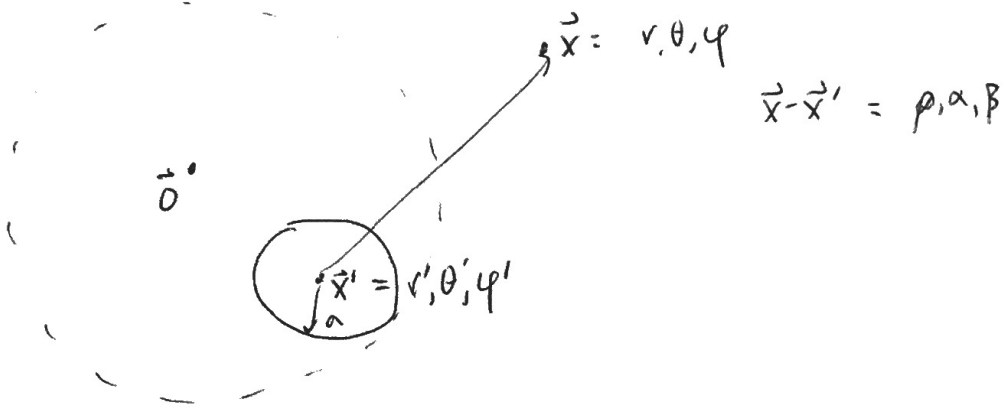


Lecture 4

- Last time: - Higher-order trace code
 - Multiple-to-multiple translations

MZM translation



Then
$$\Phi(\vec{x}) = \sum_{l,m} \frac{M'_{lm} Y_l^m(\theta, \varphi)}{\rho^{l+1}}$$

valid outside $\|\vec{x} - \vec{x}'\| > a$

\Rightarrow

$$\Phi(\vec{x}) = \sum_{l,m} M_{lm} \frac{Y_l^m(\theta, \varphi)}{r^{l+1}} \quad \text{for } \|\vec{x}\| > r' + a$$

and
$$\vec{M}_{\text{new}} = \sum_{l,m} M'_{lm}$$

~~with~~
 T^{cell} ← M2M translation operator defined

by

$$M_{lm} = \sum_{l'=0}^l \sum_{m'=-l'}^{l'} M'_{(m-m')(l-l')} J_{(m-m')(m)} A_{l'm'} A_{(m-m')(l-l')} Y_{l',m'}(\theta, \varphi)$$

$$A_{l',m'}$$

$T_{lm, l'm'}$

Proof substitute in addition formula for $\frac{Y_{l',m'}}{r^{l'}}$, collect terms. Applying T^{cell} gives $O(p^4)$ order.

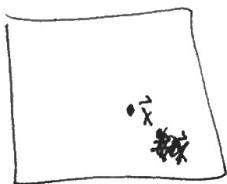
Desc. of algorithm

→ ~~format~~ 0-sort particles, set P.

1- Upward pass: form all p^{th} order multiple expansion for all boxes ~~cost~~ $O(p^4 n)$.

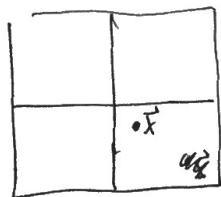
2- Evaluate: ~~for all~~ starting at level 0, evaluate multiple expansion whenever well-separated

Level 0



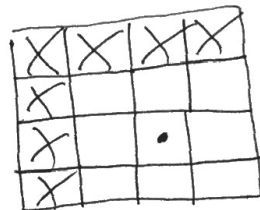
No interactions

Level 1



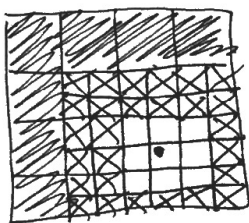
No interactions

Level 2



Exclude MPs from X's

Level 3



Eval MPs from X's.

=> Interaction list
-> "List 2"

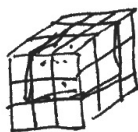
"well-separated boxes that are children of parents' neighbors."

(cf: "List 1" -> self, neighbors)

How many boxes are there in list 2?

2D: parent has 8 neighbors, they have 32 children, 9 of which will be neighbors => 27

In 3D



$3 \times 3 \times 3 = 27$ ~~boxes~~

=> 26 neighbors of parent.

=> $8 \times 26 = 208$ children

How many of the 208 will be neighbors?

