d(f_\alpha - g_\alpha) = \sum_{B \in \mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha} \int \xi d(f_\varepsilon^B - g_\varepsilon^B) - \int \xi d(\tilde{g}_\varepsilon^\alpha)_-.$$

This implies (1.5), summing (3.1) over $B \in \mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha$ and using (3.13).

In the case $\text{dist}(A_\alpha, \Omega_\varepsilon^c) > \varepsilon$ we let

$$c_\alpha = \left(\frac{\tilde{g}_\varepsilon^\alpha(A_\alpha)}{|A_\alpha|} \right)_-.$$

We deduce easily from (3.14), (3.15) and if β is small enough that $c_\alpha \leq C$ and applying Lemma 3.2 in A_α to $\tilde{g}_\varepsilon^\alpha + c_\alpha$ we obtain φ_α defined on A_α and such that $0 \leq \varphi_\alpha \leq (\tilde{g}_\varepsilon^\alpha)_+ + c_\alpha$ and, for any Lipschitz function ξ ,

$$\int_{A_\alpha} \xi d(\tilde{g}_\varepsilon^\alpha - g_\alpha) \leq C |\nabla \xi|_{L^\infty(A_\alpha)} (\tilde{g}_\varepsilon^\alpha)_-(A_\alpha), \quad \text{where } g_\alpha := \varphi_\alpha - c_\alpha.$$

Moreover $-C \leq -c_\alpha \leq g_\alpha \leq (\tilde{g}_\varepsilon^\alpha)_+$.

Then

$$\begin{aligned}
\int_{A_\alpha} \xi d(f_\alpha - g_\alpha) &= \int_{A_\alpha} \xi d(f_\alpha - \tilde{g}_\varepsilon^\alpha) + \int_{A_\alpha} \xi d(\tilde{g}_\varepsilon^\alpha - g_\alpha) \\
(3.21) \qquad \qquad \qquad &= \sum_{B \in \mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha} \int \xi d(f_\varepsilon^B - g_\varepsilon^B) + \int_{A_\alpha} \xi d(\tilde{g}_\varepsilon^\alpha - g_\alpha) \\
&\leq C |\nabla \xi|_{L^\infty(A_\alpha)} (n_\alpha + (\tilde{g}_\varepsilon^\alpha)_-(A_\alpha)),
\end{aligned}$$

where we have used (3.1) to bound the integral involving $f_\varepsilon^B - g_\varepsilon^B$. Moreover, $g_\alpha(A_\alpha) = \tilde{g}_\varepsilon^\alpha(A_\alpha)$.

If (3.14) holds, then (1.3) follows immediately from (3.21) when $\pi\beta < c/2$, with c the constant in (3.14). If (3.15) holds we deduce (1.4) from (3.21) by noting that $cn_\alpha^2 - Cn_\alpha(\log n_\alpha + 1) \geq \frac{c}{2}n_\alpha^2 - C'n_\alpha$ if C' is chosen large enough depending on c, C .

Item 4), (1.8). We adapt an argument in [St].

First, $g_\varepsilon - \sum_\alpha g_\alpha = f_\varepsilon - \sum_\alpha f_\alpha$ thus from (3.20) and since $\sum_\alpha g_\alpha \geq -C$ we find

$$(3.22) \qquad \qquad \qquad g_\varepsilon \geq \frac{3}{4} \left(e_\varepsilon \mathbf{1}_{\mathcal{B}_\varepsilon^c} + \sum_{B \in \mathcal{B}_\varepsilon} g_\varepsilon^B \right) - C.$$

Then, assuming $U_\alpha \subset \Omega_\varepsilon$, denote by $\mathcal{B}_\varepsilon^{r,\alpha}$ the set of balls in $\mathcal{B}_\varepsilon^r$ which are included in some ball belonging to $\mathcal{B}_\varepsilon^\alpha \cap \mathcal{B}_\varepsilon$, so that $\nu_\alpha(\mathcal{B}_\varepsilon) = \nu_\varepsilon(\mathcal{B}_\varepsilon^\alpha \cap \mathcal{B}_\varepsilon) = \nu_\varepsilon(\mathcal{B}_\varepsilon^{r,\alpha})$. Applying Remark 3.1 for some $r \in (\sqrt{\varepsilon}, \rho)$ with $\xi = 1$ and summing (3.7) over $B \in \mathcal{B}_\varepsilon^{r,\alpha}$ we find $e_\varepsilon(\mathcal{B}_\varepsilon^{r,\alpha}) \geq \Lambda_\varepsilon^{\alpha,r} \nu_\varepsilon(\mathcal{B}_\varepsilon^{r,\alpha})$ and then

$$\begin{aligned}
e_\varepsilon(\mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha \setminus \mathcal{B}_\varepsilon^{r,\alpha}) &\leq e_\varepsilon(\mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha) - \Lambda_\varepsilon^{\alpha,r} \nu_\varepsilon(\mathcal{B}_\varepsilon) + (\Lambda_\varepsilon^\alpha - \Lambda_\varepsilon^{\alpha,r}) \nu_\alpha(\mathcal{B}_\varepsilon) \\
&= \sum_{B \in \mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha} g_\varepsilon^B(B) + \frac{1}{2} \log \frac{1}{r} \nu_\alpha(\mathcal{B}_\varepsilon),
\end{aligned}$$

where we have used the fact that $f_\varepsilon^B(B) = g_\varepsilon^B(B)$. It follows using (3.22) that

$$(3.23) \qquad \qquad \qquad e_\varepsilon(\mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha \setminus \mathcal{B}_\varepsilon^{r,\alpha}) \leq C \left((g_\varepsilon)_+(U_\alpha) + n_\alpha \log \frac{1}{r} + 1 \right).$$

Then comes the argument in [St]: For any integer k , let $r_k = 2^{-k}\rho$, and let \mathcal{C}_k be the intersection of $\mathcal{B}_\varepsilon^{r_k} \setminus \mathcal{B}_\varepsilon^{r_{k+1}}$ and $\mathcal{B}_\varepsilon^\alpha$. Then $|\mathcal{C}_k| \leq C2^{-2k}\rho^2$, since $\rho2^{-k}$ bounds the total radius of the balls in $\mathcal{B}_\varepsilon^{r_k} \cap \mathcal{B}_\varepsilon^\alpha$. Moreover $j_\varepsilon = (iu_\varepsilon, \nabla u_\varepsilon - iA_\varepsilon)$ and thus assuming $|u_\varepsilon| \leq 1$ we have $|j_\varepsilon|^2 \leq 2e_\varepsilon$. Then using Hölder's inequality in \mathcal{C}_k and (3.23) we find for $p < 2$

$$\begin{aligned}
\int_{\mathcal{C}_k} |j_\varepsilon|^p &\leq |\mathcal{C}_k|^{1-p/2} (e_\varepsilon(\mathcal{C}_k))^{p/2} \leq |\mathcal{C}_k|^{1-p/2} (e_\varepsilon(\mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha \setminus \mathcal{B}_\varepsilon^{r_{k+1}}))^{p/2} \\
&\leq C_p 2^{-(2-p)k} (e_\varepsilon(\mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha \setminus \mathcal{B}_\varepsilon^{r_{k+1}}))^{p/2} \\
&\leq C_p 2^{-(2-p)k} \left((g_\varepsilon)_+(U_\alpha) + kn_\alpha \log 2 + 1 \right)^{p/2}.
\end{aligned}$$

Using (1.10) we find

$$\int_{C_k} |j_\varepsilon|^p \leq C_p 2^{-(2-p)k} (1 + k \log 2)^{p/2} ((g_\varepsilon)_+(U_\alpha) + 1)^{p/2}.$$