=> of 26 neighbors, 7 have same parent,

=> $26 - 7 = 19$ & have other parent

=> $208 - 19 =$ 189

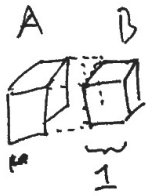
Assume uniform tree, homogeneous dist. of points

Error analysis

No theta

vs. Burns that

MP expansions are evaluated only in boxes that are well-separated:

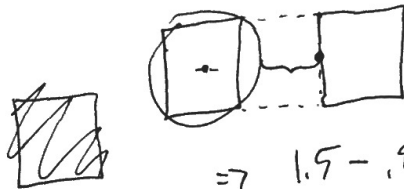


In 3D, for box of width 1,

the smallest sphere enclosing the box

has radius $r = \sqrt{\frac{1}{2^2} + \frac{1}{2^2} + \frac{1}{2^2}} = \sqrt{\frac{3}{4}} = \frac{\sqrt{3}}{2} \approx .87 = a$

The closest point in box B is in the middle of the face:



$\Rightarrow 1.5 - .87 = .63$

$\frac{3}{2} - \frac{\sqrt{3}}{2} = \frac{3-\sqrt{3}}{2} \approx .63$

$\frac{\sqrt{3/2}}{3/2} = \frac{\sqrt{3}}{3}$

~~then~~ $\frac{a}{r}$ in multiple extended ~~to~~

$$\Rightarrow \left| \phi(x) - \sum_{k=0}^p \sum_{n=-l}^l \frac{M_{kn}}{r^{kn}} Y_n^m(\theta, \phi) \right| \leq \frac{\sum |q_j|}{r-a} \left(\frac{a}{r}\right)^{p+1}$$

$$\leq \frac{\sum |q_j|}{.63} \left(\frac{\sqrt{3}}{\sqrt{3}}\right)^{p+1} \leq 2 (\sum |q_j|) (.97)^{p+1}$$

~~$\Rightarrow \log C + (p+1) \log x = \log \epsilon$~~
 ~~$p+1 \approx \log \frac{1}{\epsilon}$~~

If precision ϵ is desired in the calculation, set $p \approx O\left(\log \frac{1}{\epsilon}\right)$. No theta, p controls the accuracy.

Reducing $O(n \log n)$ to $O(n)$.

The "expensive" part of the earlier scheme are

- translation $O(p^4)$
- $\log n$ ~~total~~ MP evals per target.



Introduce Local expansions: The dual of a multipole expansion:

$$\text{Recall } \frac{1}{\|\vec{x} - \vec{x}'\|} = \sum_{l=0}^{\infty} \sum_{m=-l}^l \frac{r'^l}{r^{l+1}} Y_l^m(\theta', \phi') Y_l^m(\theta, \phi)$$

$$= \sum_{l=0}^{\infty} \sum_{m=-l}^l \left(r'^l \frac{Y_l^m(\theta', \phi')}{r^{l+1}} \right) \frac{Y_l^m(\theta, \phi)}{r^{l+1}}$$

↑ target ↑ source



Just flip the source & target:

Likewise, at \vec{x}' , for a unit source located at \vec{x} , $\|\vec{x}\| > \|\vec{x}'\|$

$$\text{we have } \frac{1}{\|\vec{x} - \vec{x}'\|} = \sum_{l=0}^{\infty} \sum_{m=-l}^l \underbrace{\left(\frac{Y_l^m(\theta, \phi)}{r^{l+1}} \right)}_{L_{lm}} r'^l Y_l^m(\theta', \phi')$$

Local Expansion

Similarly:



For source outside radius a ,
at target inside radius a ,

$$\phi(\vec{r}) = \sum_{l=0}^{\infty} \sum_{m=-l}^l L_{lm} r^l Y_l^m(\theta, \varphi)$$

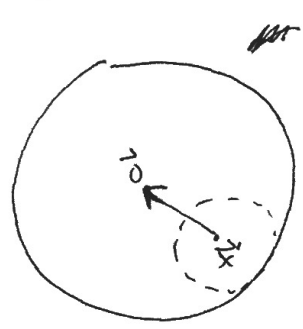
with (by superposition), $L_{lm} = \sum_{j=1}^N q_j \frac{1}{r_j^{l+1}} Y_l^m(\theta_j, \varphi_j)$

and we have the estimate:

$$\left| \phi(\vec{r}) - \sum_{l=0}^P \sum_{m=-l}^l L_{lm} r^l Y_l^m(\theta, \varphi) \right| \leq \frac{\sum |q_j|}{a-r} \left(\frac{r}{a} \right)^{P+1}$$

We will not define the local-to-local translation operators, nor multiply to ~~the~~ local translation operators, but rather only describe.