Summing these inequalities for k ranging from 0 to the largest integer K such that $r_K \geq \sqrt{\varepsilon}$ — so that in particular $r_K \leq 2\sqrt{\varepsilon}$ — we find

$$\int_{\mathcal{B}_\varepsilon \cap \mathcal{B}_\varepsilon^\alpha \setminus \mathcal{B}_\varepsilon^{2\sqrt{\varepsilon}}} |j_\varepsilon|^p \leq C_p ((g_\varepsilon)_+(U_\alpha) + 1)^{p/2},$$

where C_p is a constant times the sum of the convergent series $\sum_k 2^{-(2-p)k} (1 + k \log 2 - \log \rho)^{p/2}$. To this inequality we add

$$\int_{\mathcal{B}_\varepsilon^\alpha \cap \mathcal{B}_\varepsilon^{2\sqrt{\varepsilon}}} |j_\varepsilon|^p \leq C \varepsilon^{1-p/2} e_\varepsilon(U_\alpha)^{p/2},$$

which follows from Hölder's inequality after estimating as above $|\mathcal{B}_\varepsilon^{2\sqrt{\varepsilon}} \cap \mathcal{B}_\varepsilon^\alpha|$ by $C\varepsilon$. But since $e_\varepsilon = f_\varepsilon + \frac{1}{2}|\log \varepsilon| \nu_\varepsilon$ we may write using (1.9), (1.10),

$$(3.24) \quad e_\varepsilon(U_\alpha) \leq C(g_\varepsilon)_+(\widehat{U}_\alpha) + C|\nu_\varepsilon|(\widehat{U}_\alpha)(1 + |\log \varepsilon|) \leq C|\log \varepsilon| ((g_\varepsilon)_+(U_\alpha + B(0, 2)) + 1).$$

Thus

$$\begin{aligned} \int_{\mathcal{B}_\varepsilon^\alpha \cap \mathcal{B}_\varepsilon^{2\sqrt{\varepsilon}}} |j_\varepsilon|^p &\leq C \varepsilon^{1-\frac{p}{2}} |\log \varepsilon|^{\frac{p}{2}} \left((g_\varepsilon)_+(U_\alpha + B(0, 2))^{\frac{p}{2}} + 1 \right) \\ &\leq C \left((g_\varepsilon)_+(U_\alpha + B(0, 2))^{\frac{p}{2}} + 1 \right). \end{aligned}$$

We also add

$$\int_{U_\alpha \setminus \mathcal{B}_\varepsilon} |j_\varepsilon|^p \leq C ((g_\varepsilon)_+(U_\alpha) + 1)$$

which follows from (3.22). Finally we obtain

$$\int_{U_\alpha} |j_\varepsilon|^p \leq C_p ((g_\varepsilon)_+(U_\alpha + B(0, 2)) + 1).$$

Summing with respect to the α 's such that $E \cap U_\alpha \neq \emptyset$, this proves (1.8) and concludes the proof of Theorem 1.

Proof of Corollary 1.1. Note that

$$\int \xi d(f_\varepsilon - g_\varepsilon) = \sum_\alpha \int \xi d(f_\alpha - g_\alpha).$$

Three types of indices occur.

First we consider indices α such that $\text{dist}(A_\alpha, \Omega_\varepsilon^c) > \varepsilon$ and (1.3) holds. Since

$$(3.25) \quad g_\alpha \leq g_\varepsilon - \sum_{\beta \neq \alpha} g_\beta \leq g_\varepsilon + C,$$

we deduce from (1.3) that if $n_\alpha \geq 1$ and ε is small enough, $g_\varepsilon(A_\alpha) \geq cn_\alpha |\log \varepsilon|$ and then using (1.3) again that

$$(3.26) \quad \int \xi d(f_\alpha - g_\alpha) \leq C |\nabla \xi|_{L^\infty(A_\alpha)} (n_\alpha + \beta(g_\varepsilon)_+(A_\alpha)).$$

If $n_\alpha = 0$ the same inequality holds since from (1.3) the left-hand side is zero.

Second we consider indices α such that $\text{dist}(A_\alpha, \Omega_\varepsilon^c) > \varepsilon$ and (1.4) holds. We note that if C is large enough then $x \log x \leq \eta x^2 + C \log^2 \eta / \eta$ holds for every $x > 0$ and $\eta \leq 1$, for instance by distinguishing the cases $\eta > \frac{\log x}{x}$ and $\eta \leq \frac{\log x}{x}$. We use this and (3.25), together with (1.4) to find that if $n_\alpha \geq 1$ then

$$(3.27) \quad \int \xi d(f_\alpha - g_\alpha) \leq C |\nabla \xi|_{L^\infty(A_\alpha)} \left(n_\alpha + \eta(g_\varepsilon)_+(A_\alpha) + \frac{\log^2 \eta}{\eta} \right).$$

Again the inequality is true if $n_\alpha = 0$ since from (1.4) the left-hand side is zero in this case.

Finally we consider indices α such that $\text{dist}(A_\alpha, \Omega_\varepsilon^c) \leq \varepsilon$. In this case, noting that from Lemma 2.1 we have $n_\alpha |\log \varepsilon| \leq C e_\varepsilon(A_\alpha)$, we rewrite (1.5) as

$$(3.28) \quad \int \xi d(f_\alpha - g_\alpha) \leq C (|\nabla \xi|_{L^\infty(A_\alpha)} n_\alpha + \beta |\xi|_{L^\infty(A_\alpha)} e_\varepsilon(A_\alpha)).$$

To conclude we sum either (3.26), (3.27) or (3.28) according to the type of index α , noting that since $\text{diam}(A_\alpha) \leq 1$, we have $|f|_{L^\infty(A_\alpha)} \leq \widehat{f}$ on A_α for any function f . Since the overlap number of the A_α 's is bounded by a universal constant, we deduce (1.9).

We prove (1.10). We start by proving that when $\text{dist}(A_\alpha, \Omega_\varepsilon^c) > \varepsilon$ we have

$$(3.29) \quad \min(n_\alpha^2, n_\alpha |\log \varepsilon|) \leq C ((g_\varepsilon)_+(A_\alpha) + 1).$$

If $n_\alpha = 0$ this is trivial, if not then it follows from either (1.3) or (1.4) using (3.25).

Assume α is such that $\text{dist}(A_\alpha, \Omega_\varepsilon^c) > \varepsilon$, then since $2x \leq \eta x^2 + 1/\eta$ and since $x \leq \eta x |\log \varepsilon|$ is trivially true if $1/|\log \varepsilon| < \eta$, we deduce from (3.29) that

$$(3.30) \quad n_\alpha \leq C (\eta(g_\varepsilon)_+(A_\alpha) + 1/\eta).$$

On the other hand Lemma 2.1 implies that for any α

$$(3.31) \quad n_\alpha \leq C \frac{e_\varepsilon(A_\alpha \cap \Omega_\varepsilon)}{|\log \varepsilon|}.$$

Summing (3.30) or (3.31) according to whether $\text{dist}(A_\alpha, \Omega_\varepsilon^c)$ is $> \varepsilon$ or $\leq \varepsilon$ we deduce (1.10).

4 Proof of Theorem 2

4.1 Convergence

We study the consequences of the hypothesis

$$(4.1) \quad \forall R > 0, \quad M_R := \limsup_{\varepsilon \rightarrow 0} \int_{\mathbf{U}_R} g_\varepsilon(x) dx < +\infty.$$

and prove that it implies the convergence of the vorticities and currents in the appropriate sense.

Note that we assume $\text{dist}(0, \partial\Omega_\varepsilon) \rightarrow +\infty$ so that for every R , $\mathbf{U}_R \subset \Omega_\varepsilon$ for ε small enough. From (1.13) there exists $C > 0$ such that for any R large enough

$$B_{R/C} \subset \mathbf{U}_R \subset B_{CR}, \quad \frac{1}{C} \leq \frac{|\mathbf{U}_R|}{R^2} \leq C.$$

We now gather several easy consequences of Theorem 1 and (4.1).

Proposition 4.1. *Assume (4.1) holds, and let g_ε be as in Theorem 1. Then for any R and ε small enough depending on R we have*

$$(4.2) \quad \sum_{\alpha | A_\alpha \subset \mathbf{U}_R} \min(n_\alpha^2, n_\alpha |\log \varepsilon|) \leq C(M_{R+C} + R^2),$$

$$(4.3) \quad |\nu_\varepsilon|(\mathbf{U}_R) \leq C(M_{R+C} + R^2),$$

$$(4.4) \quad \int (f_\varepsilon - g_\varepsilon) \chi_{\mathbf{U}_R} \leq C \sum_{\alpha | A_\alpha \subset \mathbf{U}_{R+C} \setminus \mathbf{U}_{R-C}} n_\alpha (\log n_\alpha + 1) \leq C(M_{R+C} + R^2),$$

where $\{\chi_{\mathbf{U}_R}\}_R$ are any functions satisfying (1.14).

For any $1 \leq p < 2$ there exists $C_p > 0$ such that for any $R > 0$, and ε small enough

$$(4.5) \quad \int_{\mathbf{U}_R} |j_\varepsilon|^p \leq C_p(M_{R+C} + R^2).$$

Moreover, up to extraction of a subsequence, $\{j_\varepsilon\}_\varepsilon$ converges weakly in $L_{\text{loc}}^p(\mathbb{R}^2)$, $p < 2$ to some $j : \mathbb{R}^2 \rightarrow \mathbb{R}^2$; $\{\nu_\varepsilon\}_\varepsilon$ converges in the weak sense of measures to a measure ν on \mathbb{R}^2 of the form $2\pi \sum_{p \in \Lambda} d_p \delta_p$ where Λ is a discrete set and $d_p \in \mathbb{Z}$, $\{\mu_\varepsilon\}_\varepsilon$ converges to the same ν in $W_{\text{loc}}^{-1,p}(\mathbb{R}^2)$ for any $p < 2$ and $\{h_\varepsilon\}$ converges weakly in $L_{\text{loc}}^2(\mathbb{R}^2)$ to h . Moreover it holds that

$$(4.6) \quad \text{curl } j = \nu - h.$$

Proof. Assertions (4.2), (4.3) and (4.5) are direct consequences of (3.29), (1.10) and (1.8), respectively.

We prove (4.4). First we note that a consequence of (4.1) is that for every $R > 0$, if $\varepsilon > 0$ is small enough and $A_\alpha \subset \mathbf{U}_R$ then (1.4) holds. Indeed if (1.3) is true with $n_\alpha \geq 1$ (note that if $n_\alpha = 0$ then (1.3) and (1.4) are identical) then $g_\varepsilon(A_\alpha) \geq c|\log \varepsilon| - C$, using (3.25), which contradicts (4.1) if ε is small enough.

Then we use (1.3) with $\xi = \chi_{\mathbf{U}_R}$. From the fact that $\chi_{\mathbf{U}_R}$ is supported in \mathbf{U}_{R+C} and since $\text{dist}(\mathbf{U}_{R+C}, \partial\Omega_\varepsilon) \rightarrow +\infty$ we have, if ε is small enough and $A_\alpha \cap \mathbf{U}_{R+C} \neq \emptyset$, that $\text{dist}(A_\alpha, \partial\Omega_\varepsilon) > \varepsilon$. Then summing (1.3) over all such α we find

$$\int \chi_{\mathbf{U}_R} d(f_\varepsilon - g_\varepsilon) \leq C \sum_{A_\alpha \subset \mathbf{U}_{R+C} \stackrel{\alpha \text{ s.t.}}{\setminus} \mathbf{U}_{R-C}} n_\alpha (\log n_\alpha + 1),$$

which is the first inequality in (4.4). The second one then easily follows from (3.29).

We now turn to the convergence results. The weak local convergence of j_ε follows from a bound for $\int_{\mathbf{U}_R} |j_\varepsilon|^p$ valid for any ε small enough, depending on R , which is implied by (4.1) and (4.5). From (4.3), $\{\nu_\varepsilon\}_\varepsilon$ is bounded on any compact subset of \mathbb{R}^2 , hence converges (up to extraction) to a measure ν , which by (2.2) has to be of the form $2\pi \sum_{p \in \Lambda} d_p \delta_p$ where Λ is a discrete set and $d_p \in \mathbb{Z}$ for every $p \in \Lambda$ (we will prove below that $d_p = 1$).

The weak local convergence of h_ε follows from Remark 1.12 combined with the bound (4.1).

The convergence of $\{\mu_\varepsilon\}_\varepsilon$ in $W_{\text{loc}}^{-1,p}$ uses the Jacobian estimate (see [JS] or [SS4], Theorem 6.2) from which we deduce that for any $R > 0$ and any $\gamma \in (0, 1)$, and since $r(\mathcal{B}_\varepsilon^{\sqrt{\varepsilon}} \cap B_R) \leq C\sqrt{\varepsilon}$,

$$(4.7) \quad \|\mu_\varepsilon - \nu_\varepsilon\|_{(C_0^{0,\gamma}(B_R))^*} \leq C(\sqrt{\varepsilon})^\gamma (e_\varepsilon(B_R) + 1),$$

where C depends on R but not on ε .

But $\{\nu_\varepsilon\}_\varepsilon$ is bounded in B_R as measures, hence in $(C_0^{0,\gamma})^*$, and arguing again as in (3.24),

$$e_\varepsilon(B_R) \leq (g_\varepsilon)_+(B_{R+1}) + \frac{1}{2}|\log \varepsilon| |\nu_\varepsilon|(B_{R+C}) \leq C|\log \varepsilon|$$

therefore the right-hand side in (4.7) tends to 0 as $\varepsilon \rightarrow 0$ and $\{\mu_\varepsilon\}_\varepsilon$ is bounded in $(C_0^{0,\gamma}(B_R))^*$. We deduce that $\mu_\varepsilon \rightarrow \nu$ in $W_{\text{loc}}^{-1,p}$ by noting that for any $1 < p < 2$ there exists $0 < \gamma < 1$ such that $W_0^{1,p'}(B_R) \hookrightarrow C_0^{0,\gamma}$ with compact imbedding — where $1/p + 1/p' = 1$ — which implies by duality that $(C_0^{0,\gamma})^* \hookrightarrow W_0^{-1,p}$ with compact imbedding.

Finally (4.6) is obtained by passing to the limit in $\mu_\varepsilon = \text{curl } j_\varepsilon + \text{curl } A_\varepsilon$ since by Remark 1.2 we may assume (up to extraction) that $\text{curl } A_\varepsilon \rightarrow h$ weakly locally in L^2 as $\varepsilon \rightarrow 0$. \square

Remark 4.1. *From the above results, it is easy to deduce (1.18) under the stronger assumption (1.17). In this case we have $M_R \leq CR^2$ and therefore (4.3), (4.5) and Remark (1.2) imply that*

$$(4.8) \quad |\nu_\varepsilon|(\mathbf{U}_R) \leq CR^2, \quad \int_{\mathbf{U}_R} |j_\varepsilon|^p \leq CR^2, \quad \int_{\mathbf{U}_R} |h_\varepsilon|^2 \leq CR^2$$

which in turn implies (1.18).

4.2 Lower bound by the renormalized energy

We turn to the proof of the remaining statement in Theorem 2, namely that ν is of the form $2\pi \sum_{p \in \Lambda} \delta_p$ (we already know it is of the form $2\pi \sum_{p \in \Lambda} d_p \delta_p$, where the d_p 's are nonzero integers) and that under assumption (1.17) the lower bound (1.19) holds. Both are related to a lower bound of $\int \chi_R g_\varepsilon$ by the renormalized energy, where $\chi_R := \chi_{\mathbf{U}_R}$. This reproduces more or less arguments present in [BBH] and [BR]. Throughout this subsection we assume that (1.17) holds, and begin by bounding from below the integral of $(e_\varepsilon - \frac{1}{2}|\log \varepsilon| \nu_\varepsilon) \chi_R$.

Choose $R > 0$. From (4.3) we have that $|\nu_\varepsilon|$ is bounded independently of ε on the support of χ_R , thus a subsequence of $\{|\nu_\varepsilon| \mathbf{1}_{\text{Supp } \chi_R}\}_\varepsilon$ converges to a positive measure $\tilde{\nu}$ of the form $2\pi \sum_{i=1}^k k_i \delta_{a_i}$, where k_i is a positive integer for every i (the a_i 's are a subset of Λ).

From the weak convergence of j_ε to j in L^p_{loc} and using the inequality $|\nabla_{A_\varepsilon} u_\varepsilon| \geq |j_\varepsilon|$ (following from the assumption $|u_\varepsilon| \leq 1$) we have for any $r > 0$

$$(4.9) \quad \liminf_{\varepsilon \rightarrow 0} \int_{\mathbb{R}^2 \setminus \cup_{p \in \Lambda} B(p, r)} \chi_R |\nabla_{A_\varepsilon} u_\varepsilon|^2 \geq \liminf_{\varepsilon \rightarrow 0} \int_{\mathbb{R}^2 \setminus \cup_{p \in \Lambda} B(p, r)} \chi_R |j_\varepsilon|^2 \geq \int_{\mathbb{R}^2 \setminus \cup_{p \in \Lambda} B(p, r)} \chi_R |j|^2.$$

Indeed either the left-hand side is equal to $+\infty$ and the statement is true, or there is weak L^2 convergence of the currents on the complement of $\cup_p B(p, r)$ and (4.9) follows by weak lower semicontinuity of the integrand. Similarly, by weak convergence of h_ε to h we have

$$(4.10) \quad \liminf_{\varepsilon \rightarrow 0} \int_{\mathbb{R}^2 \setminus \cup_{p \in \Lambda} B(p, r)} \chi_R h_\varepsilon^2 \geq \int_{\mathbb{R}^2 \setminus \cup_{p \in \Lambda} B(p, r)} \chi_R h^2.$$

Then consider any $\eta \in (0, 1)$ small enough so that the balls $B(a_i, 2\eta)$ are disjoint. Note that since the limit of $|\nu_\varepsilon|$ on the support of χ_R is a sum of Dirac masses concentrated at the points $\{a_i\}_i$ we have for ε small enough

$$|\nu_\varepsilon|(\text{Supp } \chi_R \setminus \cup_i B(a_i, \eta)) = 0, \quad \nu_\varepsilon(B(a_i, \eta)) = 2\pi d_i,$$

where $2\pi d_i = \nu(a_i)$.