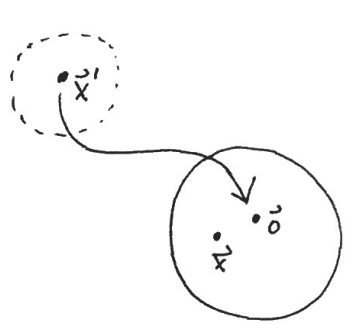
M2M:



$$\sum_{lm} M'_{lm} O_l^m(\rho, \alpha, \beta)$$

$$\rightarrow \sum_{l,m} M_{lm} I_l^m(r, \theta, \varphi)$$

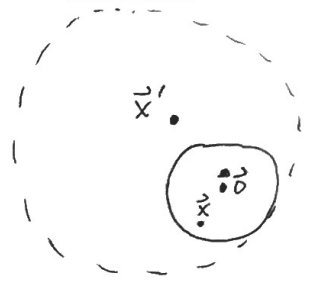
M2L



$$\sum M'_{lm} O_l^m(\rho, \alpha, \beta)$$

$$\rightarrow \sum L_{lm} I_l^m(r, \theta, \varphi)$$

L2L



$$\sum L'_{lm} I_l^m(\rho, \alpha, \beta)$$

$$\rightarrow \sum L_{lm} I_l^m(r, \theta, \varphi)$$

where $I_l^m(r, \theta, \varphi) = r^l Y_l^m(\theta, \varphi) \cdot C_l^m$

$O_l^m(r, \theta, \varphi) = \frac{Y_l^m(\theta, \varphi)}{r^{l+1}} D_l^m$

so that formulas work better (Epton, Dembert)

I_l^m = inner solid harmonics (incoming)

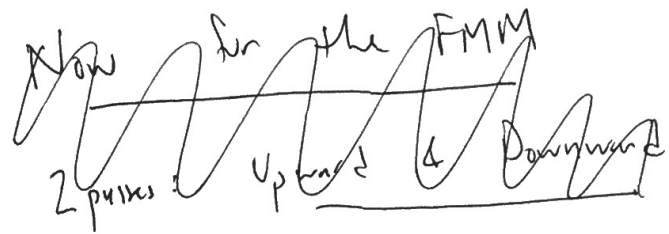
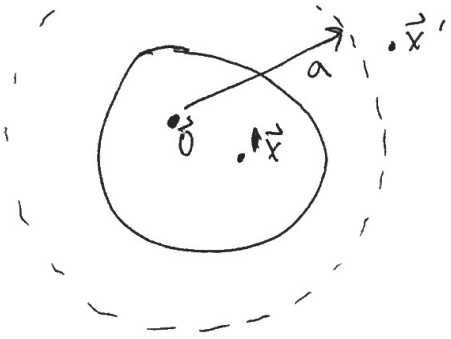
O_l^m = outer solid harmonics (outgoing)

Add in now:

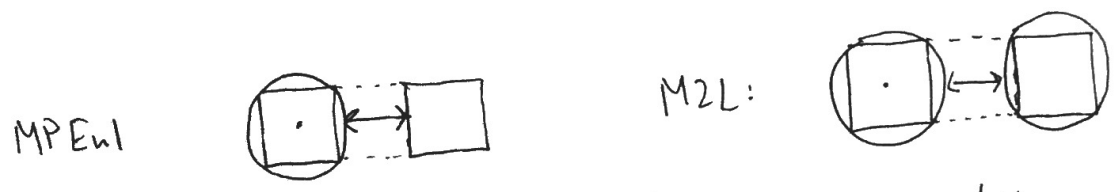
$$\frac{1}{\|\vec{x}-\vec{x}'\|} = \sum_{l,m} (-1)^l I_l^{-m} O_l^m$$

other add'n fins are much simpler.

Local expansions satisfy similar error estimates
as do multiple estimates: $O\left(\left(\frac{r}{a}\right)^{p+1}\right)$

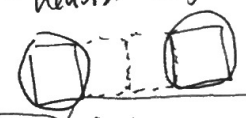


Error estimate for a M2L is worse than for an MP Est:



the separation distance is ~~is~~ less.

For small detail... To improve accuracy, "2nd nearest neighbor"



The fast multiple method:

Discuss: Computational Tradeoff

Upward pass: For levels $L, L-1, \dots, 0$, form the multiple expansion for each box, by translating & merging children expansions.