We use two distinct lower bounds for the integral of $\chi_R(e_\varepsilon - \frac{1}{2}|\log \varepsilon| \nu_\varepsilon)$ on balls. We distinguish the set I of indices such that $B(a_i, 2\eta) \subset \{\chi_R = 1\}$ and the remaining indices J . Note that if $i \in J$ then $B(a_i, 2\eta)$ intersects the set where $\chi_R \neq 1$ and the support of χ_R , thus $B(a_i, 2\eta) \subset \mathbf{U}_{R+C} \setminus \mathbf{U}_{R-C}$ for some $C > 0$ independent of $R > 0$, $\eta \in (0, 1)$ and i .

In the case $i \in I$ we use

$$(4.11) \quad \int_{B(a_i, \eta)} e_\varepsilon \geq \pi |d_i| \log \frac{\eta}{\varepsilon} + C_{|d_i|} + o_{\eta, \varepsilon}(1),$$

where C_d is a constant depending only on d such that $C_1 = \gamma$, (where γ is defined after Theorem 1), where $C_0 = 0$, and where

$$\lim_{\eta \rightarrow 0} \limsup_{\varepsilon \rightarrow 0} o_{\eta, \varepsilon}(1) = 0.$$

We postpone the proof of this well-known statement. It is very similar to analogous ones found in [BBH] or [BR]. Then we deduce from (4.11) that for any $i \in I$ and letting $C_{d_i} = +\infty$ if $d_i < 0$,

$$(4.12) \quad \liminf_{\varepsilon \rightarrow 0} \int_{B(a_i, \eta)} \left(e_\varepsilon - \frac{1}{2} |\log \varepsilon| d\nu_\varepsilon \right) \geq \pi d_i \log \eta + C_{d_i} + o_\eta(1),$$

where $\lim_{\eta \rightarrow 0} o_\eta(1) = 0$.

In the case $i \in J$ we have to introduce the weight χ_R which is no longer constant on the ball. Then we resort to Remark 3.1. Consider the family of balls \mathcal{C}_ε consisting of the balls B in $\mathcal{B}_\varepsilon^{\eta/2}$ which intersect the support of χ_{R+1} , and such that $|\nu_\varepsilon|(B) \neq \emptyset$. For any $B \in \mathcal{C}_\varepsilon$, since $|\nu_\varepsilon|(B) \neq 0$ and $|\nu_\varepsilon| \rightarrow 2\pi \sum_i k_i \delta_{a_i}$, and since $r(B) \leq \eta/2$, we have for ε small enough depending on R that there is some index i for which $B \subset B(a_i, \eta)$. Let $\mathcal{C}_\varepsilon^i$ denote the balls included in $B(a_i, \eta)$ and partition $\mathcal{C}_\varepsilon^i$ as $\cup_\alpha \mathcal{C}_\varepsilon^{i, \alpha}$ and \mathcal{C}_ε as $\cup_\alpha \mathcal{C}_\varepsilon^\alpha$ where the superscript α corresponds to the balls which are included in a ball $B \in \mathcal{B}_\varepsilon^\alpha$ (we assume $\eta/2 < \rho$).

From (3.7), for every $B \in \mathcal{C}_\varepsilon^{i, \alpha}$

$$(4.13) \quad \int_B \chi_R (e_\varepsilon - \Lambda_\varepsilon^{\alpha, \eta/2} d\nu_\varepsilon) \geq -C |\nabla \chi_R|_\infty |\nu_\varepsilon|(B) \geq -C |\nu_\varepsilon|(B).$$

Now we note that since (1.17) holds, then for ε small enough $\overline{C}_\alpha = Mn_\alpha$, for otherwise we would have $e_\varepsilon(\mathcal{B}_\varepsilon^\alpha) \geq \frac{M}{3} n_\alpha |\log \varepsilon|$ and then that

$$\sum_{B \in \mathcal{B}_\varepsilon^\alpha} g_\varepsilon^B(B) = \sum_{B \in \mathcal{B}_\varepsilon^\alpha} f_\varepsilon^B(B) = \sum_{B \in \mathcal{B}_\varepsilon^\alpha} (e_\varepsilon - \Lambda_\varepsilon^\alpha \nu_\varepsilon)(B) \geq \left(\frac{M}{3} - \pi \right) n_\alpha |\log \varepsilon| \xrightarrow{\varepsilon \rightarrow 0} +\infty,$$

if we choose $M > 3\pi$ and since $n_\alpha \geq 1$. This is a contradiction with (1.16) since $g_\varepsilon \geq \sum_B g_\varepsilon^B - C$ by (1.6), proving that $\overline{C}_\alpha = Mn_\alpha$.

Then we have from (2.5) that

$$\Lambda_\varepsilon^{\alpha, \eta/2} - \frac{1}{2} |\log \varepsilon| = \frac{1}{2} \log \eta + \Delta, \quad \text{where } |\Delta| \leq C (\log n_\alpha + 1)$$

and

$$\left| \int_B (\chi_R - \chi_R(a_i)) d\nu_\varepsilon \right| \leq C \eta |\nu_\varepsilon|(B).$$

Hence with (4.13)

$$\begin{aligned} \int_B \chi_R \left(e_\varepsilon - \frac{1}{2} |\log \varepsilon| d\nu_\varepsilon \right) &= \int_B \chi_R (e_\varepsilon - \Lambda_\varepsilon^{\alpha, \eta/2} d\nu_\varepsilon) + \left(\frac{1}{2} \log \eta + \Delta \right) \int_B \chi_R d\nu_\varepsilon \\ &\geq -C |\nu_\varepsilon|(B) + \frac{\log \eta}{2} \chi_R(a_i) \nu_\varepsilon(B) - \frac{\eta}{2} \log \eta |\nu_\varepsilon|(B) - |\Delta| |\nu_\varepsilon|(B) \\ &\geq \frac{\log \eta}{2} \nu_\varepsilon(B) \chi_R(a_i) - C |\nu_\varepsilon|(B) (1 + \log n_\alpha). \end{aligned}$$

Summing over $B \in \mathcal{C}_\varepsilon^{i,\alpha}$ and then over α and $i \in J$ we find, since

$$\sum_{B \in \mathcal{C}_\varepsilon^i} \nu_\varepsilon(B) = \nu_\varepsilon(B(a_i, \eta)) \rightarrow \nu(B(a_i, \eta)) = 2\pi d_i,$$

that,

$$\liminf_{\varepsilon \rightarrow 0} \int_{\cup_{i \in J} B(a_i, \eta)} \chi_R \left(e_\varepsilon - \frac{1}{2} |\log \varepsilon| d\nu_\varepsilon \right) \geq \pi \sum_{i \in J} d_i \chi_R(a_i) \log \eta - C\Delta(R),$$

where

$$\Delta(R) = \limsup_{\varepsilon \rightarrow 0} \sum_{U_\alpha \subset \mathbf{U}_{R+C} \setminus \mathbf{U}_{R-C}} n_\alpha (\log n_\alpha + 1).$$

Summing (4.12) over $i \in I$ and adding the above and (4.9)–(4.10), we deduce

$$(4.14) \quad \liminf_{\varepsilon \rightarrow 0} \int \chi_R \left(e_\varepsilon - \frac{1}{2} |\log \varepsilon| d\nu_\varepsilon \right) \geq \frac{1}{2} \int_{\mathbb{R}^2 \setminus \cup_{p \in \Lambda} B(p, \eta)} \chi_R (|j|^2 + h^2) \\ + \sum_{i \in I} \chi_R(a_i) (\pi d_i \log \eta + C_{d_i}) + \sum_{i \in J} \chi_R(a_i) \pi d_i \log \eta - C\Delta(R) - o_\eta(1).$$