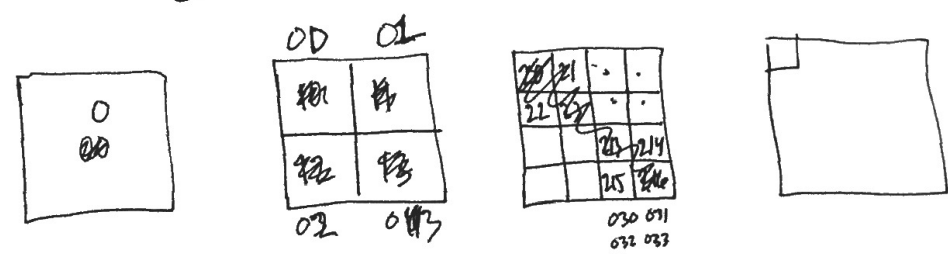
$$O(8p^4 n)$$

Downward pass: Step 1. From levels 0, ..., L, ~~translate~~ for each box, perform multiple to local translation for boxes in List 2. List 2 has 189 boxes, $\Rightarrow O(189p^4n)$

Step 2 From levels 0, ..., L, for each ~~p~~ box, perform L2L from parent (self) to children, add to existing local expansion. $O(8p^4n)$

Eval: For each box on ~~the~~ Level L, evaluate the local expansion and potential due to all neighbors. $O(p^2n)$

The $L=3$



(Levels 0, 1 are numbered).

Form MP: 20, 21, ..., 216

M2M: ~~10, 11, 12, 13, 14, 15, 16, 17, 18, 19, 20~~

(20...23) \rightarrow 10

(213...216) \rightarrow 13

(10...13) \rightarrow 00

Form MP: 0000 \rightarrow 0333 } Level 3

M2M: 0000...0003 \rightarrow 000 } Lev 3 \rightarrow 2
0010...0013 \rightarrow 001

M2L: 000 \rightarrow 010...033 } Lev 2

L2L } Lev 2 \rightarrow 3

Note: L2L or M2M or M2L:

adding multiple expansion means adding the coefficients in the expansions:

M2M



$$\phi = \sum_{lm} M_{lm} O_l^m = \sum_{lm} (\bar{M}_{lm}^1 + \dots + \bar{M}_{lm}^4) O_l^m$$

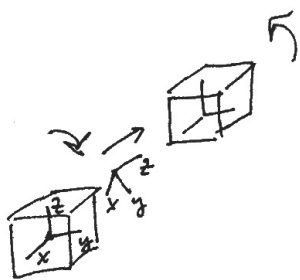
$$\phi_1 = \sum_{lm} M_{lm}^1 O_l^m \rightarrow \phi_1 = \sum_{lm} \bar{M}_{lm}^1 O_l^m$$

$$\vdots$$
$$\phi_4 = \sum_{lm} M_{lm}^4 O_l^m \rightarrow \phi_4 = \sum_{lm} \bar{M}_{lm}^4 O_l^m$$

That's the FMM.

Accelerations: The dominant cost of the FMM is the application of the translation operators. $O(p^4 n)$

Method 1: Point and shoot: $p^4 \rightarrow p^3$



- aim z-axis at target (rotate)
- translate
- rotate back

Procedurally:

We want to construct an operator

which maps:

$$\sum_0^P \sum_{m=-l}^l \alpha_{lm} O_{lm}(r, \theta, \varphi) = \sum_0^P \sum_{m=-l}^l \beta_{lm} \theta_{lm}(r, \theta, \varphi) \quad (*)$$

Why? If shifting along the z-axis (translation), then

Ex: $M_{lm} = \sum_{l'=m}^{\infty(p)} C_{lm}^{l'} \cdot M'_{l'm}$ (i.e. the m-modes don't mix)

Since $\varphi_{old} = \varphi_{new}$.

Don't need $C_{lm}^{l'}$... this means that for each lm , we do $\mathcal{O}(p)$ work $\Rightarrow \mathcal{O}(p^3)$ to translate.

Cost for rotation: Since (*) must hold for all r ,

we need that $\sum_{m=-l}^l \alpha_{lm} Y_l^m(\theta, \varphi) = \sum_{m=-l}^l \beta_{lm} Y_l^m(\theta', \varphi')$

Algorithm 1 (directly, quantum mechanics formulae)