We will now take the limit $\eta \rightarrow 0$ on the right-hand side. For that we use a Hodge decomposition of j in $B(a_i, \eta_0)$, writing $j = -\nabla^\perp H + \nabla K$, with $H = 0$ on $\partial B(a_i, \eta_0)$. Then since $-\Delta H = \nu - h = 2\pi d_i \delta_{a_i} - 1$ we have $H(x) = d_i \log |x - a_i| + F$, where F is in H^2 in the neighbourhood of a_i , in particular $H \in W^{1,p}$ for any $p < 2$, and since $j \in L^p$, this implies that $K \in W^{1,p}$ also. Then an easy computation shows that

$$\lim_{\eta \rightarrow 0} \frac{1}{2} \int_{B(a_i, \eta_0) \setminus B(a_i, \eta)} \chi_R |\nabla^\perp H|^2 + \pi (\log \eta) d_i^2 \chi_R(a_i)$$

exists and is finite, while

$$\int_{B(a_i, \eta_0) \setminus B(a_i, \eta)} \chi_R |j|^2 \geq \int_{B(a_i, \eta_0) \setminus B(a_i, \eta)} \chi_R (|\nabla^\perp H|^2 + \nabla^\perp H \cdot \nabla K).$$

Decomposing H and integrating by parts we have, writing $C_{i,\eta} = B(a_i, \eta_0) \setminus B(a_i, \eta)$,

$$\int_{C_{i,\eta}} \nabla^\perp H \cdot (\chi_R \nabla K) = \int_{C_{i,\eta}} \nabla^\perp F \cdot (\chi_R \nabla K) - d_i \int_{C_{i,\eta}} K \nabla^\perp \log \cdot \nabla \chi_R,$$

and this remains bounded as $\eta \rightarrow 0$, using the regularity of χ_R , F , and the boundedness of H , K , \log in $W^{1,p}$. We may then deduce that

$$\liminf_{\eta \rightarrow 0} \frac{1}{2} \int_{B(a_i, \eta_0) \setminus B(a_i, \eta)} \chi_R |j|^2 + \pi (\log \eta) d_i^2 \chi_R(p)$$

is not equal to $-\infty$.

As a consequence, writing $d_i = d_i^2 - (d_i^2 - d_i)$ in the right-hand side of (4.14), and this right-hand side being bounded above independently of η , we have that $\sum_i (d_i^2 - d_i) \chi_R(a_i) \log \frac{1}{\eta}$ is bounded above as $\eta \rightarrow 0$. Thus we have $d_i \in \{0, 1\}$ for any i such that $\chi_R(a_i) \neq 0$ and then $d_i = 1$ since d_i was assumed to be nonzero. In view of this, (4.14) can be rewritten as

$$\liminf_{\varepsilon \rightarrow 0} \int \chi_R(e_\varepsilon - \frac{1}{2} |\log \varepsilon| d\nu_\varepsilon) \geq \frac{1}{2} \int_{\mathbb{R}^2 \setminus \cup_{p \in \Lambda} B(p, \eta)} \chi_R(|j|^2 + h^2) + \sum_{p \in \Lambda} \chi_R(p) (\pi \log \eta + \gamma) - C\Delta(R) - o_\eta(1),$$

where we recall that $\gamma = C_1$ and we have absorbed $C_1 \sum_{i \in J} \chi_R(a_i)$ in $C\Delta(R)$.

Letting $\eta \rightarrow 0$ we thus find (see (1.15))

$$\liminf_{\varepsilon \rightarrow 0} \int \chi_R(e_\varepsilon - \frac{1}{2} |\log \varepsilon| d\nu_\varepsilon) \geq W(j, \chi_R) + \frac{1}{2} \int \chi_R h^2 + \sum_{p \in \Lambda} \chi_R(p) \gamma - C\Delta(R).$$

From (4.4) we may replace $e_\varepsilon - \frac{1}{2} |\log \varepsilon| \nu_\varepsilon$ by g_ε , with an error term which may be absorbed in $C\Delta(R)$ hence

$$(4.15) \quad \liminf_{\varepsilon \rightarrow 0} \int \chi_R dg_\varepsilon \geq W(j, \chi_R) + \frac{1}{2} \int \chi_R h^2 + \sum_{p \in \Lambda} \chi_R(p) \gamma - C\Delta(R).$$

Now, under hypothesis (1.17) and using (4.2), we have

$$\limsup_{\varepsilon \rightarrow 0} \sum_{\alpha | A_\alpha \subset \mathbf{U}_R} n_\alpha^2 \leq CR^2$$

and thus

$$\limsup_{R \rightarrow \infty} \limsup_{\varepsilon \rightarrow 0} \frac{1}{R^2} \sum_{\alpha | A_\alpha \subset \mathbf{U}_{R+C} \setminus \mathbf{U}_{R-C}} n_\alpha^\alpha |\log n_\alpha^\alpha| = 0.$$

Indeed, using Hölder's inequality, and bounding the number of α 's involved in the above sum by CR , we find

$$\sum_{\alpha | A_\alpha \subset \mathbf{U}_{R+C} \setminus \mathbf{U}_{R-C}} [n_\alpha^\alpha]^{3/2} \leq (CR)^{1/4} \left(\sum_{\alpha | A_\alpha \subset \mathbf{U}_{R+C}} n_\alpha^2 \right)^{3/4} \leq CR^{1/4+3/2}.$$

It follows, since $U_\alpha \subset A_\alpha$, that

$$(4.16) \quad \limsup_{R \rightarrow +\infty} \frac{\Delta(R)}{R^2} = 0$$

and in particular $\nu(\mathbf{U}_{R+C} \setminus \mathbf{U}_{R-C}) = o(R^2)$. Then we write, using $\nu = \text{curl } j + h$,

$$\sum_{p \in \Lambda} \chi_R(p) = \frac{1}{2\pi} \int \chi_R d\nu = \frac{1}{2\pi} \int \chi_R h - \frac{1}{2\pi} \int \nabla^\perp \chi_R \cdot j.$$

Let $E_R = \{0 < \chi_R < 1\}$. Then since $E_R \subset \mathbf{U}_{R+C} \setminus \mathbf{U}_{R-C}$ we have $|E_R| \leq CR$ and using (4.8) together with Hilbert's inequality we find

$$\int_{E_R} \chi_R h \leq |E_R|^{\frac{1}{2}} \left(\int_{E_R} h^2 \right)^{\frac{1}{2}} \leq CR^{3/2},$$

and a similar bound for $\int \nabla^\perp \chi_R \cdot j$ using (4.8) again, since it is equal to $\int_{E_R} \nabla^\perp \chi_R \cdot j$. Therefore

$$\sum_{p \in \Lambda} \chi_R(p) = \frac{1}{2\pi} \int_{\{\chi_R=1\}} h + o(R^2) = \frac{1}{2\pi} \int_{\mathbf{U}_R} h + o(R^2),$$

the second equality being proved again with the help of (4.8) and Hilbert's inequality. Together with (4.16) and (4.15), this proves (1.19).

There remains to prove (4.11). For this it is convenient to blow-up $B(a_i, \eta)$ to the unit ball B_1 . Then (4.11) becomes

$$(4.17) \quad \frac{1}{2} \int_{B_1} \left(|\nabla_B v|^2 + \left| \frac{\text{curl } B}{\eta} \right|^2 + \frac{(1 - |v|^2)^2}{2\varepsilon'^2} \right) \geq \pi |d_i| \log \frac{1}{\varepsilon'} + C_{d_i} + o_{\eta, \varepsilon}(1),$$

where $v(x) = u_\varepsilon(\eta x)$, $B(x) = \eta A_\varepsilon(\eta x)$ and $\eta \varepsilon' = \varepsilon$, so that ε' tends to 0 with ε . Note that (v, B) depends on ε but we omit this in the notation for the rest of the proof.

Since $\text{curl } A_\varepsilon \rightarrow h$ weakly in L^2_{loc} , it follows that $\|\text{curl } B\|_{L^2(B_1)} \leq 2\eta \|\text{curl } A_\varepsilon\|_{L^2(B_\eta)} \leq C\eta$. Then, choosing to work in the gauge $\text{div } B = 0$, $B \cdot \tau = \text{constant}$ on ∂B_1 , we have $\|B\|_{H^1(B_1)} \leq C\eta$. Since $j(u_\varepsilon, A_\varepsilon)$ is bounded in $L^p_{\text{loc}}(\mathbb{R}^2)$ for any $p < 2$, we deduce immediately that $\|j(v, B)\|_{L^p(B_1)} \leq C\eta^{1-2/p}$. But by Sobolev embedding, $\|B\|_{L^q(B_1)} = O(\eta)$ for any $q > 1$ hence the integral of $B \cdot j(v, B)$ on B_1 is $o_\eta(1)$. Then, since

$$|\nabla_B v|^2 = |\nabla v|^2 - 2B \cdot j(v, B) + |B|^2 |v|^2,$$

(4.17) will follow if we show that

$$(4.18) \quad \frac{1}{2} \int_{B_1} \left(|\nabla v|^2 + \frac{(1 - |v|^2)^2}{2\varepsilon'^2} \right) \geq \pi |d_i| \log \frac{1}{\varepsilon'} + C_{d_i} + o_{\eta, \varepsilon}(1).$$

To prove (4.18) we modify B in order for the current to be divergence free: As before we use the Hodge decomposition $j(v) := (iv, \nabla v) = -\nabla^\perp H + \nabla K$ with $H = 0$ on ∂B_1 , and let $\tilde{v} = v e^{-iK}$. Then denoting $e(v)$ the integrand in (4.18) we have

$$e(\tilde{v}) = e(v) - \nabla K \cdot j(v) + \frac{|v|^2}{2} |\nabla K|^2.$$

We replace $j(v) = -\nabla^\perp H + \nabla K$ and note that, integrating by parts, $\nabla K \cdot \nabla^\perp H$ integrates to 0 on B_1 . Therefore

$$\int_{B_1} e(\tilde{v}) = \int_{B_1} \left(e(v) + \left(\frac{|v|^2}{2} - 1 \right) |\nabla K|^2 \right) \leq \int_{B_1} e(v).$$

Thus if we show the lower bound (4.18) for \tilde{v} , then we are done. For this we may assume, without loss of generality, that the upper bound

$$(4.19) \quad \frac{1}{2} \int_{B_1} \left(|\nabla \tilde{v}|^2 + \frac{(1 - |\tilde{v}|^2)^2}{2\varepsilon'^2} \right) \leq \pi |d_i| \log \frac{1}{\varepsilon'} + C_{d_i}$$

holds.

The advantage is that now we have

$$j(\tilde{v}) = -\nabla^\perp H + (1 - |\tilde{v}|^2) \nabla K.$$

But $\lim_{\varepsilon \rightarrow 0} (1 - |v|^2) = 0$ in $L^q(B_1)$ for any $q > 1$, being bounded in L^∞ and tending to 0 in L^2 . Moreover, we have seen that $\|j(v, B)\|_{L^p(B_1)} \leq C\eta^{1-2/p}$, and that $B = O(\eta)$ in every L^p , so

$$(4.20) \quad j(v, B) - j(v) = |v|^2 B = O(\eta)$$

and therefore $j(v) = O(\eta^{1-2/p})$ in L^p , which implies that H and K are $O(\eta^{1-2/p})$ in $W^{1,p}$. It follows from the above that

$$(4.21) \quad j(\tilde{v}) + \nabla^\perp H = o_{\eta,\varepsilon}(1).$$

in $L^p(B_1)$, for every $p < 2$.

Moreover, since $\operatorname{curl} j(u_\varepsilon, A_\varepsilon) + h_\varepsilon \rightarrow 2\pi d_i \delta_{a_i}$ in $W^{-1,p}$ as $\varepsilon \rightarrow 0$, we have that $\operatorname{curl} j(v, B) + \eta \operatorname{curl} B \rightarrow 2\pi d_i \delta_0$. Hence using (4.20) we deduce $-\Delta H = \operatorname{curl} j(v) \rightarrow 2\pi d_i \delta_0 + o_\eta(1)$ as $\varepsilon \rightarrow 0$ in $W^{-1,p}$. Since $H = 0$ on ∂B_1 we then have

$$(4.22) \quad H(x) = -2\pi d_i \log |x| + o_\eta(1)$$

in $W^{1,p}$.

From (4.21), (4.22) we may find radii $\{r_\varepsilon\}_\varepsilon$ such that

$$\text{i) } \lim_{\varepsilon \rightarrow 0} r_\varepsilon = 1, \quad \text{ii) } \|j(\tilde{v}) + \nabla^\perp H\|_{L^p(\partial B_{r_\varepsilon})} = o_{\eta,\varepsilon}(1), \quad \text{iii) } \|H + 2\pi d_i \log\|_{W^{1,p}(\partial B_{r_\varepsilon})} = o_\eta(1).$$

We may further require that $\rho := |\tilde{v}| \rightarrow 1$ uniformly as $\varepsilon \rightarrow 0$ on $\partial B_{r_\varepsilon}$. Indeed from (4.19) we have

$$\frac{1}{2} \int_{B_1} |\nabla \rho|^2 + \frac{1}{2\varepsilon'^2} (1 - \rho^2)^2 \leq C \log \varepsilon'$$

thus a mean value argument easily implies that r_ε may be chosen such that

$$\frac{1}{2} \int_{\partial B_{r_\varepsilon}} |\nabla \rho|^2 + \frac{1}{2\varepsilon'^2} (1 - \rho^2)^2 \leq C(\log \varepsilon')^2.$$

This in turn implies using (5.1) that $\|\rho - 1\|_{L^\infty(\partial B_{r_\varepsilon})} \rightarrow 0$ as $\varepsilon \rightarrow 0$.

Then, writing $\tilde{v} = \rho e^{i\varphi}$, we have $j(\tilde{v}) = \rho^2 \nabla \varphi$ and the above implies that for some $\theta_0 \in \mathbb{R}$,

$$\tilde{v} = (1 + \tilde{\rho}) e^{i(\theta_0 + d_i \theta + \tilde{\varphi})}, \quad \text{where } \|\tilde{\varphi}\|_{W^{1,p}(\partial B_1)} = o_{\eta,\varepsilon}(1) \text{ and } \|\tilde{\rho}\|_{L^\infty(\partial B_1)} = o_\varepsilon(1).$$

Without going into further detail (see for instance [BBH], Chapter VIII), the above implies that

$$\begin{aligned} & \frac{1}{2} \int_{B_1} \left(|\nabla \tilde{v}|^2 + \frac{(1 - |\tilde{v}|^2)^2}{2\varepsilon'^2} \right) \\ & \geq \min \left\{ \frac{1}{2} \int_{B_1} \left(|\nabla u|^2 + \frac{1}{2\varepsilon'^2} (1 - |u|^2)^2 \right) \mid u = e^{id_i\theta} \text{ on } \partial B_1 \right\} + o_{\eta,\varepsilon}(1). \end{aligned}$$

From [BBH], the right-hand side is precisely equal to $\pi|d_i| \log \frac{1}{\varepsilon'} + C_{|d_i|} + o_\varepsilon(1)$, where the constant C_d is equal to γ if $d = 1$. Thus we have proved (4.18), and then (4.11).

5 Proof of Proposition 2.1

The proof of Proposition 2.1 is based on the ball construction of R. Jerrard [Je], hence we will only emphasize the points which need some modification, mostly to take into account the presence of the magnetic potential A the way we do in [SS4]. We will denote by c , C , respectively, a small and a large generic universal constant. We will number the constants we need to keep track of. Throughout this section U is a bounded domain in \mathbb{R}^2 and (u, A) are defined on U .

The first ingredient is a lower bound for the energy of $|u|$ on a circle ([Je] Lemma 2.3). It is valid for any $\varepsilon > 0$.

Lemma 5.1. *Assuming $2r \geq \varepsilon > 0$ and x are such that the closed ball $B(x, r) \subset U$ we have*

$$(5.1) \quad \frac{1}{2} \int_{\partial B(x,r)} |\nabla |u||^2 + \frac{(1 - |u|^2)^2}{2\varepsilon^2} \geq c_0 \frac{(1 - m)^2}{\varepsilon},$$

where $m = \min_{\partial B(x,r)} |u|$.

In contrast to [Je] and because we wish to work with constants independent of U we introduce

$$U_\varepsilon = \{x \in U \mid \text{dist}(x, U^c) > \varepsilon\}.$$

Then $u : U \rightarrow \mathcal{C}$ being given we introduce, following [Je], $S = \{x \in U_\varepsilon \mid |u| \leq 1/2\}$. Assuming u is continuous the connected components of S which are included in U_ε are compact, and $u/|u|$ has a well defined degree, or winding number on their boundary. Then we let

$$S_E = \{\text{Union of the components of } S \text{ with nonzero boundary degree}\}.$$

Still following [Je], for any compact $K \subset U$ such that $\partial K \cap S_E = \emptyset$ we let

$$\text{deg}_E(u, \partial K) = \sum_{S_i \text{ component of } S_E} \text{deg}(u, \partial S_i).$$

Note that this degree is defined even if $|u|$ vanishes on ∂K , provided the points where it vanishes are not in S_E .

The previous lemma implies (see [Je], Proposition 3.3)

Lemma 5.2. *There exists a collection of disjoint closed balls B_1, \dots, B_k of radii r_1, \dots, r_k such that*

1. $\forall i, r_i \geq \varepsilon$
2. $S_E \cap U_\varepsilon \subset \cup_{i=1}^k B_i$
3. $\forall i, e_\varepsilon(U \cap B_i) \geq c_1 r_i / \varepsilon$.

Proof. We only sketch the proof. If $x \in S_E$ then either $\partial B_r(x)$ intersects $\{|u| \leq 1/2\}$ for every $\varepsilon/2 \leq r \leq \varepsilon$ and the above lemma implies that $e_\varepsilon(U \cap B(x, \varepsilon)) \geq c$ or there exists $\varepsilon/2 \leq r \leq \varepsilon$ such that $|u| > 1/2$ on $\partial B_r(x)$ and then the connected component of x in S_E , which has nonzero degree, is included in $B(x, r)$. The nonzero degree implies again (see [Je]) $e_\varepsilon(U \cap B(x, \varepsilon)) \geq c$. We thus have a cover of S_E by balls which satisfy $e_\varepsilon(B) \geq cr(B)/\varepsilon$.

From Besicovitch's Lemma, there exists a disjoint subcollection $\{B_k\}_k$ such that $\{\tilde{B}_k\}_k$ covers S_E , where $\tilde{B}_k = CB_k$ and C is a universal constant. These balls still satisfy $e_\varepsilon(B) \geq cr(B)/\varepsilon$, though with a smaller constant. Then, grouping the balls which intersect in larger ones as in [Je] (see also [SS2]) allows to obtain a disjoint cover of S_E with the same property. Item 1) is trivially verified since the balls we started with had radius ε . Note also that the balls obtained here only depend on S_E hence on u . □

Still following [Je], we have:

Proposition 5.1. *Choose $c_2 \in (0, c_1)$ small enough and let*

$$\lambda_\varepsilon(x) = \min \left(\frac{c_2}{\varepsilon}, \frac{\pi}{x} \frac{1}{1 + \frac{x}{2} + \frac{\pi\varepsilon}{c_0 x}} \right).$$

Then, assuming that $B_r \subset U_\varepsilon$, that $\partial B_r \cap S_E = \emptyset$ and that $\varepsilon \leq r \leq |d|/2$, where $d = \deg_E(u, \partial B_r)$ is assumed to be different from 0, we have

$$(5.2) \quad \frac{1}{2} \int_{\partial B_r} |\nabla_A u|^2 + \frac{1}{2} \int_{B_r} |\operatorname{curl} A|^2 + \frac{1}{4\varepsilon^2} \int_{\partial B_r} (1 - |u|^2)^2 \geq \lambda_\varepsilon \left(\frac{r}{|d|} \right).$$

Moreover, the primitive function $\Lambda_\varepsilon(x) = \int_0^x \lambda_\varepsilon$ is increasing, $s \mapsto \Lambda_\varepsilon(s)/s$ is decreasing,

$$\lim_{s \searrow 0} \frac{\Lambda_\varepsilon(s)}{s} = \frac{\min(c_0, c_2)}{\varepsilon} < \frac{c_1}{\varepsilon}, \quad \frac{\Lambda_\varepsilon(\varepsilon)}{\varepsilon} \geq \frac{c_3}{\varepsilon}$$

and finally, for any $\varepsilon \leq s \leq 1/2$, and for some $C_0 > 0$,

$$(5.3) \quad \Lambda_\varepsilon(s) \geq \pi \log \frac{s}{\varepsilon} - C_0.$$

Proof. First, in the case where ∂B_r intersects $\{|u| \leq 1/2\}$ we deduce from (5.1) that (5.2) is satisfied with $c_2 = c_0/4$.

When on the contrary $|u| > 1/2$ on ∂B_r we have $\deg_E(u, \partial B_r) = \deg(u, \partial B_r)$. Then we bound from below $\frac{1}{2} \int_{\partial B_r} |u|^2 |\nabla \varphi - A|^2$, where $u = |u|e^{i\varphi}$ as follows: Still denoting $m = \min_{\partial B_r} |u|$, using the Cauchy-Schwarz inequality we have

$$\frac{1}{2} \int_{\partial B_r} |u|^2 |\nabla \varphi - A|^2 \geq \frac{m^2}{2} \frac{1}{2\pi r} \left(\int_{\partial B_r} \frac{\partial \varphi}{\partial \tau} - A \cdot \tau \right)^2 = \frac{m^2}{4\pi r} (2\pi d - X)^2$$

where we write $X := \int_{B_r} \operatorname{curl} A = \int_{\partial B_r} A \cdot \tau$. On the other hand, by Cauchy-Schwarz again

$$\frac{1}{2} \int_{B_r} |\operatorname{curl} A|^2 \geq \frac{1}{2\pi r^2} \left(\int_{B_r} \operatorname{curl} A \right)^2 = \frac{X^2}{2\pi r^2}$$

Adding the two relations we obtain

$$\frac{1}{2} \int_{\partial B_r} |u|^2 |\nabla \varphi - A|^2 + \frac{1}{2} \int_{B_r} |\operatorname{curl} A|^2 \geq \frac{1}{2\pi r} \left(\frac{m^2}{2} (2\pi d - X)^2 + \frac{1}{r} X^2 \right).$$

Minimizing the right-hand side with respect to X yields

$$(5.4) \quad \frac{1}{2} \int_{\partial B_r} |u|^2 |\nabla \varphi - A|^2 + \frac{1}{2} \int_{B_r} |\operatorname{curl} A|^2 \geq \frac{\pi d^2}{r} \frac{m^2}{1 + \frac{m^2 r}{2}}.$$

Adding (5.1) we deduce for $r \geq \varepsilon$ that

$$(5.5) \quad e_\varepsilon(\partial B_r) \geq \frac{\pi |d|}{r} \frac{|d|}{\frac{1}{m^2} + \frac{r}{2}} + c_0 \frac{(1-m)^2}{\varepsilon}.$$

If $|d| > 1$, then either $m^2 < 2/3$ and we find $e_\varepsilon > c/\varepsilon$ for a well chosen $c > 0$ or $m^2 \geq 2/3$ and, since $r/2 < |d|/4$, we have $m^{-2} + r/2 \leq 3/2 + |d|/4 \leq |d|$ implying $e_\varepsilon \geq \pi |d|/r$. Thus, if $|d| > 1$, (5.2) is satisfied. If $|d| = 1$ then minimizing the right-hand side of (5.5) with respect to m yields

$$e_\varepsilon(\partial B_r) \geq \frac{\pi}{r} \frac{1}{1 + \frac{r}{2} + \frac{\pi \varepsilon}{c_0 r}},$$

so that in every case we have $e_\varepsilon(\partial B_r) \geq \lambda_\varepsilon(r/|d|)$, if c_2 is chosen small enough.

We now turn to the properties of Λ_ε . Since λ_ε is positive, decreasing, then Λ_ε is increasing and $\Lambda_\varepsilon(s)/s$ is decreasing. It is clear that as $s \rightarrow 0$, we have $\lambda_\varepsilon(s) \sim \min(c_0, c_2)/\varepsilon \sim \Lambda_\varepsilon(s)/s$. Moreover, if $x > c\varepsilon$, with $c = \pi/c_2$, then

$$\lambda_\varepsilon(x) = \frac{\pi}{x} \frac{1}{1 + \frac{x}{2} + \frac{\pi \varepsilon}{c_0 x}}$$

hence, if $s \geq c\varepsilon$,

$$\begin{aligned} \Lambda_\varepsilon(s) &\geq \int_{c\varepsilon}^s \frac{\pi}{x} \frac{1}{1 + \frac{x}{2} + \frac{\pi \varepsilon}{c_0 x}} dx \geq \int_{c\varepsilon}^s \frac{\pi}{x} \left(1 - \frac{x}{2} - \frac{\pi \varepsilon}{c_0 x} \right) dx \\ &\geq \pi \log \frac{s}{\varepsilon} - C_0, \end{aligned}$$

for some constant C_0 . If $s < c\varepsilon$ then the inequality remains true if C_0 is chosen large enough, since $\Lambda_\varepsilon(s) \geq 0$.

Finally, $\Lambda_\varepsilon(\varepsilon) \geq \varepsilon\lambda_\varepsilon(\varepsilon) \geq c_3$, if $c_3 > 0$ is chosen small enough. \square

From there, the ball construction procedure (growing and merging of balls) from [Je] (or see [SS2] Prop 3.1) allows to deduce

Proposition 5.2. *For any $0 < s < 1/2$ there exists a family of disjoint closed balls $\mathcal{B}(s)$ (depending only on u_ε) such that*

1. *The family of balls is monotonic i.e. if $s < t$, we have $\mathcal{B}(s) \subset \mathcal{B}(t)$. Moreover, denoting by $r(B)$ the radius of B , the function $s \rightarrow \sum_{B \in \mathcal{B}(s)} r(B)$ is continuous.*

2. *For any s we have $S_E \subset \mathcal{B}(s)$.*

3. *For any $B \in \mathcal{B}(s)$*

$$e_\varepsilon(U \cap B) \geq r(B) \frac{\Lambda_\varepsilon(s)}{s}.$$

4. *If $B \in \mathcal{B}(s)$ and $B \subset U_\varepsilon$ then, letting $d_B = \deg_E(u_\varepsilon, \partial B)$, we have $r(B) \geq s|d_B|$.*

Proof. We let $\mathcal{B}(s_0)$ be the family of balls given by Lemma 5.2, where we choose s_0 small enough so that items 3 and 4 are satisfied (item 2 obviously is). We let $\mathcal{B}(s) = \mathcal{B}(s_0)$ for every $s \leq s_0$. For $s \geq s_0$ we apply the method of growing and merging of [Je] which we sketch briefly: It consists in continuously increasing the parameter s and at the same time making those balls included in U_ε such that $r(B) = s|d_B|$ grow so that the equality remains satisfied. When balls touch, the parameter s is stopped and the balls are merged into a larger ball with radius the sum of the radii of the merged balls, and this is repeated if the resulting family is still not disjoint. This does not change the total radius and when it is done, i.e. when the family is disjoint again, the increasing of s is resumed, etc... This yields a family of disjoint closed balls which is monotonic, such that $s \rightarrow \sum_{B \in \mathcal{B}(s)} r(B)$ is continuous and such that $r(B) \geq s|d_B|$ for every ball included in U_ε . Obviously $S_E \cap U_\varepsilon \subset \mathcal{B}(s)$ for every s . Also the growing and merging process depends only on the initial balls and the degrees of u_ε , hence on u_ε .

The lower bound $e_\varepsilon(U \cap B) \geq r(B)\Lambda_\varepsilon(s)/s$ is true initially and is preserved through the merging process, it is also preserved through the growing process as long as (5.2) remains valid, i.e. $r(B) < |d_B|/2$ for every $B \subset U_\varepsilon$ such that $d_B \neq 0$. This results from the properties of Λ_ε , as detailed in [Je]. Then for the process to stop, there must be a ball B for which $r(B) = s|d_B|$, i.e. a growing ball, with $r(B) \geq |d_B|/2$, hence we must have $s \geq 1/2$. \square

We may now deduce

Proof of Proposition 2.1. We begin by constructing a family $\mathcal{B}'(s)$ which contains S_E instead of $\{x \in U_\varepsilon \mid |u| \leq 1/2\}$ but satisfies items 2-3 and then modify it.

Under the hypotheses, Proposition 5.2 applies, and yields for every $0 < s < 1/2$ a family of balls $\mathcal{B}'(s)$ satisfying the 4 items stated. Choosing s_0 small enough we have $\Lambda(s_0)/s_0 \geq c/\varepsilon$ hence, letting r_0 denote the total radius of the balls in $\mathcal{B}'(s_0)$,

$$\varepsilon^{-\beta} \geq G_\varepsilon(u, A) \geq \frac{cr_0}{\varepsilon}$$

and therefore $r_0 \leq C\varepsilon^{1-\beta}$.

Let $r \in (C\varepsilon^{1-\beta}, 1/2)$, and let r_1 denote the total radius of the balls in $\mathcal{B}'(1/2)$. If $r > r_1$ then $\mathcal{B}'(1/2)$ satisfies item 2 trivially and moreover for any $B \in \mathcal{B}'(1/2)$ we have from Proposition 5.2 and using (5.3) that

$$e_\varepsilon(B) \geq |d_B| \Lambda_\varepsilon\left(\frac{1}{2}\right) \geq \pi |d_B| \left(\log \frac{1}{2\varepsilon} - C\right) \geq \pi |d_B| \left(\log \frac{r}{\bar{C}_\alpha \varepsilon} - C'\right),$$

for any $r \leq 1/2$ and any $\bar{C}_\alpha \geq 2$, proving item 3 in this case.

If $r < r_1$ then there exists $s \in (s_0, 1/2)$ such that $\mathcal{B}' := \mathcal{B}'(s)$ satisfies $r(\mathcal{B}') = r$. Then item 2 of the proposition is satisfied for this collection. Let us check item 3.

Assume then $e_\varepsilon(\mathcal{B}') \leq \bar{C} \log(r/\varepsilon)$, with $2 \leq \bar{C} \leq (r/\varepsilon)^{\frac{1}{2}}$. We show by contradiction that if M is chosen large enough, then

$$s \geq \frac{r}{M\bar{C}}.$$

Since $e_\varepsilon(\mathcal{B}') \geq r\Lambda_\varepsilon(s)/s$ and since $\Lambda_\varepsilon(s)/s$ is decreasing, if $s < r/(M\bar{C})$ and $r/(M\bar{C}) \leq \frac{1}{2}$ then

$$\bar{C} \log \frac{r}{\varepsilon} \geq M\bar{C} \Lambda_\varepsilon\left(\frac{r}{M\bar{C}}\right) \geq \pi M\bar{C} \log\left(\frac{r}{\varepsilon M\bar{C}}\right) - C_0 M\bar{C}.$$

It follows that

$$(1 - \pi M) \log \frac{r}{\varepsilon} + \pi M \log \bar{C} + \pi M \log M - C_0 M \geq 0,$$

which yields a contradiction for $M = 3/\pi$ and $r \geq C\varepsilon$, with C large enough, recalling that $\bar{C} \leq (r/\varepsilon)^{\frac{1}{2}}$. Therefore $s \geq \pi r/(3\bar{C})$ and then for every $B \in \mathcal{B}'$ such that $B \subset U_\varepsilon$ we have

$$e_\varepsilon(B) \geq r(B) \frac{\Lambda_\varepsilon(s)}{s} \geq |d_B| \Lambda_\varepsilon(s) \geq |d_B| \Lambda_\varepsilon\left(\frac{\pi r}{3\bar{C}}\right),$$

which in view of (5.3) yields $\forall B \in \mathcal{B}'$ such that $B \subset U_\varepsilon$

$$e_\varepsilon(B) \geq \pi |d_B| \left(\log \frac{r}{\varepsilon \bar{C}} - C\right),$$

if C is chosen large enough.

It remains to modify $\mathcal{B}'(s)$ so that $S := \{x \in U_\varepsilon \mid |u| \leq 1/2\} \subset \mathcal{B}(r)$. First we note that a well known application of the coarea formula yields rather easily (see [SS4],

Proposition 4.8) that S can be covered by a collection of disjoint closed balls \mathcal{C} such that $r(\mathcal{C}) \leq C\varepsilon G_\varepsilon \leq C\varepsilon^{1-\beta}$. Then for every s we do the merging of the balls in $\mathcal{C} \cup \mathcal{B}'(s)$ as in the proof of Proposition 5.2 to obtain $\mathcal{B}(s)$. If we chose s such that $r(\mathcal{B}'(s)) = r/2$ with $C\varepsilon^{1-\beta} < r < 1$ and C large enough, then $r(\mathcal{B}(s)) \leq r$ since $r(\mathcal{C}) \leq C\varepsilon^{1-\beta}$. Moreover, if $B \in \mathcal{B}(s)$ is such that $B \subset U_\varepsilon$ then $\deg(u, \partial B)$ is the sum of $\deg_E(u, \partial B')$ for $B' \in \mathcal{B}'(s)$ and $B' \subset B$. Then, if $e_\varepsilon(B) \leq \bar{C} \log(r/2\varepsilon)$ the same bound holds for the B' 's and summing the above lower bounds we find

$$e_\varepsilon(B) \geq \pi |d_B| \left(\log \frac{r}{2\varepsilon \bar{C}} - C \right).$$

Changing the constant C we can get rid of the factor 2 and $\mathcal{B}(s)$ has all the desired properties. \square

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